
Advanced NMR Techniques for the Investigation of Small Molecules in Solution

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für meine Familie

I like the cold weather. It means you get work done.

Noam Chomsky

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ABBREVIATIONS

Asp	aspartic acid
<i>ax</i>	<i>axial</i>
BuLi	butyllithium
CDF	combined distribution function
CIP	contact ion pair
CLIP-HSQC	clean-inphase HSQC
COSY	correlation spectroscopy
CPMG	<i>Carr-Purcell-Meiboom-Gill</i>
CpLi	cyclopentadienyllithium
DCM	dichloromethane
DFT	density functional theory
DHPC	dihexanoyl phosphatidylcholine
DMPC	dimyristoyl phosphatidylcholine
DOSY	diffusion ordered spectroscopy
DMG	directing metallation group
DVB	divinylbenzene
EFG	electric field gradient
ESP	electrostatic potential
<i>eq</i>	<i>equatorial</i>
FID	free induction decay
His	histidine
HMQC	heteronuclear multi quantum coherence
HSQC	heteronuclear single quantum coherence
MD	molecular dynamics
MeLi	methyllithium
Me ₂ O	dimethylether
MRI	magnetic resonance imaging
MW	molecular weight
NMR	nuclear magnetic resonance
NOESY	nuclear <i>Overhauser</i> enhanced spectroscopy
OPLS	optimized potential for liquid simulations
PAA	<i>p</i> -azoxyansiol
PAS	principal axis system
PBA	polybutylacrylate
PBBA	<i>p</i> -butoxy-benzoic acid

Abbreviations

PMDTA	<i>N,N,N',N'',N'''</i> -pentamethyldiethylenetriamine
PiLi	piperidinylithium
PS	polystyrene
Q.E.COSY	quadrupolar exclusive COSY
QSC	quadrupolar splitting constant
RCSA	residual chemical shift anisotropy
RDC	residual dipolar coupling
RI	rapid injection
RQC	residual quadrupolar coupling
SAG	strain induced alignment in a gel
SDF	spatial distribution function
SSE	slice-selective excitation
SVD	singular value decomposition
T_1 IR	T_1 inversion recovery
THF	tetrahydrofurane
TMEDA	<i>N,N,N',N'</i> -tetramethylethylenediamine
Tyr	tyrosine
Val	valine
XRD	X-ray diffraction

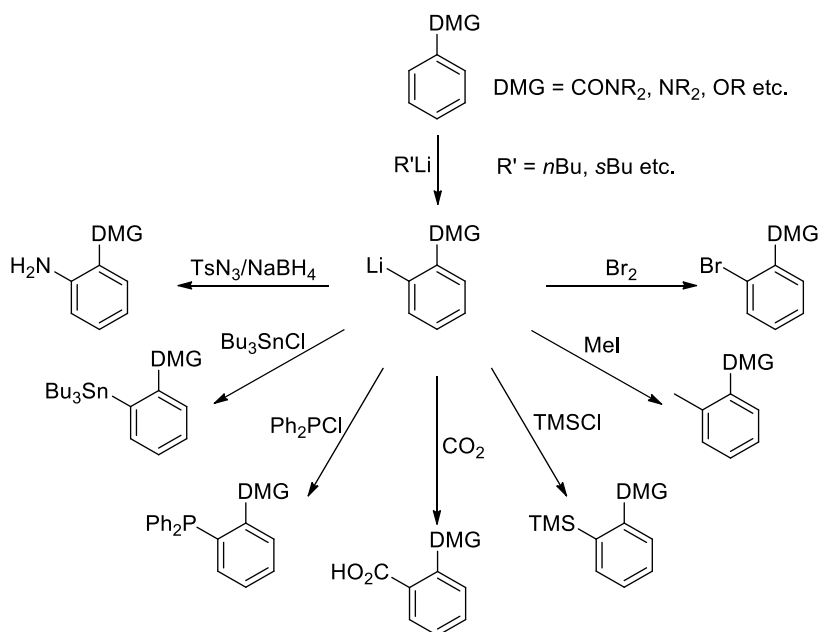
1 INTRODUCTION

High resolution nuclear magnetic resonance (NMR) spectroscopy is arguably the most powerful and versatile method for the observation of molecules in solution. Applications range from detailed structure elucidation to the characterization of dynamic reaction processes with seemingly infinite possibilities in between. This thesis focuses on the development of new methods in NMR spectroscopy of small molecules with an emphasis on lithium organic compounds.

1.1 Lithium Organic Compounds

1.1.1 Applications

The outstanding reactivity of lithium organic compounds is frequently highlighted and requires specific working procedures under inert conditions as introduced by *Schlenk* and *Holtz* in 1917.^[3] This reactivity reflects in the ability to form (thermodynamically challenging) C–C bonds which in turn has expedited research ever since. Nowadays, lithium organics resemble a versatile class of chemical reagents, and applications range from the industrial production of rubber^[4] to modern regio- and stereoselective synthesis of pharmaceuticals.^[5] For instance, the directed *ortho*-lithiation is an omnipresent procedure in synthesis that combines selectivity and reactivity towards a broad field of substrates (Scheme 1-1).



Scheme 1-1: Directed *ortho*-lithiation of an aromatic compound containing a directing metalation group (DMG) and subsequent conversions with $TsN_3/NaBH_4$ ^[6], Bu_3SnCl ^[7], Ph_2PCl ^[7], CO_2 ^[8], $TMSCl$ ^[9], MeI ^[10] and Br_2 ^[11].

Gilman^[12] and *Wittig*^[13] independently discovered that aromatic compounds bearing directing metalation groups (DMGs) such as amides, tertiary amines or ethers can be selectively lithiated in *ortho*

position. These lithiated compounds can be converted in subsequent electrophilic substitution reactions.

Despite their routine usage, lithium organic synthesis protocols are largely based on the empirical determination of reaction behaviour, rather than the derivation of mechanisms from first principles. Yet, rational design of synthetic strategies must be an objective of modern chemical research and can only be achieved by correlating structure and reactivity.

1.1.2 Structure Determination

Structure Determination in Solid State

Comprehensive structure determination in the solid state has long been aggravated by the sensitive, pyrophoric nature of lithium organic compounds. This changed especially with the development of cryo-techniques for single crystal X-ray diffraction (XRD) in the early 90s.^[14] Since then, a plethora of 3D structures with variations in coordination number, ligand identity and molecular symmetry has been obtained which are reviewed in detail elsewhere.^[15] Seminal studies from *Stalke et al.* revealed basic reoccurring motifs that are formed by the lithium atoms.^[16] Most prominent examples are the octahedral *n*-butyllithium ($n\text{BuLi}$)₆ or tetrahedral *tert*-butyllithium ($t\text{BuLi}$)₄ structures (Figure 1-1).

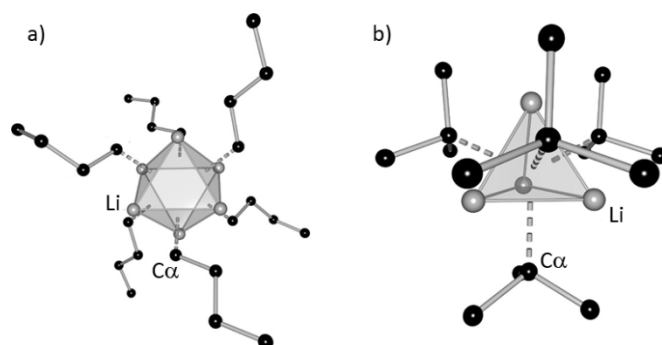


Figure 1-1: Hexameric *n*-butyllithium ($n\text{BuLi}$)₆ (a) and tetrameric *tert*-butyllithium ($t\text{BuLi}$)₄ (b) as basic reoccurring structural motif of lithium organics in the solid state. Hydrogen atoms are omitted for clarity. Reprinted with permission from reference [15c]. Copyright 2014 Wiley VCH.

It is a common perception that the formation of these polyhedral aggregates is largely owed to the ionic character of the Li–C bond.^a Hence, the polarity of solvents exerts a huge influence on the aggregation: it has frequently been shown that charge-stabilizing, polar solvents or specific addition of polydentate, electron-donating *Lewis* bases enforces the disaggregation of polyhedral aggregates into smaller subunits (Figure 1-2).

^a These structural motifs are usually described as contact ion pairs. A further, highly interesting, class of organometallic compounds is classified as solvent separated ion pairs. These compounds are characterized through distinct, solvated cation and anion subunits. Although, these molecules feature a range of interesting properties, this work will not further elaborate on their characteristics and rely on studies given elsewhere.^[17]

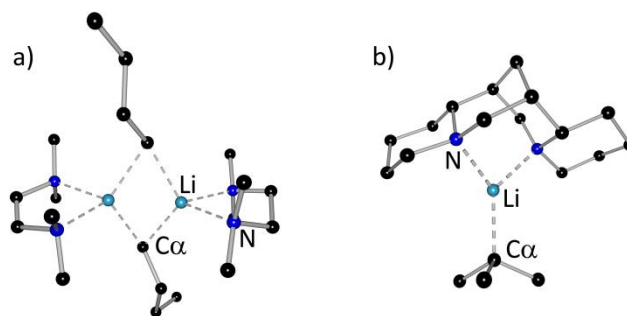


Figure 1-2: Dimeric *n*BuLi-TMEDA aggregate (TMEDA·*n*BuLi)₂ (a) and monomeric *t*BuLi (–)-sparteine aggregate [(–)-sparteine·*t*BuLi] (b). Hydrogen atoms are omitted for clarity. Reprinted with permission from reference [15c]. Copyright 2014 Wiley VCH.

For example, *N,N,N',N'*-tetramethylethylenediamine (TMEDA) allows the preparation of dimeric (TMEDA·*n*BuLi)₂^[18] (Figure 1-2 a), and chiral (–)-sparteine even yields a monomeric complex with *t*BuLi^[19] (Figure 1-2 b). The latter one enables asymmetric reaction steps as required in the synthesis of pharmaceuticals.^[5, 20] Accordingly, it has been well documented that the reactivity of lithium organics increases by switching from apolar solvents (e.g. toluene, *n*-hexane) to polar solvents (e.g. diethylether, tetrahydrofuran) to solvents with additional donor bases (e.g. TMEDA, (–)-sparteine).^[21] From this behaviour a structure-reactivity relationship of lithium organics has been deduced, which states that the increase in reactivity is achieved through disaggregation.^a

Due to adverse molecular properties, single crystals may not always be available. In these cases contributions have perpetually been made by powder XRD; e.g. the tetrameric methyllithium structure (MeLi)₄ was determined by *Weiss et al.* from powder data in 1964.^[23] Similarly, solid state NMR spectroscopy has regularly complemented XRD with a range of valuable structural information (*vide infra*).^[24] Yet, it must be emphasized that chemical syntheses are primarily performed in solution. Since it cannot be assured that structural motifs in solution and solid state coincide, liquid state techniques must support or ensure results from XRD, solid state NMR and likewise.^b

Structure Determination in Solution

It is obvious that the crucial role of solution structure determination for reaction design is widely recognized and has boosted methods development especially in the field of NMR techniques. Among these methods is heteronuclear *Overhauser* spectroscopy which has been employed for the determination of distances between H and Li e.g. by *Bauer*^[26] or *Gschwind*^[27]. Very recently, *Oulyadi* and co-workers showed how scalar Li-Li coupling constants can be extracted and utilized to obtain structural information about mixed lithium structures.^[28] Yet, the most widely used NMR techniques are rapid-injection^[29] (RI) and diffusion-ordered spectroscopy^[30] (DOSY).

^a This concept is challenged with several exceptions.^[22] Usually, these divergent cases can be explained through additional counteracting effects, such as disaggregation and extensive shielding of the polarized carbon lithium bond. Hence, the structure-reactivity relationship of lithium organics still remains a valuable and successfully applied hypothesis.

^b Additionally, state-of-the-art structure elucidation and reactivity studies of lithium organics heavily depend on quantum mechanical calculations. Several excellent reviews focus on this field of research.^[25]

RI NMR is usually performed to deduce the kinetic behaviour of analytes. This technique is based on an apparatus that usually injects the substrate into a solution containing the reagent and achieves a full and homogeneous sample mixing in the timescale of milliseconds.^[29] These rapid injections are instantaneously followed by the acquisition of spectra for a detailed monitoring of the reaction progress between both substrates. Variations in concentration and temperature then allow an ample deduction of the kinetic behaviour of the analytes as detailed in several excellent reviews.^[31]

DOSY enables the highly accurate determination of self-diffusion coefficients D through NMR techniques.^[30] The dependency of diffusion coefficients on molecular hydrodynamic radii is given through the *Stokes-Einstein* equation.^[32] By furthermore correlating the hydrodynamic radii with molecular weights (MW) it is possible to achieve a linear connection between the logarithmic diffusion coefficient $\log(D)$ and the logarithmic molecular weight $\log(MW)$.^[33] Therefore, preparation of calibration curves from model compounds enables *DOSY MW estimation* of unknown analytes from their diffusion coefficients (Figure 1-3).

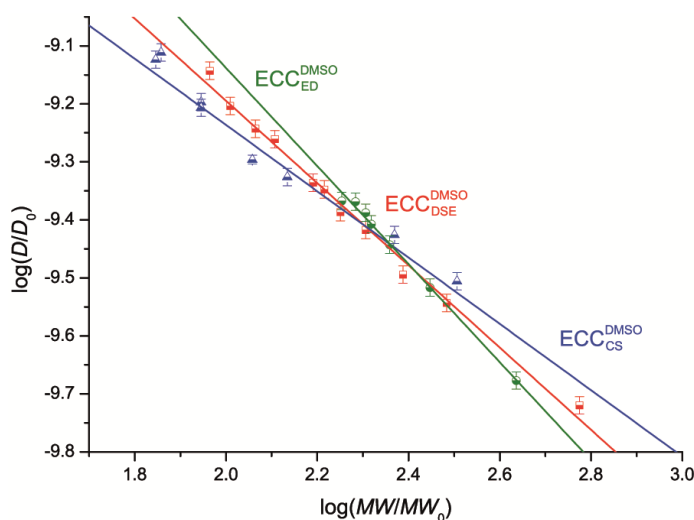


Figure 1-3: Straight lines represent the linear correlation between logarithmic diffusion coefficient $\log(D/D_0)$ as determined from DOSY experiments and logarithmic molecular weight $\log(MW/MW_0)$ of a range of simple organic molecules in dimethylsulfoxide (DMSO). These external calibration curves (ECCs) were additionally categorized according to the shape of the molecules that are employed: Expanded discs (ED, green), dissipated spheres and ellipsoids (DSE, red) and compact spheres (CS, blue). *Reprint with permission from reference [34]. Copyright 2016 Wiley VCH.*

These *estimated weights* can then be compared with *hypothetical weights* of a range of possible aggregates. Hence, differentiation e.g. between monomeric and dimeric species can be made with ease.^[35] Recently, *Neufeld et al.* and *Bachmann et al.*^[34] have made significant advancements through adapted, universally applicable calibration curves (Figure 1-3).^[36]

Although these techniques have proven to be highly successful, there is a constant challenge to develop further elaborate NMR techniques and overcome remaining limitations. For instance, various possible aggregates (structural proposals generated from combinations of lithium organic compounds and coordinating donor bases) may have very similar MW s and are (unlike monomeric and dimeric species) hardly distinguishable via DOSY MW estimation. In these cases complementary (NMR) techniques

have to provide further evidence on structural features to verify/falsify the proposals. Among these methods are slice-selective excitation^[37] (SSE) and anisotropic NMR spectroscopy^[38], which are detailed in the following section.

1.2 Slice-Selective Excitation

1.2.1 Principles

SSE is an NMR technique to obtain spatial resolution while acquiring high-resolution spectra. It can be readily illustrated through the *Zeeman* splitting of nuclei with spin $I = 1/2$ in the presence of a locally dependent magnetic field \vec{B}_z (Figure 1-4).

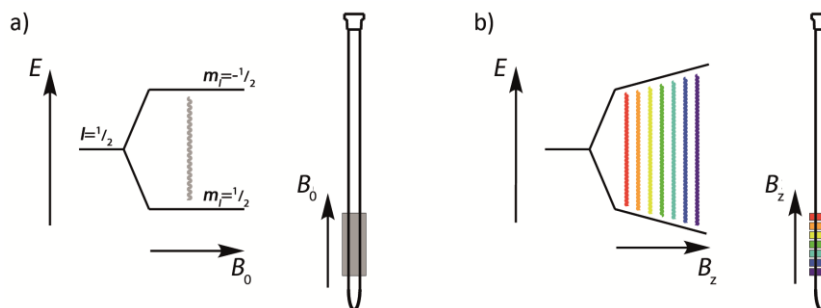


Figure 1-4: a) Zeeman splitting of a spin $1/2$ nucleus in a constant magnetic field B_0 leads to the excitation of the whole active sample volume (grey) via a hard pulse. b) A spatially dependent magnetic field B_z and frequency-shifted soft pulses allow the excitation of discrete slices (colored).

Figure 1-4 a) shows that in a static magnetic field \vec{B}_0 a single *hard pulse* with matching frequency excites the full active volume of the sample. The application of a field gradient along z (Figure 1-4 b) leads to a locally dependent magnetic field \vec{B}_z which in turn causes a locally dependent *Zeeman* splitting. This enables excitation of sample slices through shaped *soft pulses* with varying frequencies that match the energy splitting of the desired slice.

The basic SSE pulse sequence is simple and robust and therefore easily implemented for standard liquid state NMR instruments (Figure 1-5).

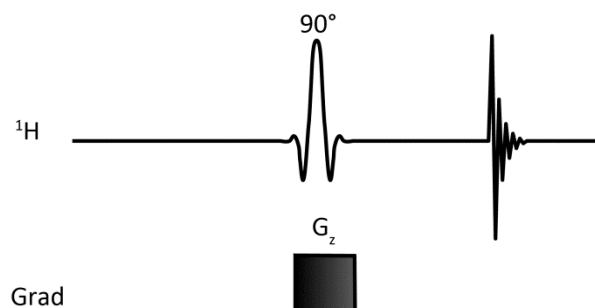


Figure 1-5: Basic SSE pulse sequence consisting of a soft 90° shaped pulse that is delivered simultaneously with a pulsed gradient G_z .

The pulse program is composed of a frequency-encoding pulsed magnetic field gradient G_z and a simultaneous soft 90° shaped pulse followed by the acquisition of the signal. The strength of the gradient G_z effects the spatially dependent magnetic field B_z . Thus, pulses with frequency offsets ν_z excite nuclei with gyromagnetic ratio γ of slices z where $z' = 0$ mm is at the centre of the gradient coil (1-1).

$$\nu_z = \frac{\gamma \cdot G_z \cdot z'}{2\pi} \quad (1-1)$$

Negative z' refer to the slices below and positive z' to the slice above. Note that in the following chapters z substitutes z' where $z = 0$ mm refers to the topmost slice and slices below are numbered consecutively. Additionally, ν_z was incremented in steps to excite slices at 0 mm, 1 mm, 2 mm, etc. for convenience, these slices are referred to as slice 0, 1, 2, etc. i.e. metric unit can be dropped. Besides frequency offsets that determine the spatial location of the centre of the slice, it is important to determine the thickness of a slice Δz to achieve sufficiently intense signals while simultaneously avoiding an overlap of slices. Δz is given by equation (1-2):

$$\Delta z = \frac{2\pi \cdot \Delta\nu_{BW}}{\gamma G_z} \quad (1-2)$$

$\Delta\nu_{BW}$ is the excitation bandwidth of the soft pulse which depends on the individual pulse shape and pulse duration. This also illustrates that a hard pulse is not appropriate as it exhibits a too wide excitation bandwidth for selective excitation (at standard gradient strengths). Usually, simulation protocols as included in NMR software suites (e.g. *nmrsim* or *stdisp* in Bruker Topspin) are applied to calculate these excitation bandwidths.

1.2.2 Applications

Magnetic resonance imaging (MRI) as known from medicinal diagnostics and research resembles the most prominent application of spatially resolved NMR spectroscopy.^[39] While the technical details of modern MRI get increasingly complex and require elaborate theoretical treatment, its basic principle can be transferred from SSE. The spatial selection of slices is extended by a phase and frequency encoding gradient that enables 3D imaging (Figure 1-6).

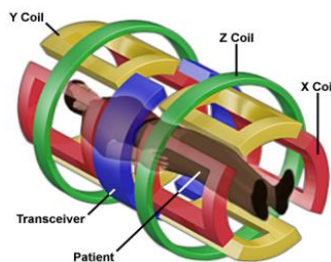


Figure 1-6: Schematic setup of an MRI scanner with a transmitter/receiver (transceiver) unit and three magnet coils (x , y and z) that encode frequencies in 3D space. Reprint from reference [40]. Copyright 2015 National High Magnetic Field Laboratory.

Note that magnetic information as acquired from MRI is usually restricted to relaxation rates and signal intensities of the H₂O resonance.^[39] In contrast, liquid state NMR spectroscopy allows slice-selection combined with the acquisition of high resolution data that includes information like the chemical shift and *J*-couplings. In recent years, this approach led to a number of surprisingly diverse experiments in which SSE has been utilized to

- i) reduce recycle delays and accelerate acquisition,
- ii) achieve homo-decoupling,
- iii) investigate heterogeneous samples along their spatial dimensions.

i) Reduction of recycle delays and acceleration of acquisition

Wagner et al. have shown that a fast reaction can be observed without inter-scan delays by consecutively exciting nuclei in different slices.^[41] This means that after successive acquisitions of signals (no *dead-time*) from different slices, full longitudinal magnetization of the *initial* slices has usually been reached again and the scans can be directly repeated to ensure an ample observation of the reaction.

This method is closely related to the reduction of recycle delays in order to decrease the overall acquisition time of 2D experiments as has been done in correlation spectroscopy (COSY) or^[42] heteronuclear multi quantum coherence (HMQC)^[43] spectroscopy. Furthermore, especially *Frydman* has developed elaborate pulse sequences in which *t*₁ increments are detected along the spatial dimensions of the sample in *one* instance to allow *single-shot (ultra-rapid)* 2D spectroscopy such as acquisition of a COSY spectrum in 0.22 s.^[44]

ii) Homo-decoupling

Very recently, the suppression of proton-proton couplings has been gaining importance in NMR based research.^[45] In crowded spectra of mixtures^[46] or biomolecules homo-decoupling converts overlapping multiplets into distinct interpretable signals (Figure 1-7).^[47]

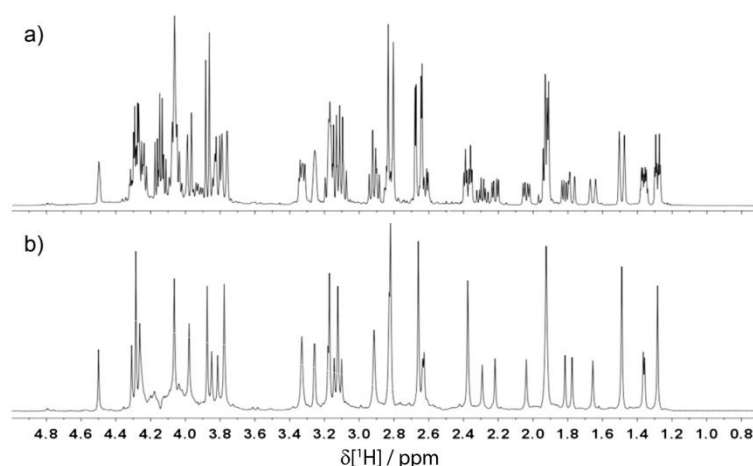


Figure 1-7: A regular a) and slice-selectively homo-decoupled spectrum b) of a mixture of strychnine and its degradation product. *Reprint with permission from reference [47a]. Copyright 2015 Elsevier B. V.*

Hence, so-called pure shift NMR spectroscopy virtually circumvents the resolution restriction of NMR which is otherwise only achieved by increasing the (expensive) magnetic strength of NMR

spectrometers.^[48] Homo-decoupling can be achieved through several types of experiments as detailed elsewhere.^[47] One approach that depends on slice-selection has been pointed out by *Zangger* and *Sterk*^[49], picked up by the research community and driven further developments. Nowadays, homo-decoupling has been combined with frequently used (pseudo) 2D techniques which in turn significantly profit from the improved resolution. Popular examples are pure shift DOSY^[50], pure shift heteronuclear single quantum coherence HSQC^[51], pure shift COSY^[52] and pure shift nuclear *Overhauser* enhanced spectroscopy (NOESY)^[52].

iii) Investigation of heterogeneous samples along their spatial dimensions

The third kind of SSE-based experiments outlined here is closely related to MRI. Within heterogeneous samples slices are excited to correlate chemical information with spatial resolution. In similar studies *Pöppler et al.*^[37] and *Trigo-Mouriño et al.*^[53] observed the distribution of small molecules within crosslinked polymer gel samples (Figure 1-8).

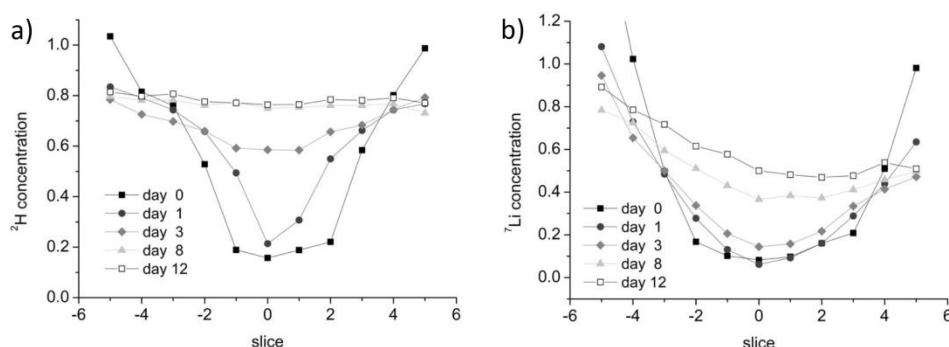


Figure 1-8: ^2H (a) and ^7Li (b) concentration within swollen polystyrene polymer sticks after various days of immersion in solutions of lithiumhexamethyldisilazane in toluene- d_8 . Concentrations were determined from signal intensities in SSE spectra. Reprint with permission from reference [37]. Copyright 2013 Wiley VCH.

Pöppler et al. were able to describe the swelling behaviour and homogeneity of a polystyrene (PS) gel, based on its uptake of solvent (toluene- d_8) and solute (lithium hexamethyldisilazane). Both compounds were characterized by ^2H (Figure 1-8, a) and ^7Li signals (Figure 1-8, b) of sample slices over the course of several days. *Trigo-Mouriño et al.* observed CDCl_3 in polymethylmethacrylate and emphasized the implications of the solvent signal splitting on the shimming procedure.^[53]

Other experiments aimed at the description of biphasic systems via SSE. This led to a remarkable study in which the distribution coefficients of pharmaceutical ingredients (e.g. ibuprofen and paracetamol) in biphasic octanol-water (pH = 7.4) samples were determined and found to be coherent with results from classical time consuming methods like the so-called *shake-flask* technique.^[54] Hence, this application can be readily transferred to pharmaceutical screening processes since distribution coefficients resemble prevalent parameters in toxicity studies.^[55]

It is worth noting that by decreasing the excited volume, magnetic field imperfections along z are compensated.^[56] Hence, SSE is naturally a most convenient technique the examination of heterogeneous samples which exhibit particularly strong inhomogeneity.^[57]

1.3 Anisotropic NMR Spectroscopy

NOESY^[58] and J -coupling analysis^[59] are well established techniques for 3D structure determination. These methods allow the detection of through-space connectivity and provide vital angular information, respectively. Nevertheless, several disadvantages such as the low limit on distances r between spins for the detection of NOE contacts ($\sim 1/r^6$) and the dependency on well-connected spin-systems for J -couplings increase the demand for complementary techniques.

In modern solid state NMR spectroscopy, dipolar^[60] and quadrupolar couplings^[61] supply additional, rich information on the 3D structure of molecules. This information is, at first sight, lost in isotropic liquid state NMR spectroscopy for the benefit of high resolution.^a However, in 1963, *Saupe et al.* showed that dipolar as well as quadrupolar couplings can be residually recovered by partially aligning molecules in so-called uniaxial environments such as solutions of nematic liquid crystals (Figure 1-9).^[63]

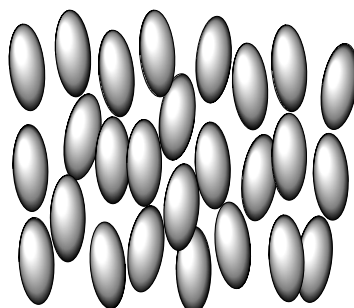


Figure 1-9: Schematic representation of molecules constituting a nematic liquid crystal phase. Note that a net orientation of the *long* molecular axes induces an alignment of potential solutes (framed within the blank space) through hindered rotation.

1.3.1 Residual Dipolar Couplings

Theoretical Basics of (Residual) Dipolar Couplings

The dipolar coupling D between two adjacent nuclei T and S with spin $I = 1/2$ can be described through equation (1-3):

$$D = -\frac{3}{8\pi^2 R^3} \gamma_T \gamma_S \mu_0 \hbar \left(\cos^2 \theta - \frac{1}{3} \right) \quad (1-3)$$

R is the distance between nuclei T and S, γ_T and γ_S are the gyromagnetic ratios of nuclei T and S, respectively, μ_0 is the (magnetic) vacuum permeability, \hbar is the Planck's constant divided by 2π and θ is the angle between the magnetic field vector \vec{B}_0 and the vector connecting S and T (Figure 1-10).^[64]

^a Indeed, *long range* dipolar couplings indirectly become visible in solution as well, since dipolar couplings cause the NOE effect.^[62]

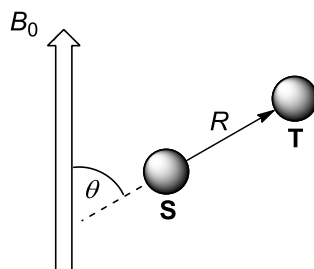


Figure 1-10: The angle θ is defined by the internuclear vector between S and T with length R and the vector for the magnetic field \vec{B}_0 .

Note that D becomes strongest at an angle of 0° , which is similar to bar magnets that adopt a head to tail configuration as the lowest energy configuration (Figure 1-11).

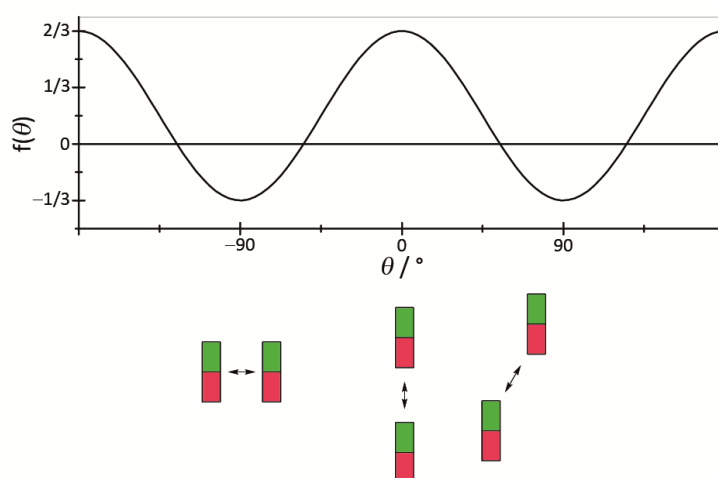


Figure 1-11: The angle term $f(\theta) = \cos^2(\theta) - 1/3$ imposes the limits for the strength of the magnetic dipolar coupling. This is similar to the interaction of bar magnets that show the strongest repulsive interaction at 90° (left) and strongest attractive interaction at 0° (middle). Note that the value of the *magic angle* (right), characterized by the vanishing interaction, depends on the distance and spatial dimensions of the bar magnets. For nuclei in molecules and their magnetic moments the magic angle is at 54.7° .

In contrast, D vanishes to give one sharp signal if θ is at an angle of 54.7° which is called the *magic angle*. A molecule in solution uniformly *tumbles* around the Cartesian x , y , and z axes. The averaged *composite tumbling* of the molecule is then around a room diagonal of a cube which is at 54.7° which in turn leads to the vanishing dipolar coupling.

Dipolar couplings have long been in the focus of solid state NMR spectroscopy and the so-called *Pake* pattern signal of a powdered sample (Figure 1-12) is a prominent trait within a solid state spectrum.^[60]

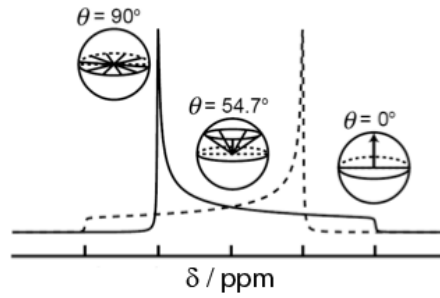


Figure 1-12: Typical *Pake* pattern of a powdered sample containing nuclei with $I = 1/2$. Orientation with $\theta = 0^\circ$ is less probable than $\theta = 90^\circ$. This corresponds to the minimum intensity (“feet”) and the maximum of the signal, respectively.

In a powder θ can attain all values from 0° to 90° but the exact orientation along the magnetic field is statistically less probable (only one exact parallel and one antiparallel orientation) than the perpendicular orientation (all angles between 0° and 360° within the perpendicular plane). Therefore $\theta = 0^\circ$ and $\theta = 90^\circ$ correspond to the *feet* and the maximum of the *Pake* pattern, respectively. The coupling strengths between these features enable the direct determination of the distance between both spins according to equation (1-3).

While a powder exhibits a homogeneous distribution of orientations, uniaxial environments induce physical barriers that lead to an averaged, small, net orientation of (non-spherical) molecules. Hence, the orientation of molecules in these anisotropic media has to be expressed through a 3D probability space which is known as the alignment tensor \mathbf{A} (1-4).

$$\mathbf{A} = \begin{pmatrix} A_{xx} & A_{xy} & A_{xz} \\ A_{yx} & A_{yy} & A_{yz} \\ A_{zx} & A_{zy} & A_{zz} \end{pmatrix} \quad (1-4)$$

\mathbf{A} is a second rank, traceless ($A_{xx} + A_{yy} = -A_{zz}$), symmetric ($A_{ij} = A_{ji}$) tensor and therefore fully described through five independent matrix elements. \mathbf{A} is closely related to the probability tensor \mathbf{P} via equation (1-5):

$$\mathbf{A} = \mathbf{P} - \frac{1}{3} \mathbf{1} \quad (1-5)$$

Here, $\mathbf{1}$ is the identity matrix. In the principal axis system (PAS) of \mathbf{P} (i.e. the diagonalized \mathbf{P}_{PAS} , see equation (1-6)) the principal values $P_{\bar{x}}$, $P_{\bar{y}}$ and $P_{\bar{z}}$ correspond to the physical probabilities to find the magnetic field vector \vec{B}_0 along the associated axis x , y and z , respectively.

$$\mathbf{P}_{\text{PAS}} = \begin{pmatrix} P_{\bar{x}} & 0 & 0 \\ 0 & P_{\bar{y}} & 0 \\ 0 & 0 & P_{\bar{z}} \end{pmatrix} \quad (1-6)$$

Note that these preferences for certain orientations are usually very weak, hence, values for A_i within typical modern alignment media (*vide infra*) are in the orders of 10^{-5} to 10^{-4} . This leads to so-called residual dipolar couplings (RDCs) in the order of several Hz as can be obtained by equation (1-7).

$$\bar{D} = -\frac{3}{8\pi^2 R^3} \gamma_T \gamma_S \mu_0 \hbar (\vec{r}^T \mathbf{A} \vec{r}) \quad (1-7)$$

Here, the angular term of equation (1-3) is substituted by the product of the alignment tensor \mathbf{A} and the unit vector \vec{r} that characterizes the orientation between S and T and its transposed \vec{r}^T . Equation (1-7) is the most fundamental expression from which

- i) RDCs can be (back-)calculated for a known alignment tensor
- ii) the alignment tensor can be derived from a set of known experimental RDCs for a hypothetical structure.^a

Note that the bar above D which signals the averaging of the dipolar coupling to an RDC in anisotropic media is omitted in the subsequent discussion.

Experimental Determination of RDCs

Most commonly, RDCs between directly bound C and H ($^1D_{CH}$) are measured. Furthermore, N and H couplings ($^1D_{NH}$) are frequently observed in the backbone of protein structures.^[65] These one-bond RDCs can be collected by NMR spectroscopic methods in which C–H (N–H etc.) couplings are either recorded in the direct (F2) or indirect (F1) dimension (Figure 1-13).^[66]

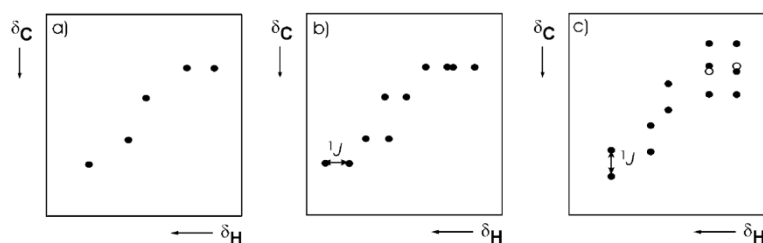


Figure 1-13: Schematic illustration of spectra from HSQC type experiments. a) Decoupled HSQC. b) HSQC coupled in direct (F2) dimension. c) HSQC coupled in indirect (F1) dimension. *Reprint with permission from reference [66]. Copyright 2008 Wiley VCH.*

The observed coupling (signal splitting) $^1T_{XH}$ ($X=C,N$) in anisotropic media is always a sum of the scalar $^1J_{XH}$ and dipolar $^1D_{XH}$ coupling as both inseparably develop in the course of the experiment.

The spectrum of an isotropic sample yields only scalar couplings such as $^1J_{CH}$ (Figure 1-13 a) and it is important to emphasize that the sign of this J -coupling cannot be directly obtained from the HSQC experiment. However, one-bond J -couplings between nuclei that exhibit the same sign for the gyromagnetic ratio are known to be always positive.^[67] This is the case for $^1J_{CH}$. The anisotropic sample yields the combined coupling (Figure 1-13 a and b).^b

^a This corresponds to setting up an equation system from (1-3), with experimental RDCs and the vectors describing C–H bonds in the proposed structure. The latter are then multiplied with the (five) unknown tensor elements of \mathbf{A} . Hence, five linearly independent equations can be used to determine the alignment tensor components.

^b Note that further definitions can be found in the literature with $^1T_{CH} = ^1J_{CH} + 2 \cdot ^1D_{CH}$.^[66] Accordingly, equation (1-7) then contains an additional factor $\frac{1}{2}$. However, from an applied point, this is of minor importance since only ratios of RDCs are employed for further structure analysis.

$${}^1T_{\text{CH}} = {}^1J_{\text{CH}} + {}^1D_{\text{CH}} \quad (1-8)$$

Simple subtraction (${}^1T_{\text{CH}} - {}^1J_{\text{CH}}$) yields the RDC ${}^1D_{\text{CH}}$ with the correct sign (as long as ${}^1D_{\text{CH}} < {}^1J_{\text{CH}}$ since only the magnitude of ${}^1T_{\text{CH}}$ is measured). Acquisition of spectra in F1 and F2 has individual benefits and drawbacks: Most prominently, F2-resolved couplings offer good resolution at short experimental times. However, as remarked above, typical ${}^nJ_{\text{XH}}$ scalar couplings (usually between protons)^a evolve during acquisition as well and may lead to broad signals that complicate the extraction of the dipolar couplings (unless no additional homo-decoupling is applied). In contrast, F1-coupled spectra only exhibit sharp signals in the corresponding dimension since virtually no two ${}^{13}\text{C}$ nuclei are neighbored/scalar coupled at natural abundance. A significant drawback of F1-coupled acquisition is the aggravated extraction of couplings between C and diastereotopic protons which is hinted in (Figure 1-13 c). Small differences in RDCs lead to an overlap and unreliable values.

One of the most prominent types of experiments to achieve F2 coupled spectra is the clean-in-phase (CLIP) HSQC experiment which has been developed by *Luy et al.* (Figure 1-14).^[68]

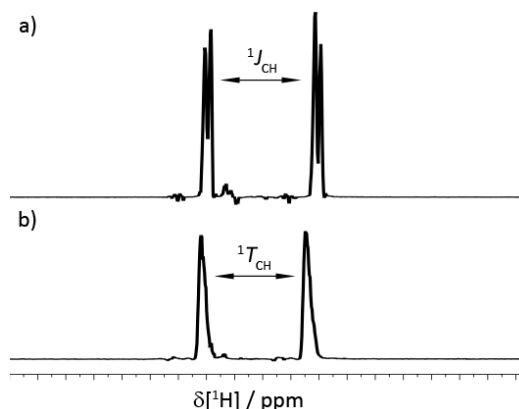


Figure 1-14: Exemplary slices of two CLIP-HSQC spectra for C2-H2 couplings (*ortho* H and carbon) of toluene molecules. a) Isotropic ${}^1J_{\text{CH}}$ coupling in neat dichloromethane. b) Anisotropic ${}^1T_{\text{CH}}$ coupling in crosslinked PS immersed in dichloromethane. Note that the change in chemical shift results predominantly from changes of the chemical environment/susceptibility.

Further experiments to extract long range dipolar couplings ${}^nD_{\text{CH}}$ and ${}^nD_{\text{HH}}$ may add valuable information but require a more elaborate treatment, since they are considerably weaker and the sign of the scalar coupling constant has to be determined.^[69]

Applications of Residual Dipolar Couplings

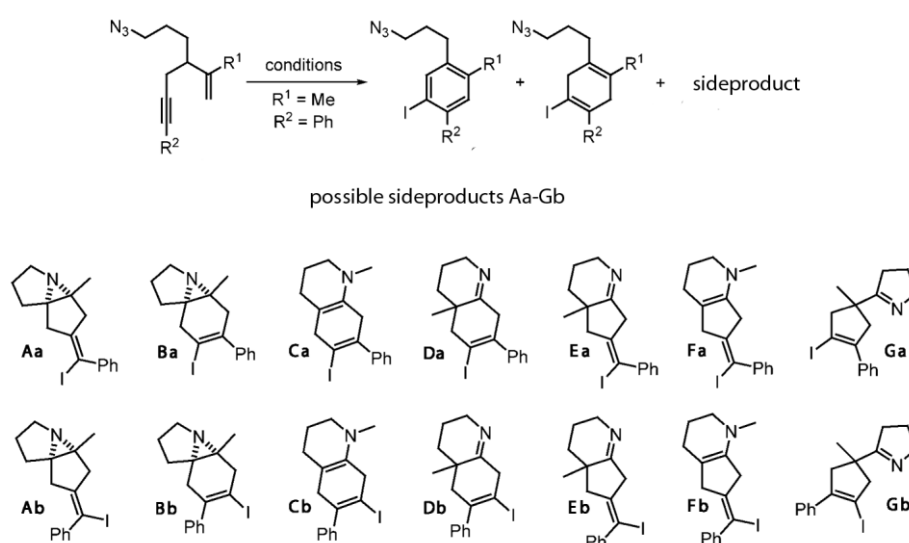
Until today, RDCs have been increasingly employed in the study of large biomolecules^[70] as well as small- to medium-sized synthetic and naturally occurring substances.^[66, 71] Frequent applications include the analysis of relative orientations between molecular fragments^[72], structure refinement of XRD data^[73], the analysis of (constrained) molecular dynamics data^[74] and structure verification^[75]. The first three applications are almost exclusively encountered in biomolecular studies. The most

^a Dipolar couplings between protons may also add to these multiplets.

widespread application of RDCs in studies on small molecules is (3D) structure validation. After experimental acquisition of RDCs as described above, four basic steps constitute this procedure:

- i) Generation of a comprehensive set of possible molecular structures.
- ii) Determination of alignment tensors from RDCs and structural proposals.
- iii) Back-calculation of RDCs from structural proposals and alignment tensors.
- iv) Assessment of the validity of the structural proposal based on experimental and back-calculated RDCs.

In an interesting study *Kummerlöwe et al.* showed how the unknown structure of a sideproduct in an electrophilic cyclization of an azide-containing 1,5-enyne was determined via RDCs.^[76] With only little knowledge restricted to basic building blocks, the authors created a diverse set of 14 trial structures (Scheme 1-2).



Scheme 1-2: Electrophilic cyclization reaction between an azide-containing 1,5-enyne leading to two main products and one unknown sideproduct. 14 possible structures for the sideproduct were created and verified/falsified through RDC analysis. *Reprint with permission from reference [76]. Copyright 2011 Wiley VCH.*

Next, the authors determined the alignment tensor **A** from these structural proposals and experimental RDCs D_{exp} by setting up and solving linear equation systems from (1-7). In turn, the alignment tensor enabled back-calculation of RDCs (D_{calc}). From the highest compliance between D_{exp} and D_{calc} *Kummerlöwe et al.* finally concluded that the reaction led to “the unexpected formation of the tricyclic compound” **Ba** (Scheme 1-2). This structure verification procedure including the setting up of appropriate equation systems is described in more detail in section 4.2.

Note that the simplest approach to utilize RDCs is to estimate relative orientations of prominent molecular features (e.g. helices in biomolecules) from comparison between suitable (parallel) RDCs. However, common RDC based studies heavily rely on the determination of the alignment tensor **A**. This tensor can be either calculated from experimental RDCs and trial structures (as done by *Kummerlöwe et al.* and described in detail in section 4.2) or estimated from software relying on first principles such as the moments of inertia (Tramite)^[77] or the molecular shape combined with simulated steric obstruction (PALES)^[78]. Both methods resemble very intuitive assessments on molecular

orientation either by approximating that molecules in elongated pores preferably rotate around axes with small moments of inertia or that molecular rotation is hindered by a physical wall. Both approaches have been used in biomolecular studies.^[79] For small molecules benchmark molecular dynamics (MD) simulations have been successfully applied to examine their alignment in nematic liquid crystal solutions, but are so far not regularly employed in small molecules structure elucidation.^[80]

However, the ever growing popularity and applicability of RDC based structure determination has enormously capitalized on the development of supportive and comprehensive software. PALES^[78, 81] and REDCAT^[82] for biomolecules and MSpin^[83] for small molecules today initialize RDC analysis implementation in routine NMR.

1.3.2 Residual Quadrupolar Couplings

Theoretical Basics of (Residual) Quadrupolar Couplings

Quadrupolar nuclei possess a nuclear spin with $I > 1/2$ and feature non-spherical charge distributions. The resulting electric quadrupole moment Q (usually denoted in the quantity eQ where e is the elementary charge) interacts with the electric surrounding (neighbouring atoms and electrons) that is characterized by the electric field gradient (EFG). The EFG can be fully described by a second-rank traceless, symmetric tensor \mathbf{V} (1-9).

$$\mathbf{V} = \begin{pmatrix} V_{xx} & V_{xy} & V_{xz} \\ V_{yx} & V_{yy} & V_{yz} \\ V_{zx} & V_{zy} & V_{zz} \end{pmatrix} \quad (1-9)$$

In its principal axis system and with ordering $|V_{\bar{x}}| \leq |V_{\bar{y}}| \leq |V_{\bar{z}}|$ the asymmetry parameter η describes the *cylindricity* of the EFG (1-10):

$$\eta = \frac{V_{\bar{x}} - V_{\bar{y}}}{V_{\bar{z}}} \quad (1-10)$$

The quadrupolar nucleus may be illustrated by an axially symmetric ellipsoid localized within a negative electric charge distribution of surrounding electrons (Figure 1-15 a).

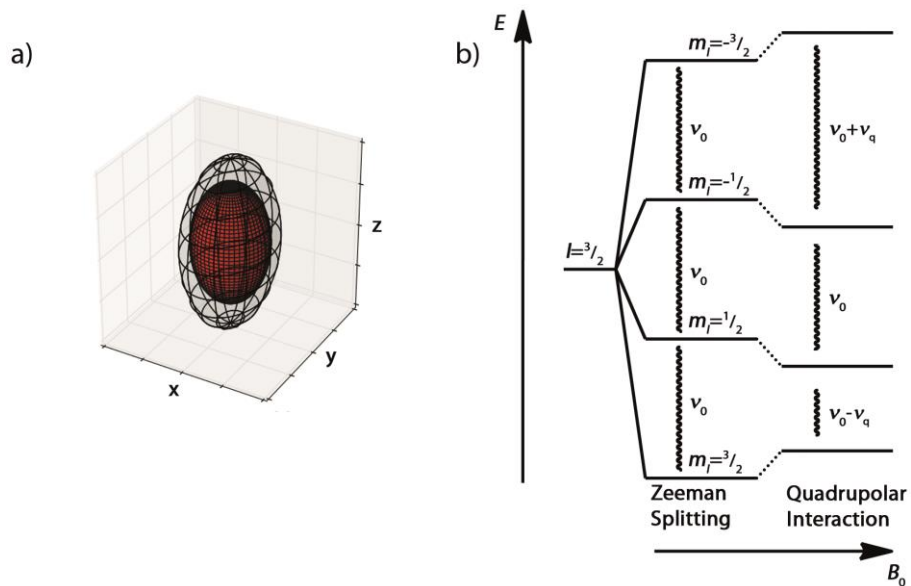


Figure 1-15: a) Schematic illustration of the quadrupolar nucleus (red) within the negative electric charge (colorless). b) The splitting of energy levels of a single nucleus with spin $I = 3/2$ in an external magnetic field (Zeeman splitting) is further modified through the quadrupolar interaction.

The spin-states m_I of a quadrupolar nucleus in a magnetic field experience an energetic attenuation or amplification through interaction with the EFG (Figure 1-15 b). This results in $2I$ spectroscopic transitions with $\Delta m_I = 1$. Hence, for a nucleus with $I = 3/2$ three resonances can be observed: ν_0 is the central transition and $\nu_0 + \nu_q$ as well as $\nu_0 - \nu_q$ are the satellite transitions. According to this, the quadrupolar coupling ν_q is the frequency difference between the central transition and either one of the satellite transitions. The following definition of the quadrupolar coupling and its orientation dependency is adopted from the conventions used by Man:^{[61b]a}

$$\nu_q = \frac{3\chi}{4I(2I-1)} (3\cos^2\beta - 1 + \eta\sin^2\beta\cos 2\alpha) \quad (1-11)$$

Here, χ is the quadrupolar coupling constant, I is the nuclear spin, η is the asymmetry parameter of the EFG and α and β are the Euler angles of the magnetic field vector in the PAS of the EFG (for a detailed definition and discussion of these angles see reference [84]). The quadrupolar coupling constant χ is usually defined by equation (1-12).

$$\chi = \frac{e^2qQ}{h} \quad (1-12)$$

Here, eQ is the nuclear electric quadrupole moment, eq is the largest component of the EFG tensor (V_z) and h is the Planck's constant. Usually, χ and η can be readily determined from solid state powder spectra as depicted in Figure 1-16.

^a Note that further definitions can be found in the literature with an additional factor $1/2$. In these cases the satellite transitions in Figure 1-15 b) exhibit frequencies $\nu_0 \pm 2\nu_q$.

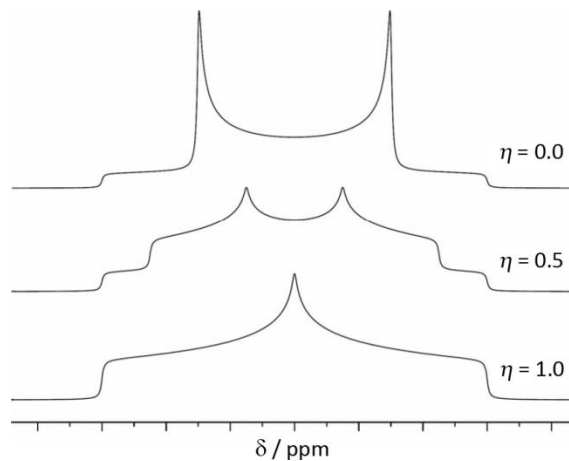


Figure 1-16: Typical powder pattern of a sample containing nuclei with $I = 1$ for EFGs with increasing asymmetry parameter η . Reprint with permission from reference [61c]. Copyright 2013 Wiley VCH.

In the case of no asymmetry ($\eta=0$) the EFG can be treated like a vector and is analogous to dipolar couplings. χ can then be directly read of the difference between the maxima of the peak or its *feet*. Increasing asymmetry is usually well determined from fitting procedures and comparison of frequency differences between feet and maxima.

In the case of partial alignment, no continuous distribution of orientations/frequencies is observed, but an averaged partial alignment from which a number of satellite resonances result according to the nuclear spin I . The EFG tensor elements contribute to these frequency shifts but are modulated by their probability of being aligned along the magnetic field vector. This relationship defines the RQC ν_{rq} as given in equation (1-13):

$$\nu_{\text{rq}} = \frac{3eQ}{2I(2I-1)\hbar} \sum_{i,j=x,y,z} V_{ij}A_{ij} \quad (1-13)$$

The equation contains a set of constants already known from above. Additionally, eQ , η and the angular term as included in (1-11) are substituted by a sum over all Cartesian products V_{ij} and A_{ij} which are elements of the EFG and the alignment tensor, respectively. Note the distinction between this Cartesian definition of the subscripts i and j (i.e. a shared reference frame for both tensors such as a molecular fixed frame or either one of the PAS of the tensors) and the usual conventions for naming and ordering the PAS elements of both tensors. This may easily lead to confusion and misunderstandings but will be acknowledged specifically at the appropriate passages of this thesis.

Applications of (Residual) Quadrupolar Couplings

The majority of studies on quadrupolar nuclei employs solid state NMR spectroscopy and subsequent lineshape analysis of e.g. satellite transitions. Suitable simulation software like Dmfit (QUASAR)^[85] or SIMPSON^[86] allow straightforward determination of eq (χ) and η (i.e. the EFG \mathbf{V}) from fitting simulated to experimental spectra. Corresponding techniques and detailed applications have been extensively reviewed elsewhere.^[61] Notably, within the group of *Günther* various studies on ⁷Li-containing compounds have been conducted and dependencies between the quadrupolar coupling constant $\chi(^7\text{Li})$ and angles X-Li-X (X = C, N) were derived (Figure 1-17).^[87]

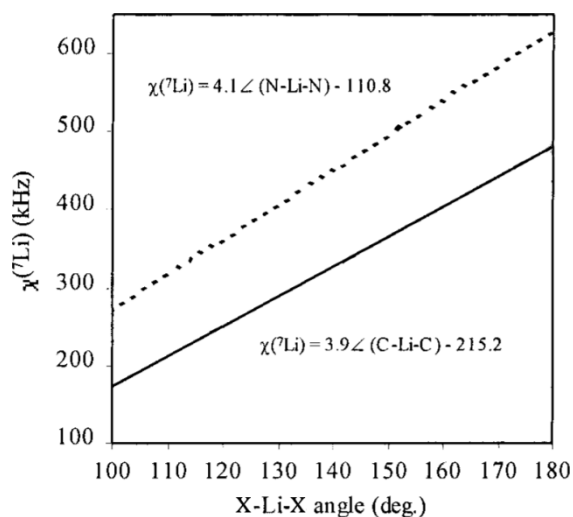


Figure 1-17: Correlation between the quadrupolar coupling constant $\chi(^7\text{Li})$ as determined from solid state NMR and the structural N–Li–N angle (dotted) or C–Li–C angle (solid) as determined from X-ray data. *Reprinted with permission from reference [87b]. Copyright 2000 American Chemical Society.*

These correlations may offer viable structural information for molecular compounds that do not crystallize. In contrast to RDCs, applications of RQCs are limited due to adverse atomic properties. In particular, strong quadrupolar moments induce fast relaxation and linebroadening and thereby prohibit the observation of RQCs in cases such as ⁷⁵As with $I = 3/2$ and $Q = 31.5 \text{ fm}^2$ which is about eight times higher than for ⁷Li ($Q = -4.01 \text{ fm}^2$). Nevertheless, the applicability of RQCs in anisotropic solutions was recognized early on together with upcoming RDC techniques. Until today, especially the use of ²H spectroscopy ($I = 1$; $Q = 0.286 \text{ fm}^2$) has been attracting some attention since it can be employed for enantiomeric differentiation of deuterated solutes in chiral nematic liquid crystals.^[88] In these cases deuterons are assumed to experience a nearly axially symmetric EFG with $|V_{zz}|$ along the sigma bonds (similar to RDCs *vide supra*). Thus the RQCs are a direct and most precise measure of the molecular alignment which, in chiral alignment media, is dependent on the absolute configuration. Still, this usage is limited as it requires the challenging determination of RQC signs to concisely draw conclusions on absolute configuration. Recently, some progress to overcome this remaining challenge was made by *Tzvetkova* and *Luy* through the development of a quadrupolar exclusive COSY (Q.E.COSY) experiment.^[89]

Besides these quantitative approaches Pöppler *et al.*^[38] showed that RQCs can be exploited in an innovative manner to distinguish between lithium aggregates. The authors persuasively reasoned that the values of RQCs ν_{rq} as defined in equation (1-13) become zero either if $\mathbf{A} = 0$, i.e. the aggregate is highly symmetric (O_h , T_d etc.) and no alignment within the gel is possible or if $\mathbf{V} = 0$, i.e. the EFG around ${}^7\text{Li}$ is symmetric (Figure 1-18).

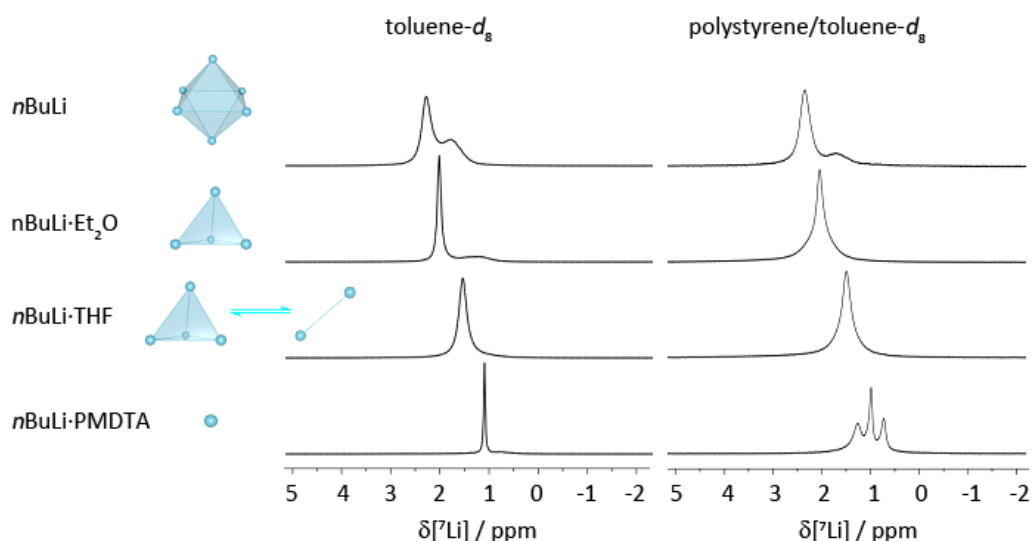
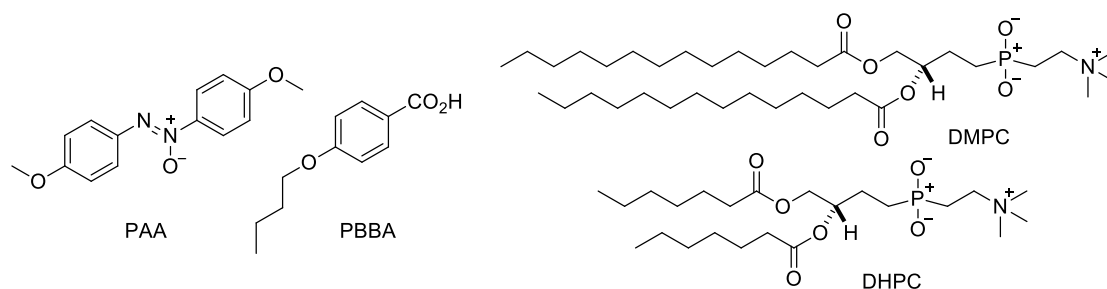


Figure 1-18: Aggregates of *n*BuLi and associated ${}^7\text{Li}$ spectra in isotropic (toluene- d_8) and anisotropic (PS/toluene- d_8) solution. Reprinted with permission from reference [38]. Copyright 2012 Wiley VCH.

Hence, this qualitative approach represents a smart and straightforward method to distinguish symmetric from asymmetric lithium aggregates. Yet, a more comprehensive determination of the EFG for molecules with $\mathbf{A} \neq 0$ will amend the possibilities of solid state NMR spectroscopy and facilitate structure determination as shown by Günther *et al.*^[87] This concept is explored in section 4.3.

1.3.3 Alignment Techniques

In 1963 Saupe *et al.* used solutions of *strong alignment media* like *p*-azoxyanisole (PAA) or *p*-butoxybenzoic acid (PBBA) which led to the first observations of anisotropic interactions (Scheme 1-3). However, these strongly orienting liquid crystals proved impractical for the early utilization of residual couplings since the RDCs and RQCs were too large to be determined with standard liquid state NMR techniques. Hence, the increase of research into anisotropic NMR spectroscopy was only initiated in the mid-90s when *weak alignment media* such as mixtures of dimyristoyl phosphatidylcholine (DMPC) and dihexanoyl phosphatidylcholine (DHPC) (Scheme 1-3) were first described.^[90]



Scheme 1-3: Typical examples for molecular components of nematic liquid crystals: *p*-azoxyanisole (PAA), *p*-butoxybenzoic acid (PBBA), dimyristoyl phosphatidylcholine (DMPC), dihexanoyl phosphatidylcholine (DHPC).

Molecules oriented in these media showed RDCs in the order of J -couplings which could be experimentally determined from known NMR spectroscopic methods (*vide supra*).^[90] Alternatively, in 2000 Tycko *et al.*^[70a] and Sass *et al.*^[91] independently revived an alignment concept already introduced by Deloche and Samulski in 1981.^[92] Various crosslinked polymers can take up huge amounts of solvents to give polymeric gels with massively increased volume. Swelling of small cylindrical sticks that neatly fit into NMR tubes leads to an extension that is restricted to the orientation along the tube (also see Figure 1-19 b). Accordingly, this leads to polymer strands that are predominantly oriented along the direction of the swelling and thus yield elongated cavities within the gel. Molecules dispensed within the gel (cavities) exhibit degrees of alignment that can be easily adjusted by parameters such as gel type, crosslinking percentage or solvent.^[93] Hence, so-called strain induced alignment in a gel (SAG) initiated several studies from which a multitude of possible polymeric alignment media resulted (Table 1-1).

Table 1-1: Selected polymeric alignment media, appropriate solvents and special remarks.

	Solvents	Special remarks
Polyacrylamide ^[70a, 91]	Strongly polar (e.g. H ₂ O)	Seminal work
Gelatin ^[94]	Polar (e.g. dimethyl sulfoxide, H ₂ O)	Chiral, can be made from gummibears
Polydimethylsiloxane ^[95]	Weakly polar (e.g. THF, <i>n</i> -hexane, benzene)	Single polymer resonance at 0.1 ppm
Polystyrene ^[93a]	Weakly polar (e.g. THF, toluene)	Chemically inert

Furthermore, applying external mechanical stress on polymer sticks is a convenient method to tune the alignment. Several stretching and compression devices have been developed that allow rapid preparation of samples. Note that uptake and sufficient equilibration of solvents in a gel may take several days to weeks. Application of outside forces reduces the time for sample preparation to a minimum. Figure 1-19 a) shows the components of a compression device as developed by Gayathri *et al.* in the group of Gil.

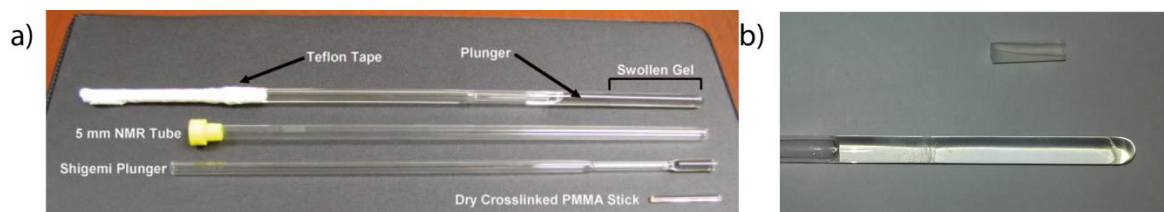


Figure 1-19: a) Compression device as applied by *Gayathari et al.* to mechanically achieve alignment within a gel. Reprint with permission from reference [37]. Copyright 2016 Wiley VCH. b) Swollen and non-swollen cylindrical PS.

The device is based on a *Shigemi* plunger that is pressed onto the polymer and thereby depletes the cavities within the gel, which in this case is a polymethyl methacrylate (PMMA) matrix. In Figure 1-19 b) non-swollen and confined, swollen cylindrical PS sticks can be visually compared.

Within the range of partial alignment media described above, suitable conditions can be found to study

- i) RDCs between adjacent nuclei with spin $I = 1/2$,
- ii) Residual interactions between the electric quadrupole moment of a nucleus with spin $I > 1/2$ and the surrounding EFG,
- iii) changes in the observed chemical shift.

While RDCs and RQCs have been described in more detail in the sections above the residual chemical shift anisotropy (RCSA) is not in the focus of this thesis. Excellent reviews are found elsewhere.^[96]

2 OUTLINE

Within the *Stalke* group the strong expertise in solid state structure determination of small metal-organic molecules^[97] has been increasingly complemented by liquid state NMR methods such as DOSY^[34, 36] or NMR spectroscopy in anisotropic media.^[38] Comprehensive structure elucidation allows the derivation of structure-reactivity principles and in turn enables smart design of chemical reactions. The goal of this thesis is to further extend and improve the *NMR toolbox*. In particular, it aims at the utilization of SSE to observe dynamical processes and the application of anisotropic NMR spectroscopy to obtain ample structural information.

Slice-Selective Excitation

Horizontal sample slices can be excited through soft frequency-shifted pulses in the presence of a magnetic field gradient. *Pöppler et al.* have employed this technique to study the diffusion of small molecules into a PS gel matrix.^[37] In chapter 3 the SSE experiment lays the foundation for a variety of applications in which heterogeneous, dynamic systems can be observed to

- i) characterize reaction intermediates,
- ii) perform single-shot fast titrations,
- iii) perform qualitative NMR chromatography.

The basic principles of these methods, as well as their implementation will be exemplified by tracing the reaction of *N,N,N',N'',N'''*-pentamethyldiethylenetriamine (PMDTA) with *n*BuLi, performing a titration of Li⁺ with 12-crown-4 and achieving chromatographic separation of four prototypical amino acids.

Anisotropic NMR Spectroscopy

Crosslinked polymer matrices induce partial alignment of molecules in solution, which leads to additional nuclear spin interactions, namely RDCs and RQCs. These observables are increasingly employed in the determination of the configuration and conformation (i.e. the 3D structure) of small molecules.^[66, 70-71, 88] In chapter 4 possibilities are shown on how to utilize RDCs and RQCs to

- i) observe the alignment of small molecules in polymers,
- ii) determine the 3D structure of molecules,
- iii) determine the aggregation of lithium organic compounds.

These applications are exemplified by the observation of small aromatic molecules oriented in PS and polybutylacrylate, the structure elucidation of synthetic β -amino acid esters containing an indane moiety and the determination of the aggregation of the lithium organic compound cyclopentadienyl lithium (CpLi).

3 APPLICATIONS OF SLICE-SELECTIVE NMR SPECTROSCOPY

3.1 Preparation of the SSE Experiments

Although SSE is a robust and straightforward method, it requires an accurate technical preparation:

- i) The determination of the magnetic field gradient strength.
- ii) The setting of the pulse duration.
- iii) The determination of the spatial detection sensitivity.

These steps have been adopted and advanced in the course of this work.

Determination of the magnetic field gradient strength

According to equation (1-1), for a given nucleus with gyromagnetic ratio γ the position of the slice z depends on the strength of the magnetic field gradient G_z and the frequency offset ν_z of the soft-pulse.

$$\nu_z = \frac{\gamma \cdot G_z \cdot z}{2\pi} \quad (1-1)$$

In the case of Bruker NMR probes the gradient strength G_z is adjusted by an electric current I_a that runs through the gradient coils.^a

$$G_z = g_{IG} \cdot I_a \quad (3-1)$$

The corresponding proportionality constant g_{IG} is a factory-determined value that often lacks reliability, in which case frequency offsets lead to inaccurate assumptions about the positions of slices. This prohibits further quantitative analysis such as the spatial distribution of molecules. Hence, the gradient strength must be evaluated beforehand.

One common evaluation method is to reproduce literature-known diffusion coefficients through DOSY measurements which are based on magnetic field gradients.^[30] If DOSY experiments are run with sufficient care for experimental conditions like temperature, viscosity and concentration, only inaccurate gradient strengths can lead to values that deviate from literature. In turn, this deviation allows the estimation of a reliable value for g_{IG} .^[98]

An alternative technique that requires less expensive care for the experimental conditions, depends on the imaging of cylindrical discs (¹H NMR-silent solids such as perfluorinated polymers) with well-defined spatial dimensions.^[99] The longitudinal image of these cylindrical discs in the frequency domain is proportional to its spatial dimensions and G_z is subsequently derived as the proportionality constant. A very similar approach was utilized in this work that further illustrates this technique:

A capillary filled with H₂O (similar to the experimental setup known from the *Evans* method for the determination of susceptibility^[100]) is immersed in an NMR tube filled with D₂O. The capillary is then

^a Technical specifications depend on the probe, standard broadband observe (BBO) and inverse (BBI) probe allow $I_{\max} = 10$ A.

shifted between two positions (Δz) at the top and the bottom of the active (*excitable*) volume of the NMR tube (Figure 3-1 a). At each step the exact position of the capillary is measured with a sliding calliper and subsequently a standard NMR imaging sequence (see 5 Experimental Part) consisting of a spin echo combined with two pulsed gradients (gradient strength at 20% $I_{a,max}$) is applied to obtain images as shown in Figure 3-1 b.

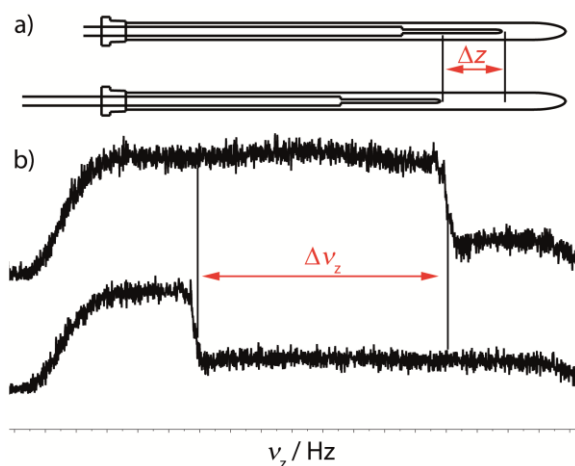


Figure 3-1: a) Two positions of the capillary within the NMR tube and their difference Δz . b) Corresponding images in the frequency domain with difference $\Delta \nu_z$.

The steep descent in the frequency domain corresponds to the bottom position of the capillary (more precisely, the definite H_2O volume) and quantitative, reproducible values can be read off at half of the maximum signal intensity. Experimental differences between positions Δz and frequencies $\Delta \nu_z$ are shown in Table 3-1.

Table 3-1: Differences between positions Δz and frequencies $\Delta \nu_z$ from three repetitions of sliding the capillary within the NMR tube as illustrated in Figure 3-1.

	Experiment 1	Experiment 2	Experiment 3
Δz /mm	8.8030	8.5770	8.8830
$\Delta \nu_z$ /Hz	38437	37321	38596

These values allow straightforward calculation of the real gradient strength G_z :

$$G_z = \frac{2\pi \cdot \Delta \nu_z}{\gamma \cdot \Delta z_z} \quad (3-2)$$

With a known $I_{a,max} = 10$ A this resulted in a spectrometer specific constant g_{IG} of 5.1 ± 0.1 $Gcm^{-1}A^{-1}$, which considerably deviates from the value (5.6 $Gcm^{-1}A^{-1}$) initially stated in the hardware specification of the instrument. Note that for the newly determined g_{IG} and according to equation (1-1), at $I = 0.5I_{max}$ a frequency offset ν_z of 98209 Hz excites a slice 9.0 mm above the center of the gradient coil. By applying the factory determined g_{IG} this slice would be assumed to be 8.2 mm above the centre of the gradient coil. This 9% difference clearly highlights the imperative to (re)determine the gradient strength for individual probes if they are used in slice-selective applications.

Setting of the pulse duration and power

Excitation of an inhomogeneous sample that exhibits a continuous variation of properties along its spatial dimension (e.g. a shifting resonance that results from a concentration gradient along the NMR tube) leads to broadened signals. A less selective (harder) pulse excites a thicker slice which gives increased sensitivity but a further broadening due to the summation of the signals originating from different parts of the sample (Figure 3-2).

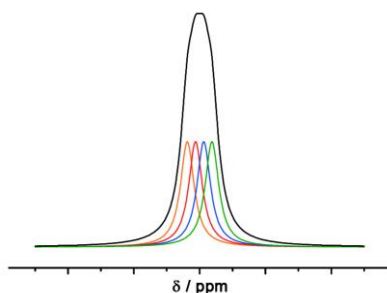


Figure 3-2: Schematic representation of signal widths in inhomogeneous samples. A changing chemical shift can be resolved by four selective excitations (colored) in the presence of a gradient. An unselective or less selective excitation leads to a *summed*, broadened signal (black).

As defined in equation (1-2), the signal thickness depends on the excitation bandwidth $\Delta\nu_{\text{BW}}$ which in turn is governed in a non-trivial manner by the pulse duration and the shape of the pulse. The Bruker spectrometer software TopSpin includes a decent simulation tool (shape tool, *stdisp*) in which *Fourier* transformations of specified pulses yield their excitation bandwidths. In this work, exclusively *Gaussian Cascade (G4)* soft-pulses were applied.^[101] This shape achieves a nearly rectangular excitation profile across the excitation bandwidth of the pulse (Figure 3-3).

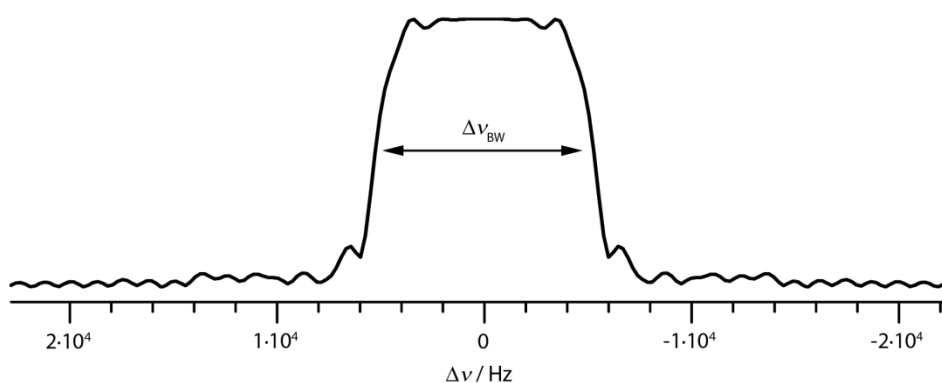


Figure 3-3: Excitation profile of an 800 μs *G4 Gaussian Cascade* pulse with 90° flip angle. The profile corresponds to a frequency bandwidth $\Delta\nu_{\text{BW}}$ of ~ 11 kHz at half maximum intensity.

Furthermore, slices were chosen to be excited succinctly, i.e. they should be separated by 1 mm and exhibit an approximate thickness of 1 mm as well to avoid overlap of slices. The necessary parameters to obtain this mode are detailed in Table 3-2.

Table 3-2: Pulse shape and pulse duration for nuclei relevant to this work that yield excitation bandwidths $\Delta\nu_{\text{BW}}$ which at a gradient strength G_z excite the desired slice thickness.

Nucleus	Pulse shape	Gradient strength $G_z/\text{G}\cdot\text{cm}^{-1}$	Pulse duration /ms	Exc. bandwidth $\Delta\nu_{\text{BW}}/\text{Hz}$	Slice width /mm
^1H	<i>G4</i>	25.6	800	~11000	~1
^2H	<i>G4</i>	51.1	4800	~1700	~0.5
^7Li	<i>G4</i>	51.1	1000	~8500	~1

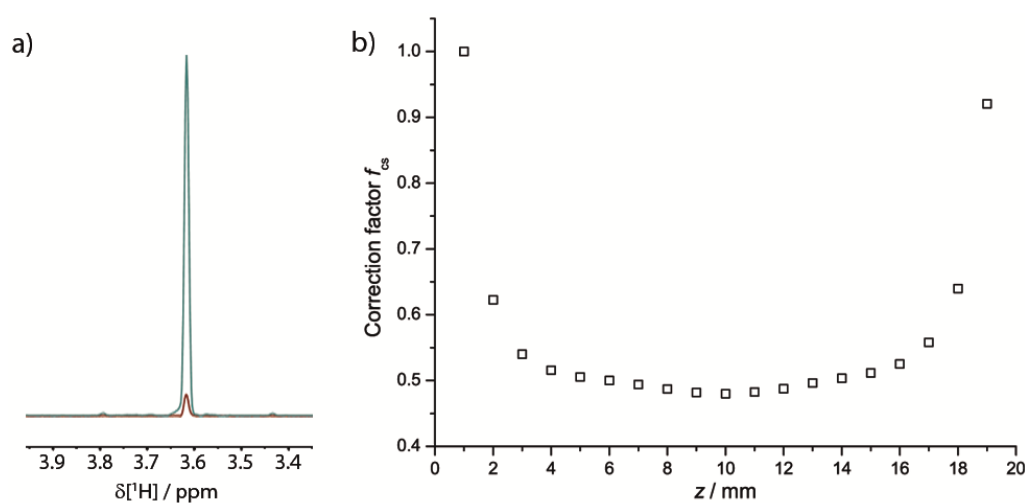
For standard, rectangular high power pulses the flip angle is set through the pulse duration. This duration cannot be varied for selective pulses without changing the excitation bandwidth. Hence, maximum intensity flip angles (90°) for soft pulses must be adjusted by setting the pulse power (amplitude of the radio frequency pulse). TopSpin enables the determination of the appropriate power for a desired flip angle at a certain pulse length by integrating the adiabatic shape of the soft pulse. The full set of parameters to obtain 90° flip angles and 1 mm slice width is given in Table 3-3.

Table 3-3: Pulse shape, pulse duration and pulse power for relevant nuclei that achieve the desired 90° flip angle.

Nucleus	Pulse shape	Pulse power /W	Pulse duration /ms	Flip angle / $^\circ$
^1H	<i>G4</i>	0.358	800	90
^2H	<i>G4</i>	3.93	4800	90
^7Li	<i>G4</i>	3.18	1000	90

Determination of the spatial detection sensitivity

The active region of the NMR probe (given by the coil dimensions) is limited to a length of around 2 cm. This allows excitation of ~20 slices à 1 mm. Signal integrals (I_{sig}) of 90° soft pulses are reduced to ~5% compared to 90° hard pulses (Figure 3-4 a). However, the sensitivity profile of the detection coil is not perfectly uniform, which has to be addressed through correction factors f_{cs} ($I_{\text{corS}} = f_{\text{cs}} \cdot I_{\text{sig}}$). Therefore, homogeneous reference samples were measured and signal integrals were determined (Figure 3-4 b).

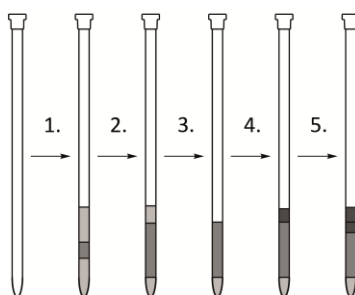
Figure 3-4: a) Comparison of a 90° hard pulse (green) and a 1 mm *G4* 90° soft pulse (red). b) Correction factors f_{cs} for the integrals of 19 slices (top to bottom) for 1 mm *G4* 90° soft pulses.

Note that in Figure 3-4 a) the intensity decreases but the full width at half maximum of the signal decreases as well. This increase of spectral resolution results from less magnetic imperfections along thinner slices in homogeneous samples. According to Figure 3-4 b) the detection coil exhibits mitigations especially at the outer regions. These mainly originate from the confined spatial dimensions of the coil and cannot be circumvented.

3.2 Observation of Reactions

Chemical reactions can be observed by monitoring substrates in a series of individual NMR spectra over time. The observation of fast reactions and short-lived intermediates has become possible through custom build NMR hardware and techniques such as stopped flow and RI NMR.^[29] In this thesis, a very different approach was made by utilizing the spatial resolution that can be achieved through SSE:

As described above, SSE has already been utilized to study diffusion of solvents and solutes into polymers. In particular, *Pöppler et al.* focused on inert crosslinked PS that can be swollen with toluene solutions containing highly reactive lithium amides.^[37] The fully equilibrated polymers (swollen over a sufficiently long period of time, usually several days) containing a reactive species is the basis for the innovative SSE reaction monitoring approach. Here, a further substrate is added on top of the gel and diffuses into the polymer to yield a reaction front between both compounds. This progress (Scheme 3-1) can be readily observed through SSE.

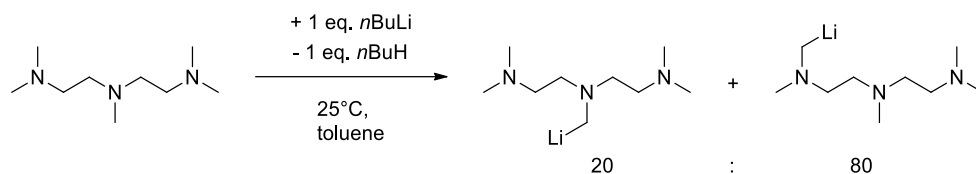


Scheme 3-1: Basic steps of sample preparation: 1. A crosslinked polymer (dark grey) is immersed in a solution which contains the reactive (lithium organic) species (bright grey). 2. The polymer is swollen (equilibrated) over a sufficiently long time period. This can be evaluated through homogeneous signal intensities via SSE as shown by *Pöppler et al.* 3. The supernatant solution above the polymer is removed. 4. A solution containing a further substrate (black) is added on top of the polymer. 5. The substrate diffuses into the gel. 6. This is then followed by SSE over an appropriate period of time.

Reaction between *n*BuLi and PMDTA

The lithiation of PMDTA with *n*BuLi was examined as a case study. PMDTA is usually applied as a donor base that allows disaggregation of lithium organic compounds into smaller aggregates.^[15a, 19, 102] The disaggregation and subsequent lithiation of substrates is usually performed at lowered temperatures to control the reactivity and delicate selectivity of the lithium organic compounds. Disaggregation at room temperature frequently leads to the undesired lithiation of the donor base. A most prominent

example is *n*BuLi in THF which easily undergoes lithiation and ether cleavage.^[103] A similar behaviour is documented for *n*BuLi and PMDTA (Scheme 3-2).^[104]



Scheme 3-2: The reaction between PMDTA and *n*BuLi in toluene at 25°C yields a mixture of lithiated PMDTA (20% central methyl group, 80% terminal methyl groups lithiated).

Previous studies have shown that at room temperature and upon addition of one equivalent *n*BuLi a mixture of products results in which either the central methyl group (20%) or one of the terminal methyl groups (80%) of PMDTA is lithiated.^[105] This reaction was chosen as a case study to test the reaction observation approach:

According to Scheme 3-1, samples were prepared by immersing a cylindrical PS stick in a solution of *n*BuLi in toluene-*d*₈. Then PMDTA was added on top of the swollen polymer. Slice-selective ¹H and ⁷Li NMR measurements started ~3 h after preparation of the sample and went on for three days (experimental details on sample preparation and SSE parameters are described in the experimental section 5.1). Figure 3-5 shows a series of 19 slice-selective ⁷Li NMR spectra as obtained three days after the addition of PMDTA to the polymer imbibed with *n*BuLi.

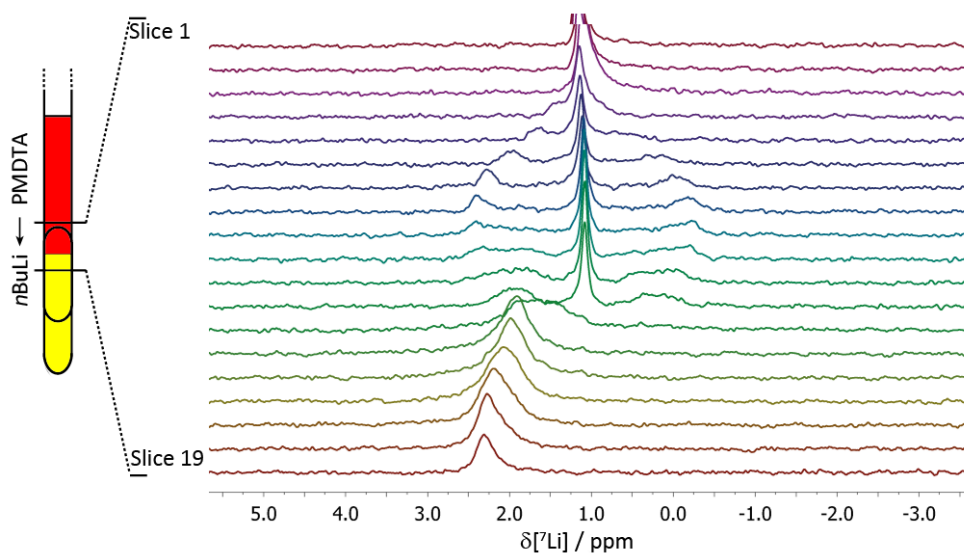


Figure 3-5: Slice-selective ⁷Li NMR spectra (slices 1 to 19) of the reaction between *n*BuLi and PMDTA within a PS gel three days after the addition of PMDTA.

Slices 19 and 18 at the bottom of the active volume still show the signal of unreacted *n*BuLi at 2.3 ppm.^[16, 106] In the upper slices 1-12 a narrow resonance at 1.1 ppm is observed that can be assigned to the products of the lithiation (Scheme 3-1). It is impossible to discriminate the terminally and internally lithiated products since their chemical environments are almost identical and the narrow chemical shift range of ⁷Li nuclei prohibit these minute distinctions.

Slices 1–4 are located outside the polymer and, hence, only a singlet is observed. Satellite peaks in slices 5–12 arise from RQCs due to the partial alignment of the product molecules inside the stretched gel (*vide infra*). For *n*BuLi in slices 18 and 19 a singlet is observed since the O_h symmetry of the hexameric $[n\text{BuLi}]_6$ aggregate in toluene does not allow any alignment.^[38]

Note that the size of the RQCs in slices 5–12 changes significantly. Various factors may contribute to this phenomenon: Most likely, the amount of strain within the gel is changing.^[37,53] In particular, the strain may decrease towards the surface between neat solvent and polymer. Accordingly, the alignment of molecules in these slices becomes weaker as can be observed for spectra of slices between 5 and 8. This is strongly supported by ^2H SSE spectra and the observed splitting for their methyl deuterons (Figure 3-6).

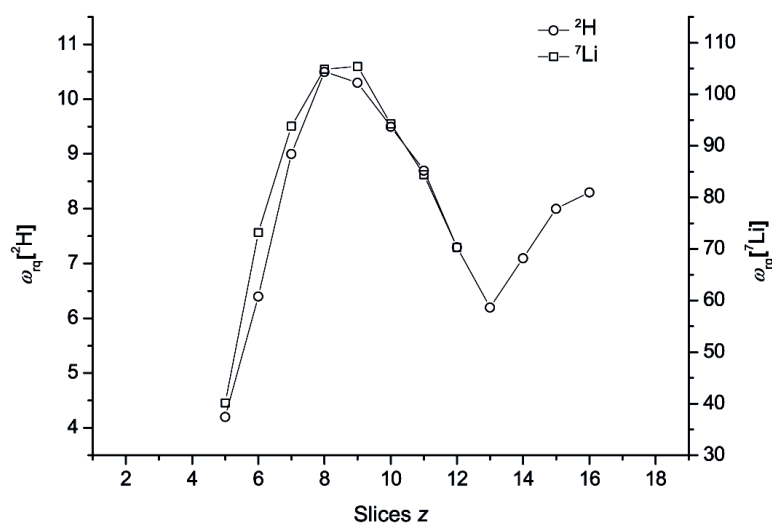


Figure 3-6: Correlation between slices and RQCs for ^2H of the methyl group in toluene- d_8 and ^7Li of the lithiation product.

Toluene does not take part in the reaction and can therefore serve as a probe for the alignment strength of the polymer. Figure 3-6 shows that the lithium specimen and toluene- d_8 experience proportional strain. Hence, it is reasonable that the change in ^7Li RQCs is due to changes in the alignment matrix. These changes presumably results from the exothermic reaction between *n*BuLi and PMDTA which disturbs the relaxed/stretched polymer network and results in the dip at the reaction front (slice 13).

Slices 9–17 in Figure 3-5 exhibit further interesting features: While the central transition is sharp in slices 9–12, the satellite transitions become increasingly broadened. This indicates a superposition of different RQCs which again mainly originate from varying alignment as supported by the ^2H SSE spectra of toluene- d_8 (though different, superimposed reaction products with similar ^7Li shift may exist). This non-equilibrium strain in the middle parts of the gel is most likely due to a response of the PS matrix upon the uptake of PMDTA.^a

^a Note that a plethora of possible parameters within hydrogels influences their swelling behaviour. These range from solute characteristics like different osmotic pressures to external stimuli such as the temperature. These responses are particularly relevant to the application of hydrogels in medicinal chemistry.^[107]

Additionally, RQCs are (in contrast to ^7Li chemical shifts) very sensitive towards changes of the EFG which is analogous to changing aggregation motifs. Hence, slices 9-12 may also be characterized by dynamical processes involving different aggregates as well as both lithiated products.

Slices 13-17 are characterized through the vanishing of satellites and increasingly broad central transitions which exhibits a changing chemical shift between 1.6 and 2.3 ppm. This impressively marks the reaction front between PMDTA and $n\text{BuLi}$. It must be highlighted, that the reaction between $n\text{BuLi}$ and PMDTA is strongly exothermic.^[104] Hence, samples without a polymer gel exhibit convection and uncontrolled mixing of both compounds. Within the PS matrix $n\text{BuLi}$ is practically immobilized and a remarkably distinct reaction front results from the predominant diffusion of the PMDTA molecules. However, the lack of resolved RQCs for ^7Li in these slices may stem from a decreased degree of alignment through increased temperature. This effect cannot be fully separated from dynamic processes that constitute the lithiation and which are in the focus of interest. Therefore, further conclusions must be drawn from the series of 19 slice-selective ^1H NMR spectra of the same sample (Figure 3-7 a).

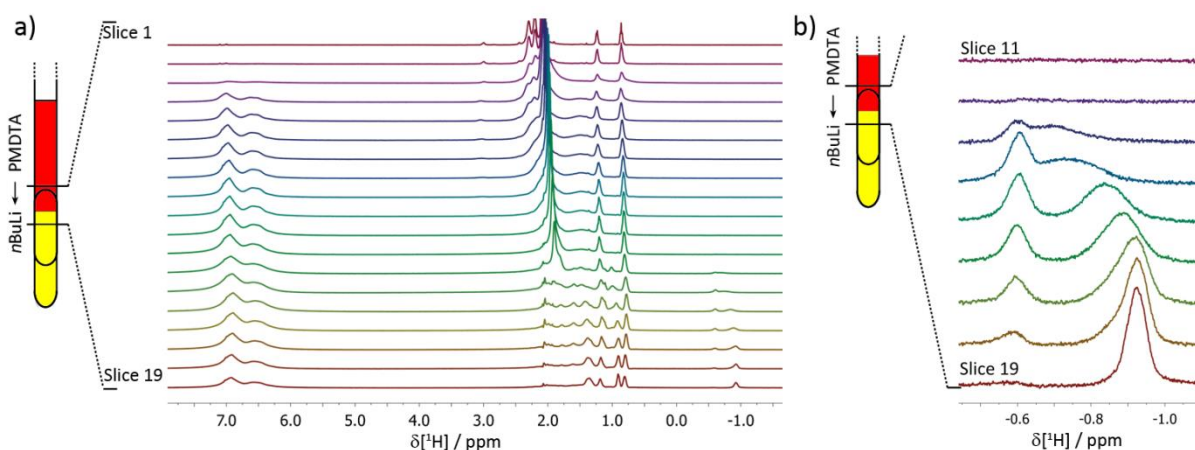


Figure 3-7: a) Slice-selective ^1H NMR spectra (slices 1–19) of the reaction between $n\text{BuLi}$ and PMDTA within the PS gel, three days after addition of PMDTA. b) Enlarged excerpt of the same spectrum for slices 11–19 at the high field limit of the spectrum.

In contrast to ^7Li , ^1H gives relatively crowded spectra with broad polymer resonances between 6.1 and 7.2 ppm and around 2 ppm. Especially between 0.6 and 2.3 ppm signals from various reaction components and the polymer disable attribution of single resonances. However, the isolated region between -0.5 and -1.1 ppm is quite informative (Figure 3-7 b). This part includes the exposed $n\text{BuLi}$ $\alpha\text{-CH}_2$ protons (-0.92 ppm)^[16, 106] and a number of interesting features: The most remarkable characteristic is the appearance of an additional, non-shifting resonance at -0.59 ppm in slices 18-13 with a maximum intensity in slice 15. The connection between the reaction progress and the appearance of this peak can be easily illustrated through a comparison of its integral and the integral of the remaining $n\text{BuLi}$ $\alpha\text{-CH}_2$ signal (although this signal is shifting, it can be fully attributed to $n\text{BuLi}$, *vide infra*) over the course of three days (Figure 3-8).

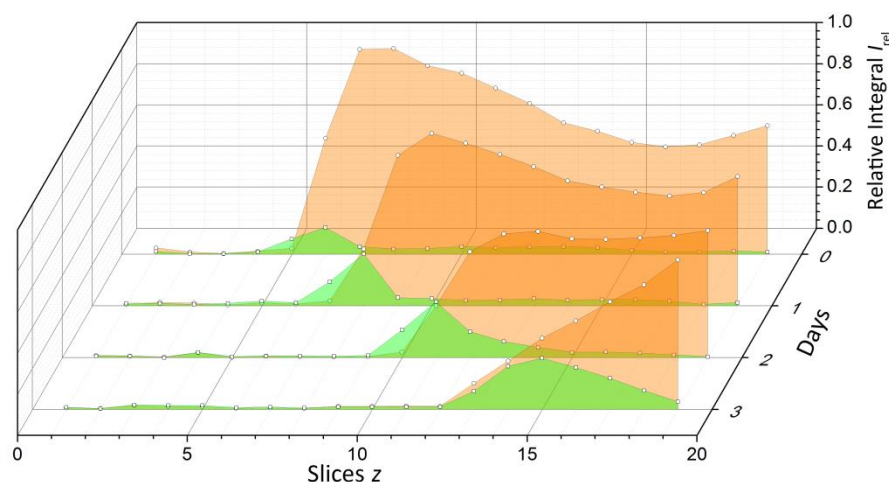


Figure 3-8: Relative integrals for the ^1H NMR signals of the intermediary species at -0.59 ppm (green) and the signal of the $n\text{BuLi}$ $\alpha\text{-CH}_2$ protons (orange) as function of slice number, 3 h, 1, 2 and 3 days after addition of PMDTA to $n\text{BuLi}$ inside the pre-swollen PS gel. All signals high-field to -0.59 ppm were attributed to $n\text{BuLi}$.

Figure 3-8 clarifies the motion of the reaction front, where the $n\text{BuLi}$ $\alpha\text{-CH}_2$ signal (orange) vanishes and the signal at -0.59 ppm appears. The latter builds up and decays exclusively at the reaction front which impressively indicates its intermediate character. In fact, this signal as well as further signals in the crowded part of the spectra (Figure 3-9 a) can be assigned to the complex $[(n\text{BuLi})_2\text{PMDTA}]_2$ which was first described by Gessner and Strohmann in 2007.^[104]

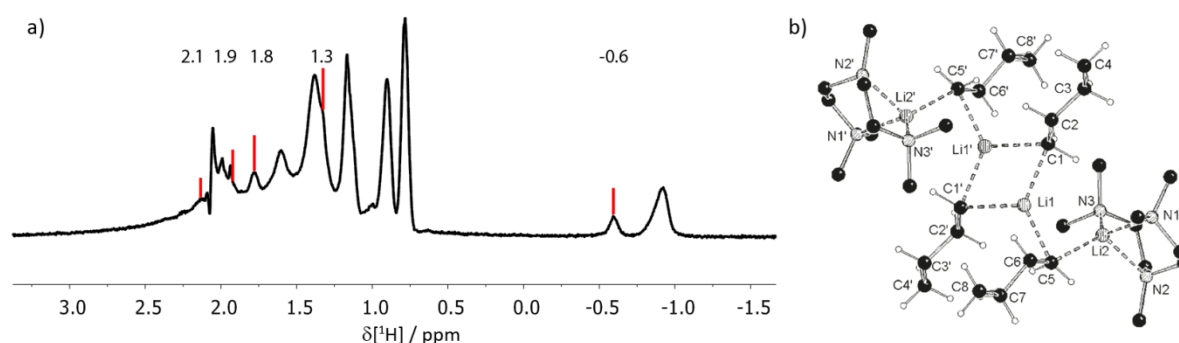


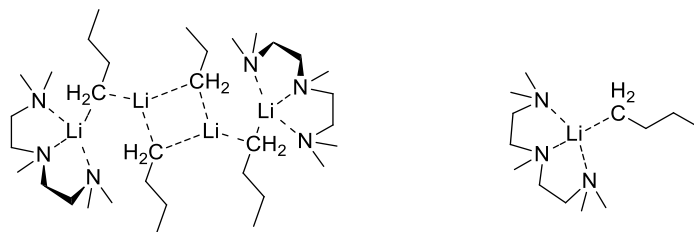
Figure 3-9: a) NMR spectrum of slice 17 after 3 d with explicit values that can be attributed to $[(n\text{BuLi})_2\text{PMDTA}]_2$ according to the report from Gessner and Strohmann.^[104] b) Crystal structure of $[(n\text{BuLi})_2\text{PMDTA}]_2$. Crystal structure reprinted with permission from reference [104]. Copyright 2007 Wiley VCH.

In their study, the authors found that remarkably stable crystals can be obtained from quenched reactions between $n\text{BuLi}$ and PMDTA. The structure of the crystallized molecules was determined via XRD (Figure 3-9 b) and single crystals were subsequently solved in toluene- d_8 to record ^1H NMR spectra. Note that the crystallization as well as XRD and NMR measurements of the stable intermediate were performed at 223 K while the authors observed that the lithiation inevitably starts above 253 K. Additional comprehensive quantum chemical calculations on $[(n\text{BuLi})_2\text{PMDTA}]_2$ and a hypothetical monomeric $[(n\text{BuLi})\text{PMDTA}]$ species together with corresponding transition states between both aggregates and the lithiated product unambiguously indicated that $[(n\text{BuLi})_2\text{PMDTA}]_2$ is the crucial intermediate in the lithiation of PMDTA.^[104]

Figure 3-8 displays several further features: While 3 h after the addition of PMDTA to the pre-swollen gel, the reaction front is relatively sharp with little formation of $[(n\text{BuLi})_2\text{PMDTA}]_2$, it becomes significantly blurred as PMDTA moves downward with time. This concentration gradient is in accordance with what would be expected from diffusion into/within a gel.^[108] Furthermore, the dip in the $n\text{BuLi}$ concentration which is initially observed in the center of the gel most likely arises from slow and incomplete diffusion of $n\text{BuLi}$ into the polymer during sample preparation.^a This again underlines that the $n\text{BuLi}$ hexamer ($M_W = 383 \text{ g}\cdot\text{mol}^{-1}$) may be regarded as rather static compared to the faster diffusing PMDTA molecules ($M_W = 173 \text{ g}\cdot\text{mol}^{-1}$).

The concluding remarks on reaction observation via SSE must address another interesting feature of Figure 3-7 b) that has not been fully acknowledged yet: the $n\text{BuLi}$ resonance at -0.92 ppm is shifting downfield and significantly broadens from slice 19 to 13. It is obvious that, due to the diffusion of PMDTA and the resulting concentration gradient, different ratios between the donor base and $n\text{BuLi}$ must occur in the slices along the NMR tube. Hence, the change in shift resembles the course of a titration experiment in which two compounds ($n\text{BuLi}$ and PMDTA) form another $[(n\text{BuLi})_x(\text{PMDTA})_y]$ complex with unknown stoichiometry. Within this work, only a hypothesis can be made about the aggregate:

The shifting signal seemingly approaches the resonance of the dimeric intermediate at -0.59 ppm which loosely indicates similar chemical environments for the unknown $[(n\text{BuLi})_x(\text{PMDTA})_y]$ and $[(n\text{BuLi})_2\text{PMDTA}]_2$. A gradual breaking down of $[n\text{BuLi}]_6$ towards the monomeric $[(n\text{BuLi})\text{PMDTA}]$ adduct has been proposed several times, but not yet been experimentally confirmed.^[15a] The experimental findings shown above indicate this disaggregation process.



Scheme 3-3: Proven intermediate $[(n\text{BuLi})_2\text{PMDTA}]_2$ (left) and potential monomeric $[(n\text{BuLi})\text{PMDTA}]$ adduct (right).

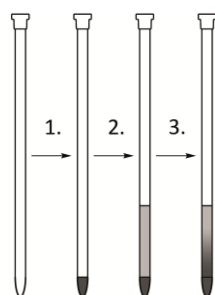
As already stated before, the latter observations are made through a setup that resembles an elegant titration experiment. The diffusion of PMDTA leads to different ratios of PMDTA and $n\text{BuLi}$ along the NMR tube. This was correlated with changing chemical shifts of nuclei in the corresponding slices. This approach is optimized and exploited in the application of SSE as a single-shot titration experiment.

^a As already mentioned, the swelling behaviour of the polymer is not in the focus of this thesis. However, preliminary studies showed that the species and the amount of solute heavily influence these dynamics. The swelling of polymers with neat solvent leads to extreme expansion of the gel and frequent *polymer ripping*. Different osmotic pressures as well as different diffusion characteristics are evident explanations for this interesting phenomenon.

3.3 Single-Shot Titration

Chemical shift titration is a powerful method to determine the stoichiometry and stability constants of complexes in coordination, supramolecular and medicinal chemistry.^[109] In such an experiment, the chemical shift of a particular resonance attributed to one component (e.g. the metal, guest or target) is monitored in a series of NMR spectra, while the concentration of the other component (e.g. the host or ligand) is systematically varied. Rather than incrementing the concentration step-by-step as in the traditional approach, the experiment on *n*BuLi and PMDTA (section 3.2) showed the possibility to perform titrations in a single experiment.

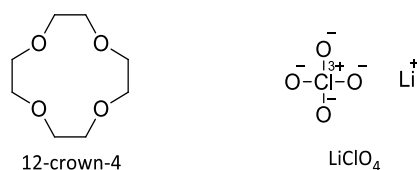
The most obvious drawbacks of a complete adoption of the approach described above are the time-consuming sample preparation (several days of swelling) and the challenging measurements within the PS gel: gel samples require extensive manual shimming *on the* free induction decay (FID) to achieve sufficiently homogeneous signals and resolution (also see experimental section 5.1). These pitfalls are accepted for SSE in reaction observation since the gel is essential to prevent the convection (only this allows the persistent formation of a reaction front) and allows the detection of anisotropic parameters (which add structural information). However, this procedure would be less applicable, simplistic and straightforward than standard titration experiments on several different samples with changing concentrations. The decisive step to achieve a smart single-shot titration experiment is to dismiss the gel and instead layer two plain compounds in solutions (Scheme 3-4). In the present study, this was achieved by solidifying the bottom compound through a decrease in temperature.



Scheme 3-4: Basic steps of sample preparation: 1. The first compound (neat) is filled into the NMR tube. 2. The second compound, immersed in deuterated solvent, is carefully added on top of the first compound. 3. The compound at the bottom then diffuses into the solution above. 4. This can be readily observed through SSE.

Titration of Li⁺ with 12-crown-4

As a case study the titration of Li⁺ with 12-crown-4 was performed.^[110] LiClO₄ was used as a Li⁺ source since it is highly soluble in organic solvents such as acetonitrile.



Scheme 3-5: Compounds for the titration experiment.

Crown ethers are versatile chemical compounds that are especially known for their selective binding towards metal cations.^[111] This enabled their application as crucial detection moieties in molecular sensors^[112] and as transfer agent in phase transfer catalysis^[113]. The ether 12-crown-4 has a particularly high affinity towards Li^+ and the complexation is well characterized.^[110] Hence, this complexation was chosen as a case study to test the single-shot titration approach:

According to Scheme 3-4 samples were prepared by first filling liquid 12-crown-4 in an NMR tube. The sample was then cooled down to 5°C to solidify the ether (melting point 16°C). Subsequently, a solution of LiClO_4 in acetonitrile- d_3 was layered on top of the solid ether. This procedure prevents mixing of both components (without a gel) prior to measurements, and only inside the NMR magnet (25°C) the ether melts and slowly diffuses into the LiClO_4 solution. Diffusion of LiClO_4 into the ether reservoir occurs likewise, but with little impact on the ^7Li concentration due to the $\sim 9:1$ volume ratio and the lowered concentration gradient (5.1 mol/L LiClO_4 solution as opposed to bulk ether). Approximately 3, 6, and 9 h after sample preparation, slice-selective ^1H and ^7Li NMR measurements were performed (experimental details on sample preparation and SSE parameters are described in the experimental section 5.1). Figure 3-10 shows the series of 19 slice-selective ^1H NMR spectra recorded after 6 h and the corresponding normalized integrals of the ^1H signal in each slice.

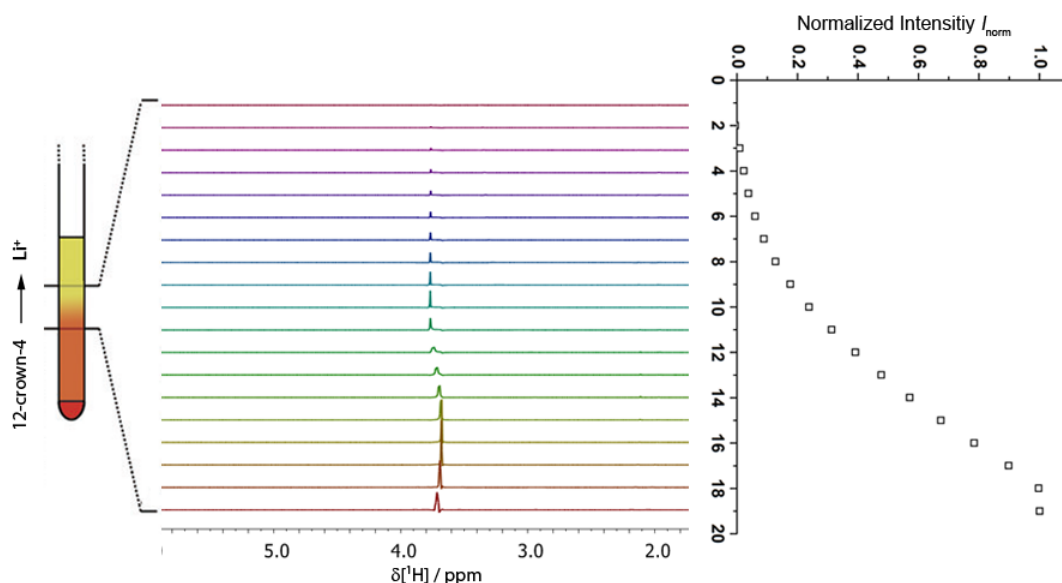


Figure 3-10: Slice-selective ^1H NMR spectra of 12-crown-4, 6 h after sample preparation and the corresponding concentration profile from corrected and normalized Intensities.

The single chemical shift of the ether signal exhibits only minor changes and is therefore not suitable for any further analysis. However, signal intensities decrease from bottom to top indicating the diffusion of 12-crown-4 in the same direction. Integration of the signals for all 19 slices and subsequent correction according to the sensitivity of the detection coil (see section 3.1) illustrates a smooth concentration gradient. In contrast to slice-selective ^1H measurements, ^7Li spectra exhibit a more pronounced change of the chemical shift for the Li^+ ion (Figure 3-11).

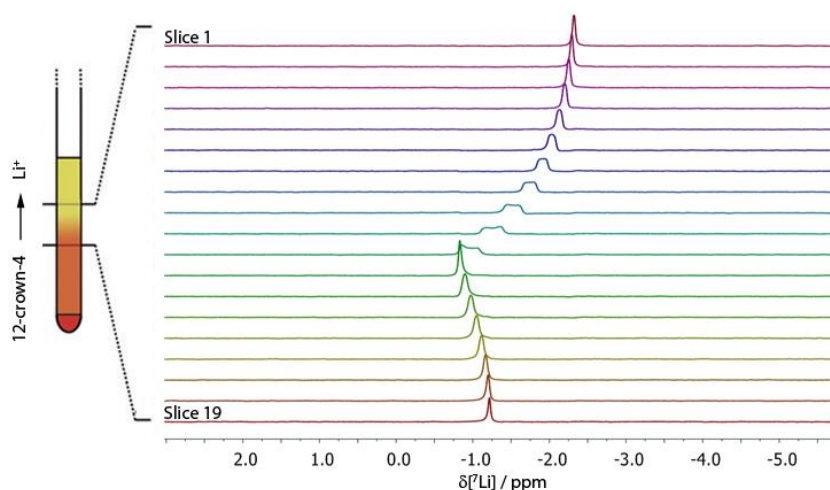


Figure 3-11: Slice-selective ${}^7\text{Li}$ NMR spectra of LiClO_4 in acetonitrile- d_3 in the presence of a gradient of 12-crown-4, 6 h after sample preparation.

The signals for ${}^7\text{Li}$ were integrated and corrected. These values are approximately constant and thereby confirm that almost no diffusion of LiClO_4 occurs. Slice 1 at the top of the active volume shows a narrow ${}^7\text{Li}$ resonance at -2.3 ppm which is typical for Li^+ coordinated by four acetonitrile molecules.^[110] As the ether concentration increases from slice 1 to 19 (Figure 3-10), the ${}^7\text{Li}$ resonance is shifted downfield and reaches a maximum (-0.8 ppm) in slice 12.^[110] Note that the ${}^7\text{Li}$ resonances in slices 6 to 11 are increasingly broadened. This is due to the strongly changing ether concentration within these slices which can only be monitored as superimposed (rectangular weighted sums) signals. This broadening could be overcome by further decreasing the slice width (see section 3.1) at the cost of sensitivity. After a maximum downfield shift is observed in slice 12, the signal moves upfield again until it reaches -1.2 ppm in slice 19.^[110]

Quantitative analysis^[114] of titration data depends on well-determined concentrations of all components involved. Calibration curves that correlate absolute signal intensity I_{abs} and concentration are a common measure to obtain these values. Instrument specific calibration curves for ${}^1\text{H}$ and ${}^7\text{Li}$ are documented in the experimental section 5.3 and were applied to convert the measured and corrected absolute integrals for each slice-selective ${}^1\text{H}$ NMR and ${}^7\text{Li}$ NMR spectrum into concentrations. Finally the specific ether/lithium ratio of a slice could be assigned to each corresponding ${}^7\text{Li}$ chemical shift (Figure 3-12).

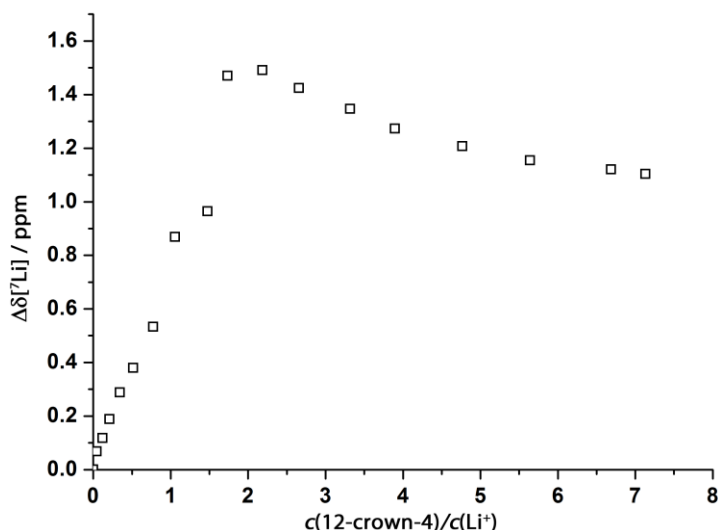


Figure 3-12: Chemical shift difference $\Delta\delta[{}^7\text{Li}]$ between pure Li^+ in acetonitrile and varying amounts of 12-crown-4 as measured 6 h after preparation of the sample.

Figure 3-12 further illustrates the changing ${}^7\text{Li}$ resonance in the presence of different ether/lithium ratios which qualitatively matches the titration curve of Li^+ by 12-crown-4 as provided by *Masiker* (Figure 3-13).^[110b]

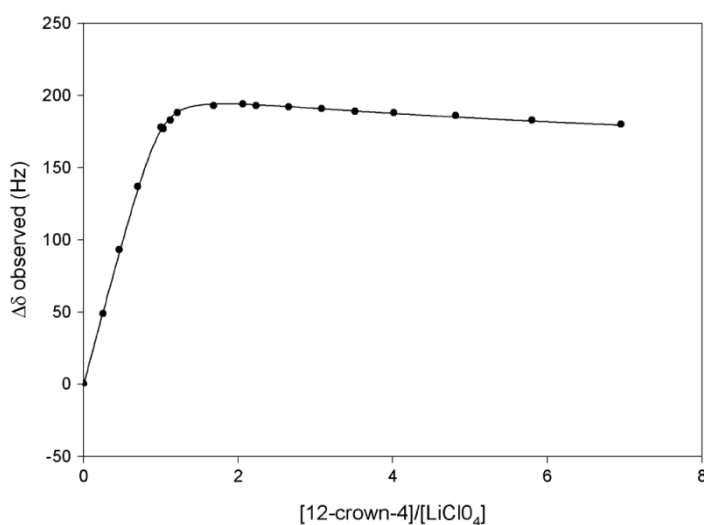


Figure 3-13: Chemical shift difference $\Delta\delta[{}^7\text{Li}]$ between pure Li^+ in acetonitrile and varying amounts of 12-crown-4. Reprint with permission from reference [110b]. Copyright 2010 John Wiley and Sons.

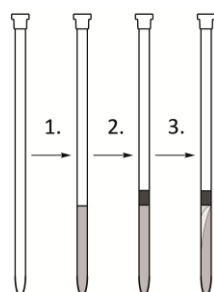
The authors verified a 1:2 binding model between Li^+ and 12-crown-4 in which lithium (M) is stepwise coordinated by the ether (L) resulting in a sandwich structure (ML_2) at higher ether/lithium ratios.^[110] For the SSE attempt decreasing the initial concentrations is the first choice to obtain better resolved results since the concentration gradient would be more shallow which increases the number of sampled data points for low ratios and, additionally, reduces peak widths. Finally, reducing slice widths by decreasing the excitation bandwidth would allow the fine tuning of peak resolution and once again sampling of more data points.

Meanwhile, the idea of utilizing concentration gradients and SSE for titration experiments was picked up and further improved by *Mitrev et al.* in the group of *Jeannerat*.^[115] In their study, the authors based

experiments on the use of agar gel samples in which one compound was dispersed. The second compound is layered on top of the gel and subsequently effects the concentration gradient through diffusion into the gel. This approach is very similar to the *n*BuLi PMDTA reaction observation in PS as described above. However, in the case of this titration experiments, agar has several advantages in comparison to a crosslinked polymeric gel such as PS: the authors state that the agar “gels have a remarkable property of causing no observable changes in the chemical shifts and lineshapes compared to normal aqueous solutions”. The gel matrix is not visible in the NMR spectrum due to its “small quantity and favourable relaxation properties” and gel samples can be very fast and easily prepared by mixing the stationary compound, water and granular agar and applying “a simple heating/cooling cycle”.^[115] Agar is commonly used in gel-electrophoresis^[116] for the separation of compounds in mixtures. This led to a case study in which SSE is exploited as a tool to perform qualitative chromatography.

3.4 Qualitative Chromatography

Characterization of complex mixtures is increasingly important for analytical chemists.^[46] Liquid chromatography and NMR can be directly coupled, e.g. via flow probes, enabling on-the-fly separation and detection of mixture compounds.^[117] Additionally, DOSY allows qualitative separation of mixture compounds according to self/tracer diffusion coefficients.^[118] This approach can be further enhanced by the application of diffusion affecting surfactants, complexing agents and likewise which is known as *matrix-assisted DOSY*.^[119] However, it must be emphasized that tracer diffusion (most commonly described through the *Stokes-Einstein* equation^[32]) is different from chemical diffusion along a concentration gradient as described through *Fick's laws*^[120] (both become only identical at infinite dilution^[121]). The applications shown above illustrate that chemical diffusion can be readily observed through SSE. Hence, SSE can be utilized as part of a complementary and expedient method to achieve a chromatographic separation. The key to achieve this separation is the preparation of a stationary phase which grants homogeneous starting conditions (all compounds start diffusing at the same time at the same boundary) and sufficient separation of compounds as they diffuse into the gel (Scheme 3-6).



Scheme 3-6: Basic steps of sample preparation: 1. The stationary phase is prepared within the NMR tube. 2. A mixture of compounds to separate is added on top of the stationary phase. 3. The compounds start diffusing into the stationary phase at different rates. 4. The diffusion progress can be readily observed through SSE.

Separation of amino acids L-aspartic acid, L-histidine, L-tyrosine and L-valine in agarose

As a case study the separation of amino acids L-aspartic acid (Asp, acidic), L-histidine (His, basic), L-tyrosine (Tyr, polar/neutral) and L-valine (Val, apolar) within agar hydrogel was examined (Figure 3-14).

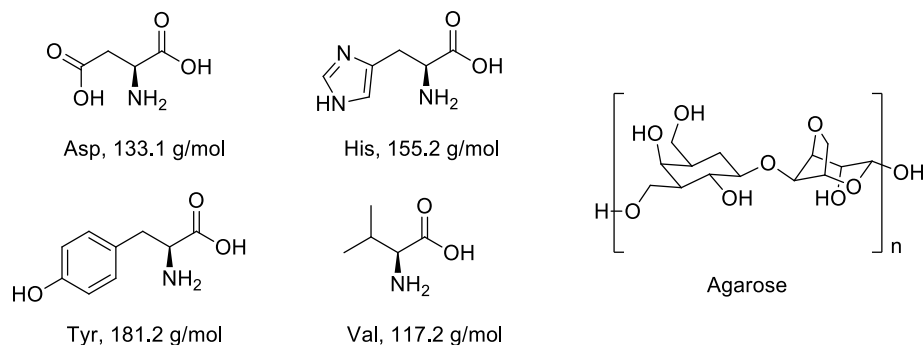


Figure 3-14: Amino acids that were used in the qualitative chromatography through SSE case-study, from left top to right bottom: L-aspartic acid, L-His, L-Tyr and L-Val. Additionally the structure of agarose as main component of agar is shown on the right.

The amino acids are similar in molecular weight (117.2 to 181.2 g/mol) and soluble in slightly basic water. As already stated above agar gel has been described as a beneficial diffusion controlling medium in an elegant titration experiment by *Mitrev et al.*^[115] Agar consists of linear agarose and agaropectin molecules dispersed between the agarose chains. When dissolved in water a so-called heterogeneous gel is formed in which strong inter-polymer interactions hinder the polymer chains from extensive movements.^[108b] The polymer is quasi static and spin relaxation is very fast, which is the origin of the NMR *silence* as observed by *Mitrev et al.* In contrast, PS is homogeneous and its polymer chains move less restricted effecting the residual polymer signals as observed in section 3.2.

According to Scheme 3-6 samples were prepared by first immersing agar in D₂O within an NMR tube. The suspension was then heated until a homogeneous gel resulted. An equimolar mixture of L-aspartic acid, L-His, L-Tyr and L-Val in slightly alkaline D₂O was added on top of the gel sample and SSE spectra were recorded immediately afterwards as well as 2, 6 and 10 h later. This procedure was repeated for gel samples with varying amounts of agar (weight-percentages of 2, 3 and 5wt%), details on sample preparation and SSE parameters are described in the experimental section 5.1).

All spectra were then processed and integrals of distinct signals (L-Asp 2.20 ppm, L-His 7.54 ppm, L-Tyr 6.48 ppm, L-Val 0.79 ppm) in each slice were measured and corrected using homogeneous reference samples (Figure 3-15).

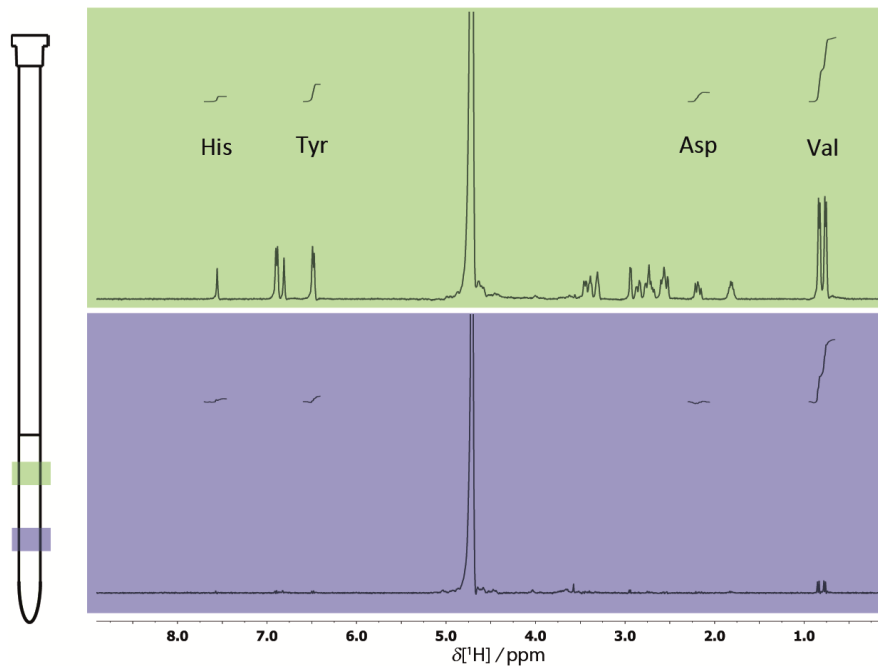


Figure 3-15: Spectra of slice 3 (top) and 10 (bottom) in the 3wt% agar sample as obtained 6h after preparation of the sample. From left to right integrals are marked for: His, Tyr, Asp and Val.

These spectra already show a pronounced decrease of signal intensity from top to bottom. Especially the signal for L-Val can be easily examined by eye, and is thereby recognized as the fastest diffusing species of the amino acid mixture. To achieve a reliable quantitative evaluation for the diffusion of the observed amino acids, concentration profiles were prepared from the measured intensities. Therefore, the integral I_0 of the slice $z_0 = 0$ mm which corresponds to the boundary between agar and neat amino acid mixture in D_2O was determined. Practically, z_0 was chosen as the slice that exhibits half of the maximum signal intensity of all SSE spectra obtained immediately after preparation. At this point, amino acid diffusion into the gel is still negligible.

Since integral I and concentration c are proportional, their ratios I/I_0 and c/c_0 involving the integral and concentration at the boundary (I_0 and c_0), respectively, is equal. Thus, profiles of the normalized concentration c/c_0 against the slice z can be directly prepared from the integrals (Figure 3-16).

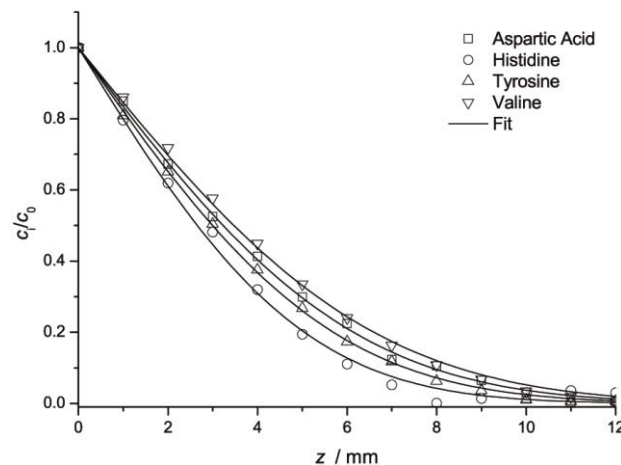


Figure 3-16: Concentration profile c/c_0 along slices z for the amino acids mixture in the 3wt% agar sample as obtained from spectra recorded 6 h after preparation of the sample. The fit is according to equation (3-3).

These profiles were fitted to obtain diffusion coefficients D_{SSE} via equation (3-3) which is the solution of Fick's 2nd law^[120] as obtained by Laplace transformation for the given boundary conditions.^[121]

$$\frac{c}{c_o} = \text{erfc}\left(\frac{z}{2\sqrt{D_{\text{SSE}}t}}\right) \quad (3-3)$$

Here, erfc is the complementary error-function and t is the elapsed time since the diffusion into the gel has begun. The obtained values for three samples with increasing amount of agarose are summarized in Table 3-4.

Table 3-4: Average diffusion coefficients \bar{D}_{SSE} and standard deviations for samples with 2wt%, 3wt% and 5wt% of agar as obtained through SSE. Values obtained after 2, 6 and 10 h were averaged.

	$\bar{D}_{\text{SSE}}(2\text{wt\%}) / 10^{-10} \text{ m}^2\text{s}^{-1}$	$\bar{D}_{\text{SSE}}(3\text{wt\%}) / 10^{-10} \text{ m}^2\text{s}^{-1}$	$\bar{D}_{\text{SSE}}(5\text{wt\%}) / 10^{-10} \text{ m}^2\text{s}^{-1}$
Asp	7.0±0.5	5.6±0.3	5.2±0.7
His	4.1±0.2	3.6±0.2	3.4±0.2
Tyr	6.0±0.4	4.7±0.1	4.4±0.2
Val	8.0±0.7	6.4±0.4	5.7±0.4

The obtained trend of diffusion coefficients is in accordance with the molecular weight of the observed amino acids (Val < Asp < His < Tyr) yet interactions between amino acid and the gel certainly also induce effects that allow a *tailoring* of the chromatographic effect of agar: diffusion coefficients from the SSE approach differ sufficiently for all evaluated amino acids proving the anticipated chromatographic separation successful. For 2wt% of agar, coefficients range from $4.1 \cdot 10^{-10} \text{ m}^2/\text{s}$ for L-His to $8.0 \cdot 10^{-10} \text{ m}^2/\text{s}$ for L-Val (maximum difference $\Delta_{\text{max}} 3.9 \cdot 10^{-10} \text{ m}^2/\text{s}$). This gap is reduced for higher percentages of agar e.g. for 5wt% corresponding values range from $3.4 \cdot 10^{-10}$ to $5.7 \cdot 10^{-10} \text{ m}^2/\text{s}$ ($\Delta_{\text{max}} 2.3 \cdot 10^{-10} \text{ m}^2/\text{s}$). Hence, the separation of components can be adjusted through different amounts of agar.

In general, the obtained diffusion coefficients decrease with increasing agar weight percentage. Since higher amounts of agar should lead to higher gel densities and viscosities, this is in accordance with the *Stokes-Einstein* equation (diffusion coefficient is inversely proportional to viscosity).^[32] Additionally, there may be numerous further parameters, which influence the diffusion within the heterogeneous agar gel e.g. the order and type of crosslinking (recall that agar is a heterogeneous gel with strong inter polymer interaction in contrast to a homogeneous gel such as PS).^[108b] This property was briefly examined in a small variational study on the relaxation behavior of molecules within the polymers by recording transverse relaxation times T_2 of H₂O in samples of 2, 3 and 5wt% agar with *Carr-Purcell-Meiboom-Gill* (cpmg) pulse sequences^[122] (experimental details on sample preparation and parameters for the cpmg experiment are described in the experimental section 5.2). It is well-known that the pore size of (mostly inorganic) porous media (e.g. zeolites and metal organic frameworks) and the relaxation time (both longitudinal, T_1 , and transverse, T_2) of probe molecules within the pores are correlated by equation (3-4).^[123]

$$\frac{1}{T_2} = \frac{1}{T_{2,\text{bulk}}} + \delta \frac{S}{V} \quad (3-4)$$

Here $T_{2,\text{bulk}}$ is the relaxation time of molecules in bulk solution (molecules in the center of a pore), δ is the surface relaxivity, S is the surface area of the pore and V is the volume of the pore. In the case of a spherical pore the equation is reduced to a linear dependence between T_2 and the pore radius. Hence, molecules in small pores exhibit shortened relaxation times. Figure 3-17 depicts how the signal intensities I_{cpmg} decay during t through transverse relaxation.

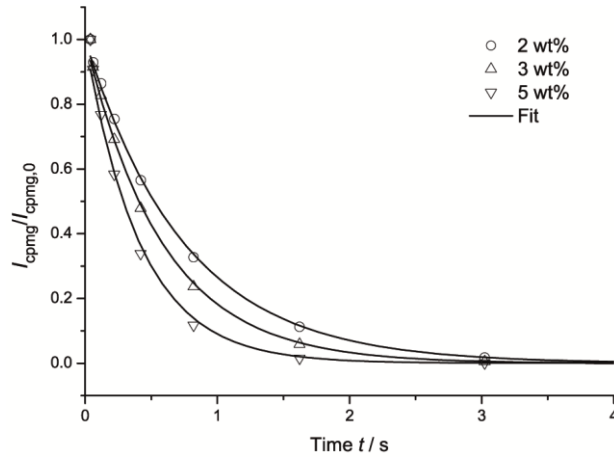


Figure 3-17: Integrals obtained from cpmg T_2 measurements of water in samples with 2wt%, 3wt% and 5wt% of agar versus elapsed time t (details are described in the experimental section 5.2).

The decay has been fitted by the mono exponential equation (3-5) in which $I_{\text{cpmg},0}$ is the signal intensity at $t = 0$.^[67]

$$\frac{I_{\text{cpmg}}}{I_{\text{cpmg},0}} = \exp\left(-\frac{t}{T_2}\right) \quad (3-5)$$

Values are reported in Table 3-5.

Table 3-5: Transverse relaxation times T_2 for water protons in samples with 2wt%, 3wt% and 5wt% of agar as obtained through cpmg pulse sequences.

	2wt%	3wt%	5wt%
T_2 (H ₂ O) /s	0.75±0.02	0.59±0.02	0.41±0.02

The T_2 values steadily decrease with increasing agar weight percentage. Accordingly, the pore size decreases towards higher amounts of agarose which is equivalent to an increase in crosslinking. Yet, these relaxation times should be treated with care since equation (3-4) also indicates that the mono-exponential decay as employed to fit the relaxation data could be insufficient. Within heterogeneous porous media a variety of pore sizes exist and the decay must be multi-exponential. This requires more sophisticated mathematical methods such as the inverse *Laplace* transformation to obtain relaxation time distributions.^[124] These methods have been successfully employed in the analysis of porous media and offer insights that may help for the characterization of polymeric gels and likewise. However, these methods also require extremely careful application are therefore beyond the scope of this work.^[125]

A final remark on values listed in Table 3-4 must address the determination of z_0 and its implications. This procedure was not fully adapted to e.g. the concave gel surface which is negligible for the *relative* discrimination between various diffusing species, but the correct assignment of z_0 has a strong impact on the obtained *absolute* diffusion coefficients. For instance using $z_1 = 0$ (i.e. assigning the slice below z_0 as the boundary slice with $I/I_0 = 1$, *vide supra*) gives $\bar{D}_{\text{SSE}} 5.5 \cdot 10^{-10} \text{ m}^2/\text{s}$ for L-Asp acid in 2wt% agar samples. This considerably deviates from $7.0 \cdot 10^{-10} \text{ m}^2/\text{s}$ as listed in Table 3-4. The literature value D_{Lit} of L-Asp acid in H_2O is $7.4 \cdot 10^{-10} \text{ m}^2/\text{s}$ at a starting concentration of 0.030 mol/l .^[126] If it is taken into account that the diffusion coefficient decreases with increasing concentration^[126] (Lit.: 0.030 mol/l ; SSE: 0.037 mol/l) and viscosity^[32] (Lit.: H_2O ; SSE: $\text{H}_2\text{O}/\text{agar}$ gel matrix), the obtained values are eventually in remarkably good accordance with the literature values.

For the evaluation of the separation performance of the SSE approach, the results discussed above were compared to the matrix assisted DOSY technique. Therefore, standard DOSY spectra of each gel sample were recorded one week after preparation to deduce diffusion coefficients (experimental details on DOSY acquisition and processing are described in the experimental section 5.2 and 5.3). Results are listed in Table 3-6.

Table 3-6: Average diffusion coefficients \bar{D}_{DOSY} and standard deviations in units of $10^{-10} \text{ m}^2/\text{s}$ for samples with 2wt%, 3wt% and 5wt% of agar as obtained through DOSY.

	$D_{\text{DOSY}}(2\text{wt\%}) / 10^{-10} \text{ m}^2\text{s}^{-1}$	$D_{\text{DOSY}}(3\text{wt\%}) / 10^{-10} \text{ m}^2\text{s}^{-1}$	$D_{\text{DOSY}}(5\text{wt\%}) / 10^{-10} \text{ m}^2\text{s}^{-1}$
Asp	4.3 ± 0.2	4.1 ± 0.3	3.7 ± 0.4
His	4.2 ± 0.1	4.4 ± 0.2	4.2 ± 0.1
Tyr	3.5 ± 0.1	3.6 ± 0.1	3.5 ± 0.1
Val	5.0 ± 0.1	4.9 ± 0.1	4.3 ± 0.1

Absolute values \bar{D}_{SSE} and D_{DOSY} differ since the former describes chemical diffusion and the second tracer diffusion (which only become identical at infinite dilution). In contrast to SSE, the absolute values D_{DOSY} are not reliable since HDO/ H_2O in the homogeneous reference solution was determined to diffuse at a rate of $1.71 \cdot 10^{-9} \text{ m}^2/\text{s}$ (obtained through analysis of the solvent residual signal). This deviates strongly from the literature value of $D_{\text{Lit}} = 2.30 \cdot 10^{-9} \text{ m}^2/\text{s}$.^[98b] Note, that the measurement of absolute diffusion coefficients via DOSY is a delicate task that depends on various experimental parameters (temperature, viscosity and likewise) and therefore, the majority of DOSY applications employ normalized/corrected diffusion coefficients that are obtained through the addition of reference substances.^[34, 36] However, correction of the values listed in Table 3-4 with a factor obtained from the HDO/ H_2O reference ($k = D_{\text{Lit}}/D_{\text{DOSY,Exp}} \approx 1.35$, $D_{\text{DOSY,Cor}} = D_{\text{DOSY}} \cdot k$) does not lead to a higher compliance of values \bar{D}_{SSE} and $D_{\text{DOSY,Cor}}$.

The essential conclusion drawn from these results is that the separation performance by D_{DOSY} is significantly worse than \bar{D}_{SSE} . For example for 2wt% of agar, coefficients range from $D_{\text{DOSY}} 3.5 \cdot 10^{-10} \text{ m}^2/\text{s}$ for L-Tyr to $5.0 \cdot 10^{-10} \text{ m}^2/\text{s}$ for L-Val. The corresponding difference $\Delta_{\text{max}} 1.5 \cdot 10^{-10} \text{ m}^2/\text{s}$ is lower than for \bar{D}_{SSE} ($\Delta = 2.0 \cdot 10^{-10} \text{ m}^2/\text{s}$) and much lower than the biggest difference for \bar{D}_{SSE} between L-His and L-Val as already pointed out above ($\Delta_{\text{max}} = 3.9 \cdot 10^{-10} \text{ m}^2/\text{s}$). In contrast to SSE, values D_{DOSY} for L-His and L-Asp acid (2wt%, 3wt%) as well as L-histidine and L-Val (5wt%) are practically indistinguishable.

Additionally, increasing weight percentage has a discernable impact on D_{DOSY} of L-Asp acid and L-Val, while for L-His and L-Tyr no changes occur at all. This is remarkable since the latter amino acids share only very few characteristics. Elucidation of this phenomenon must be based on further investigations. In particular, molecular dynamics simulations have proven to be able to probe the causes of specific diffusivities.^[127]

3.5 Summary and Outlook Slice-Selective Excitation

In the studies described above it could be shown that slice-selective NMR spectroscopy is an effective tool to perform

- i) *in situ* observation of reactions,
- ii) single-shot NMR titrations,
- iii) qualitative NMR chromatography.

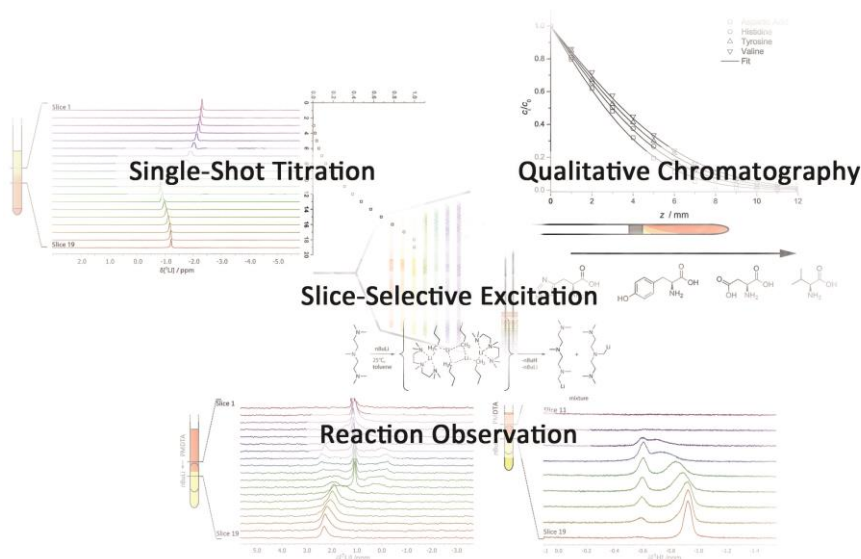


Figure 3-18: SSE on heterogeneous samples and its possible applications: i) *In situ* observation of reactions. ii) Single-shot NMR titration. iii) Qualitative NMR chromatography.

Reaction observation by slice-selective NMR spectroscopy was tested in the reaction between PMDTA and $n\text{BuLi}$, where the previously characterized intermediate $[(n\text{BuLi})_2\text{PMDTA}]_2$ could be identified.^[104] A stretched PS gel was used as medium to slow down the reaction and avoid convection, immobilize one reactant ($n\text{BuLi}$) with respect to the other (PMDTA) and principally enable the observation of anisotropic parameters as additional source of structural information. The magnitude of ^7Li RQCs can be obtained with ease from 1D SSE experiments and qualitative values have already proven to amend structure determination. Future developments should focus on the combination of SSE and 2D techniques such as the CLIP-HSQC to facilitate RDC structure determination (see section 4.2) of intermediates and derive the alignment which allows quantitative RQC analysis (see section 4.3). Additionally, the swelling process of PS gel as a more specific problem and future projects concerning this topic were addressed throughout the discussion. It is especially remarkable that the PS alignment

strength shows a response at the reaction front between PMDTA and *n*BuLi. This may result from the reaction heat from the lithiation and can be examined in future studies.

Single-shot titration was successfully tested in the complexation of a lithium salt by 12-crown-4, where a conventional ^7Li chemical shift titration curve could be reproduced that indicates an initial strong coordination of Li^+ by one crown ether followed by a weaker coordination to give the double decker type complex. It has also been highlighted that this approach was greatly improved in subsequent studies by *Mitrev et al.* who utilized agar gel as diffusion controlling medium.^[115] Note that a further advantage of the gel method (both in titration and reaction monitoring applications) is the possibility to adjust the width of the reaction front and hence to *zoom* into the interesting region. Hence, it will be possible to apply slice-selective NMR titration to a broad range of reactions to obtain information about mechanisms as well as stoichiometry of complexes and products. Once again, this approach would greatly benefit from a combination of SSE with more elaborate 2D NMR techniques. For instance, DOSY and NOESY are powerful methods to characterize complexes.^[24, 26-27, 33, 36, 50, 62d, 128] Their (fast) SSE variants would significantly enhance the quantity and quality of information.

Qualitative chromatography was illustrated in a study with agar as stationary phase and the analytes L-Asp acid, L-His, L-Tyr and L-Val. It was demonstrated that in this particular case the discrimination through SSE is superior to DOSY.^[118a, 118b, 118d] Therefore, the SSE approach to NMR chromatography is at least complementary to the established DOSY method and should be considered in cases where DOSY gives ambiguous results. Different amounts of agar gel were evaluated as stationary media and their separation behavior was attributed to pore sizes and the amount of crosslinking (through physical entanglement as well as electrostatic interactions). Support for these hypotheses was found by a small variational study on T_2 relaxation data. Relaxation data were obtained by standard cpmg pulse sequences and in general seem to be valuable observables for the examination of swollen polymeric gels. Future applications could determine the extent to which alignment parameters and relaxation rates are correlated. Diffusion measurements (both through SSE and DOSY) will amend this approach as well. In general, combination of SSE chromatography and 2D NMR techniques can serve once more as a basic but powerful method for the characterization of mixtures. Note that the use of chiral agarose, as of other chiral alignment media, principally enables the separation of racemic mixtures.^[129]

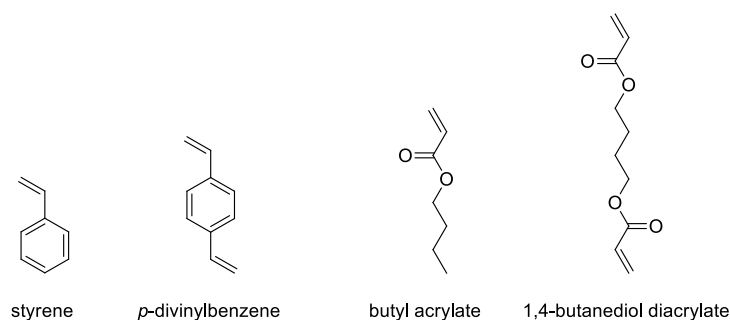
Here, further studies could aim at the evaluation of this approach in combination with molecular dynamics simulations of diffusion to enable full enantiomer discrimination. The methods shown above will add further impetus to research in slice-selective NMR spectroscopy of heterogeneous samples in particular polymeric gels.

4 ANISOTROPIC NMR SPECTROSCOPY IN POLYMERIC GELS

4.1 Alignment of Small Symmetric Molecules

Weak alignment of molecules for NMR experiments is nowadays achieved by a wide range of anisotropic media. The most common media are nematic liquid crystals^[130] and mechanically deformed gels.^[70a, 91] Originally, research started within liquid crystals and these systems have been studied very thoroughly ever since. Today, comprehensive theoretical treatment of nematic liquid crystals offers unprecedented insights and enables a comprehensive interpretation of alignment phenomena within these media.^[80, 131]

Theoretical estimations of alignment has been achieved through steric models^[81] and first principle approximations on inertia^[79e] or gyration^[79e] of molecules. These models sufficiently predict the alignment of large molecules and are therefore commonly incorporated in RDC analysis programs.^[81, 83] Meanwhile, for small molecules only very little is known about the mechanism of alignment.^[132] It is obvious that e.g. the influence of sparse electrostatic interactions must be reviewed in the case of small molecules to supplement purely geometric approaches. In this thesis, the weak alignment of benzene derivatives and polycyclic homologues within crosslinked stretched PS and polybutylacrylate (PBA) was examined to obtain an improved understanding of small molecule alignment in polymer gels (Scheme 4-1).



Scheme 4-1: (Atactic) PS is thermally polymerized from styrene with covalent crosslinking between PS polymer strands through divinylbenzene (technical grade: 80% *p*-divinylbenzene). For (atactic) PBA, thermal polymerization of butyl acrylate and 1,4-butanediol diacrylate as cross-linker is performed.

In particular, the alignment of toluene, benzoic acid, methyl benzoate, 4-methylphenol and naphthalene in dichloromethane (DCM)/PS and DCM/PBA were evaluated. First $^1J_{\text{CH}}$ couplings were measured for each molecule by CLIP-HSQC experiments of isotropic samples (in neat DCM). Subsequently, anisotropic samples of these substances in PS and PBA were prepared. Swelling and sufficient equilibration of PS and PBA in DCM was achieved after four and three days, respectively.^[75b] Further CLIP-HSQC measurements yielded couplings $^1T_{\text{CH}}$ from which one-bond RDCs were calculated ($^1D_{\text{CH}} = ^1T_{\text{CH}} - ^1J_{\text{CH}}$). Details on sample preparation and CLIP-HSQC parameters are described in the experimental sections 5.1, 5.2 and 5.3. The molecular geometry can be conveniently described by polar

coordinates. Therefore, the molecules were set into coordinate systems by aligning their C_2 axes (all functional groups are considered to rapidly rotate as to fulfill this symmetry) along z . Naphthalene was aligned along z by its *long* C_2 axis. Furthermore, the ring planes were embedded within the y - z plane. Polar angles ϑ were then determined between the z axis and the respective C–H (C–CH₃) bond vector (Figure 4-1).

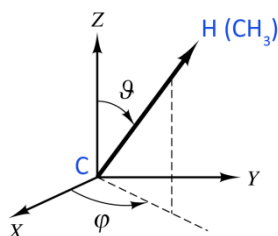


Figure 4-1: Polar angle ϑ and azimuthal angle φ for an arbitrary C–H bond vector.

The parallel arrangement between the ring and the y - z plane leads to $\varphi = 90^\circ$ for all bonds. RDCs and associated distances and angles are summarized in Table 4-1.

Table 4-1: Textbook C–H Distances R_{IS} and polar angles ϑ between the z axis (green arrow) and the particular C–H (C–CH₃) vectors (note that azimuthal angles are 90° since the ring is embedded within the y - z plane; red and green arrows). Experimental RDCs for the corresponding bonds of structures aligned in PS/DCM as well as PBA/DCM.

		$R_{\text{IS}} / \text{\AA}$	$\vartheta / ^\circ$	$^1D_{\text{CH}}(\text{PS}) / \text{Hz}$	$^1D_{\text{CH}}(\text{PBA}) / \text{Hz}$	
Toluene		1–CH ₃	1.09	0 ^a	2.5±0.2 ^a	7.2±0.9 ^a
		C2–H2	1.09	60	–25.1±1.5	–9.1±0.6
		C3–H3	1.09	120	–23.6±0.7	–10.6±0.3
		C4–H4	1.09	180	–10.4±0.6	–19.7±0.3
Benzoic acid		C2–H2	1.09	60	–18.1±1.4	–5.8±0.4
		C3–H3	1.09	120	–18.9±0.6	–8.3±0.5
		C4–H4	1.09	180	–42.7±0.9	25.2±0.6
Methyl benzoate		C2–H2	1.09	60	–30.4±0.5	–9.6±0.2
		C3–H3	1.09	120	–29.2±1.6	–9.9±0.7
		C4–H4	1.09	180	–42.4±0.8	–21.8±1.7
4-Methylphenol		C2–H2	1.09	60	–31.2±0.7	13.1±0.8
		C3–H3	1.09	120	–31.7±3.6	13.8±0.7
		4–CH ₃	1.09	180 ^a	2.6±0.6 ^a	0.4±0.1 ^a
Naphthalene		C1–H1	1.09	90 ^c	–51.7±0.4	–8.8±1.7
		C2–H2	1.09	30 ^c	–22.5±1.2	–30.9±1.2

^a RDCs for rotating methyl groups can be treated as $^1D_{\text{CH}}$ with a corresponding vector along the C_3 symmetry axis of the methyl group by scaling the original value with -3 (e.g. toluene C1–Me then yields -7.5 Hz in PS).^[133]

The investigated molecules were treated like textbook structures with aromatic C–H distances of 1.09 Å and prevalent 30°, 60°, 90°, 120° and 180° angles between bonds concerning the aromatic moiety. Note

that the averaged RDCs which are measured for a methyl group can be easily transformed into a virtual C–H RDC pointing along the C–CH₃ threefold symmetry axis by applying a scaling factor of -3 . This factor stems from the more general formula (4-1) which is valid for a freely rotating group leading to averaged RDCs.

$$\langle D \rangle = D_{||} \frac{3 \cos^2 \phi - 1}{2} \quad (4-1)$$

Here $\langle D \rangle$ represents the averaged RDC, $D_{||}$ would be the RDC of a virtual C–H bond pointing along the axis of the rotation and ϕ is the angle between this axis and each *rotating vector* of the methyl group. In the case of ideally tetrahedral angles H–C–H and C–C–H within the methyl group, ϕ is 109.5° and the (inverse) scaling factor becomes $-1/3$. Usually, angles are non-ideal and hence the scaling factor is different from $-1/3$. This becomes visible if collinear C–H vectors exist which then display couplings that are beyond the virtual RDC as in the case of toluene (*vide infra*). In fact, these considerations have allowed the investigation of subtle differences of angles in methyl as well as methylene groups.^[133-134] In the cases shown above, the methyl RDC for toluene was dismissed, since the collinear RDC for C4-H4 is much better determined. However, for 4-methylphenol the obtained methyl RDC is crucial and must be included for further investigations.

The coordinate system chosen above is very convenient, since its axes coincide with the molecular symmetry elements. Likewise, the PAS of the alignment tensor \mathbf{A}_{PAS} must contain existing symmetry elements. This means that the axes of the chosen coordinate system also span the PAS of the alignment tensor, i.e. the alignment tensor is diagonalized:^[64, 135]

$$\mathbf{A}_{\text{PAS}} = \begin{pmatrix} A_{\bar{x}} & 0 & 0 \\ 0 & A_{\bar{y}} & 0 \\ 0 & 0 & A_{\bar{z}} \end{pmatrix} \quad (4-2)$$

Thus, equation (1-7) can be transformed into polar coordinates and rewritten:

$$D = k(A_{\bar{x}} \sin^2 \vartheta \cos^2 \varphi + A_{\bar{y}} \sin^2 \vartheta \sin^2 \varphi + A_{\bar{z}} \cos^2 \vartheta) \quad (4-3)$$

For each molecule, the measured RDCs and corresponding angles and distances can then be employed to constitute a linear equation system to which the $A_{\bar{x}} + A_{\bar{y}} + A_{\bar{z}} = 0$ can be added. Therefore, only two independent alignment parameters remain. These essential effects of molecular symmetry on the number of independent alignment tensor elements are additionally summarized in Table 4-2:^[135]

Table 4-2: Independent elements of the alignment tensor for various point groups.

Point group	Independent alignment tensor elements
C_1, C_i	$A_{zz}, A_{xy}, A_{xz}, A_{yy}, A_{yz} (= 5)^a$
C_2, C_{2h}, C_s	$A_{zz}, A_{yy}, A_{xy} (= 3)^a$
C_{2v}, D_2, D_{2h}	$A_{zz}, A_{yy} (= 2)^a$
$C_{>2}, C_{>2v}, C_{>2h}, D_{>2}, D_{>2h}, D_{\geq 2d}$	$A_{zz} (= 1)^a$
T, T_d, O, O_h, K_h	–

^aNote the convention $|A_{\bar{x}}| \leq |A_{\bar{y}}| \leq |A_{\bar{z}}|$.

In general, at least five *independent* RDCs are necessary to derive an alignment tensor for molecules of arbitrary symmetry. In this work, *dependent* relates to two RDCs for parallel internuclear vectors or any more than three vectors located in one plane. According to Table 4-1, the number of *independent* RDCs exceeds the necessary independent alignment tensor elements for the observed molecules (e.g. toluene C_{2v} : two tensor elements, three *independent* experimental RDCs). These overdetermined equation systems can be solved by the least-square routine *numpy.linalg.lstsq* as included in the NumPy package for Python (script is added to the Appendix).^[136] Resulting alignment tensors for PS are displayed in Table 4-3.

Table 4-3: Alignment tensor elements $A_{\bar{x}}$, $A_{\bar{y}}$ and $A_{\bar{z}}$ for the set of molecules in swollen DCM/PS as determined by least-square techniques from overdetermined equation systems of (4-3) with values taken from Table 4-1. Additionally, the magnitude of these values and the orientation within the molecular frame are represented by arrows with corresponding length, their sign is represented by their color (blue: positive, red: negative).

	Toluene	Benzoic acid	Methyl benzoate	4-Methylphenol	Naphthalene
$A_{\bar{x}}/10^{-4}$	1.4	1.4	3.6	1.1	1.8
$A_{\bar{y}}/10^{-4}$	4.2	6.1	6.0	5.6	7.4
$A_{\bar{z}}/10^{-4}$	-5.6	-7.5	-9.6	-6.7	-9.2

Note, that solutions for linear equation systems of (4-3) were renamed according to the convention $|A_{\bar{x}}| \leq |A_{\bar{y}}| \leq |A_{\bar{z}}|$.^[64] Additionally, figures of the molecules are shown in which arrows indicate the magnitude and sign of the alignment tensor elements through length and color, respectively. Recall that the alignment tensor essentially describes the probability to find one of its principal axes oriented along the magnetic field relative to the probability in isotropic solution. A more positive principal alignment tensor element implies a higher probability to find this axis oriented along the magnetic field vector. In Table 4-3 molecules are depicted in the (exaggeratedly) preferred orientation of the molecule within the laboratory frame, i.e. with their highest positive alignment tensor element oriented along the magnetic field vector (vertical sheet axis). Hence, the orientation of toluene (i.e. its C_2 axis) in PS is found *perpendicular* or *orthogonal* to the magnetic field. Toluene is then also perpendicularly oriented towards the elongated pores within the gel that result from the swelling of the polymer. This is remarkable since these pores are usually assumed to be responsible for the alignment of the solutes according to simple

sterical/geometrical considerations. 4-methylphenol and naphthalene show an analogous mismatch between PS pore geometry and molecular shape.

Benzoic acid and methyl benzoate behave differently and show the allegedly steric orientation which will be referred to as *parallel* or *aligned* along the field/pore. Discussion of absolute alignment values is omitted, since RQCs of CD_2Cl_2 were not fully consistent for the five samples, indicating different alignment strengths. Hence, a purely qualitative classification is made. It can be assumed that the sum of i) generalized steric (repulsive, parallel) and ii) directional (attractive, orthogonal) interaromatic forces defines the alignment characteristics of these small aromatic molecules in PS. For toluene, 4-methylphenol and naphthalene the directional interaromatic forces prevail (Figure 4-2, green).

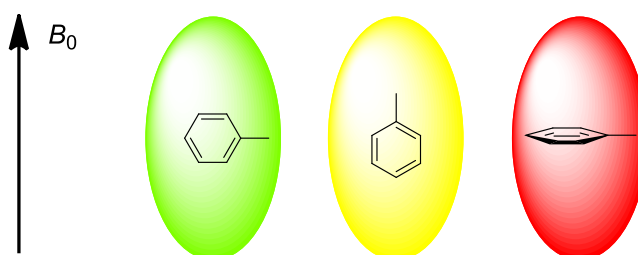


Figure 4-2: Schematic illustration of the most probable (green), less probable (yellow) and least probable orientation of toluene in PS.

Interactions between aromatic compounds are influenced by a wide range of factors that include electrostatic, hydrophobic and van-der-Waals interactions.^[137] Attribution of these factors to the situation encountered in the gel-solute mixture of small aromatic molecules and PS ring moieties requires very elaborate and extensive studies that are beyond the scope of this thesis. In particular, the orientation of the polymer ring moieties is not ascertainable by solution state NMR spectroscopic means. Therefore, common structural motifs like the stacked (also face-to-face or sandwich), offset stacked or T-shaped (also edge-to-face) configuration of aromatic rings, which in turn can be attributed to certain interactions, are not traceable.^[137]

The steric parallel force becomes clear from the second and least probable orientation (Figure 4-2 yellow and red, respectively). The orthogonal alignment of the ring plane towards the magnetic field (long pore axis) is never preferred, as the associated tensor element A_z is always the most negative tensor element (also see Table 4-3).

In contrast to toluene, in 4-methylphenol and naphthalene the parallel orientation of benzoic acid may be traced back to the formation of *long* hydrogen-bonded dimers which might overcome a tipping point at which steric demand exceeds the influence of interaromatic *orthogonal forces*. This hypothesis is limited since methyl benzoate obviously cannot form hydrogen-bonded dimers but is also in parallel orientation to the pore. However, instead of increasing the spatial extension (dimerization), reduction of the additional interaromatic attractive force may also lead to a prevalent parallel alignment, as is the situation for methyl benzoate: one difference between the molecules grouped into the parallel and orthogonal categories is the inductive effect of the substituents and the resulting electronic characteristics. The carboxylic acid and its ester are characterized by a $-M$ (mesomeric) and $-I$ (inductive) effect, whereas the other aromats are associated with a $+M$ and $+I$ or no effect. This can be

illustrated by electrostatic potential (ESP) maps, which were exemplarily obtained from B3LYP/ 6-31G(d,p) geometry optimization and subsequent analogous DFT calculations and the CHELPG method^[138] for toluene and benzoic acid. These surfaces are depicted in Figure 4-3.

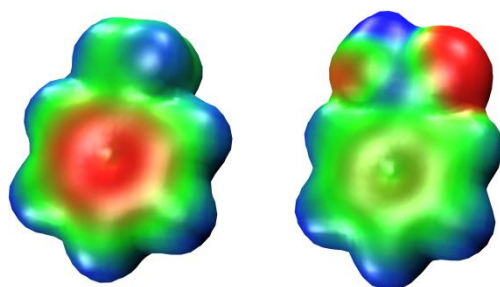

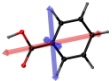

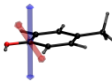


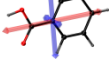

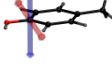








Figure 4-3: ESP maps for toluene (left) and benzoic acid (right) as derived from B3LYP/6-31G(d,p) DFT calculations and the CHELPG method. Red marks areas of high electrostatic potentials (more negative charge), blue corresponds to low electrostatic potentials (less negative charge).

The low electrostatic potential within the ring moiety may signal a decreased interaromatic interaction in the case of benzoic acid and methyl benzoate which allows domination of the steric forces that lead to parallel orientation.

In summary, only vague molecular features have been found that delimit the parallel from the orthogonal orientation: it is very likely that for small molecules the type and strength of (attractive) intermolecular interactions and more generalized steric (repulsive electrostatic) interaction are both responsible for alignment media phenomena. This hypothesis is supported by more distinct results for the orientation of aromats in PBA. The corresponding alignment tensor elements are summarized in Table 4-4.

Table 4-4: Alignment tensor elements $A_{\bar{x}}$, $A_{\bar{y}}$ and $A_{\bar{z}}$ for a set of molecules in swollen DCM/PBA as determined by SVD from overdetermined equation systems of (4-3) with values taken from Table 4-1. Additionally, the magnitude of these values and the orientation within the molecular frame are represented by arrows with corresponding length, their sign is represented by their color (blue: positive, red: negative).

	Toluene	Benzoic acid	Methyl benzoate	4-Methylphenol	Naphthalene
$A_{\bar{x}}/10^{-4}$	0.9 	1.1 	0.8 	0.2 	1.3 
$A_{\bar{y}}/10^{-4}$	2.8 	2.5 	3.1 	2.4 	5.4 
$A_{\bar{z}}/10^{-4}$	-3.7 	-3.6 	-3.9 	-2.6 	-6.7 

In PBA, the principal values of the alignment tensor for toluene, methyl benzoate and naphthalene indicate parallel orientation in agreement with steric considerations. In contrast, benzoic acid once more behaves differently and orients perpendicular towards the magnetic field as can also be observed to some extent for 4-methylphenol which has a very unique orientation. These latter compounds have the possibility to establish hydrogen bonds towards the ester group of PBA. Hence, the orientation of small molecules in PBA can be strictly derived from either a sterical match between pore and molecular shape (toluene, methyl benzoate, naphthalene), or from the molecule's ability to form hydrogen bonds

with the polymer (benzoic acid and 4-methylphenol). The results on small molecules alignment in polymeric gels have implications for applications of RDCs and RQCs in structure determination:

- i) *A priori* prediction of alignment based on simple sterical/mechanical models is often not sufficient for small molecules.
- ii) The interplay of attractive and repulsive interactions changes the energetic states of solutes within the alignment medium, similar to strong solvent effects. Hence, the population of conformers can be considerably biased compared to isotropic solutions.

The prediction of alignment will be addressed in the following subsection. The effect of conformation will be briefly discussed in section 4.2.

4.1.1 Molecular Dynamics Simulations

Starting from a proposed molecular structure one can either back-calculate the alignment tensor from experimental RDCs (*vide supra*) or estimate theoretical alignment from first principles such as the moment of inertia (Tramite)^[77] or the model's shape combined with steric obstruction (PALES).^[81] Prediction of RDCs by the latter means is successfully performed for larger biomolecules, but the results introduced above underline their limits in the case of small to medium sized molecules. In the last few years, some effort has been put into the study of the alignment of small organic solutes in liquid crystals by stochastic molecular dynamics (MD) simulations^[80, 131], but only very recently attempts were made to transfer this knowledge and simulate solutes in stretched polymeric gels. In an innovative study *Frank et al.* directly obtained averaged RDCs from trajectories of strychnine in the presence of an elongated PS strand.^[132] However, this approach must be followed by further benchmark studies to learn more about its potentials and limitations. Toluene and benzoic acid exhibited unusual, diverse alignment in the experiments described above and were therefore chosen to be evaluated in MD simulations of solute polymer systems within this thesis.

MD simulations were performed with GROMACS 5.1.2.^[139] The topology was built with the optimized potentials for liquid simulations/all-atom (OPLS) and tabulated standard OPLS charges (exemplary topology files are included in the Appendix).^[140] Similar to the original publication by *Frank et al.*, polymeric PS and PBA strands were mimicked by six atactic styrene and butyl acrylate units, respectively. First attempts revealed that the phenyl moiety at the methylene end of the strand was strongly bent away during MD runs (observed by spatial distribution functions, *vide infra*). Hence, in this work methyl groups were added to the methylene backbone end which gave a consistent orientation of the phenyl moieties. The strands were positionally restrained by adding strong isotropic forces on the methyl end groups. A typical energy optimized structure for PBA and benzoic acid prior to the start of the simulation is shown in Figure 4-4.

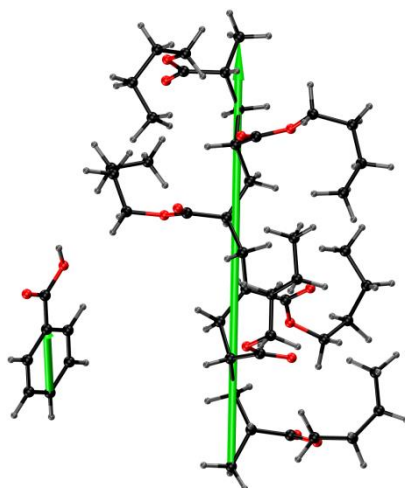


Figure 4-4: Energy-optimized starting position for the PBA strand and a benzoic acid molecule prior to the simulation. The green polymer strand arrow connects both positionally fixed methyl groups. The benzoic acid is characterized by an arrow parallel to the *quasi* C_2 (C1-C4) axis.

MD runs were performed in small cubic boxes with periodic boundary conditions and further options as suggested by *Frank et al.* Most notably, no solvent molecules were added, since the authors reported that only then alignment became non-zero during the production runs. To mimic solvent effects, *Langevin* dynamics were applied in which friction terms were added. Analogous to the original proceedings, long 400 ns production runs were executed directly after energy minimization. The *Berendsen* thermostat^[141] was set to a target temperature of 300 K but no preceding thermal equilibration run was conducted.^a Further details on the MD simulation protocol as well as analyzing tools are described in the experimental section (5.4). The production runs were analyzed by sampling $\cos^2\theta - 1/3$ where θ is the polar angle between C-H vectors and the vector defined by the two positionally restrained methyl carbons. This term was averaged over time to give $\langle \cos^2\theta - 1/3 \rangle$. The progression of this averaging for toluene and benzoic acid in PS is depicted in Figure 4-5.

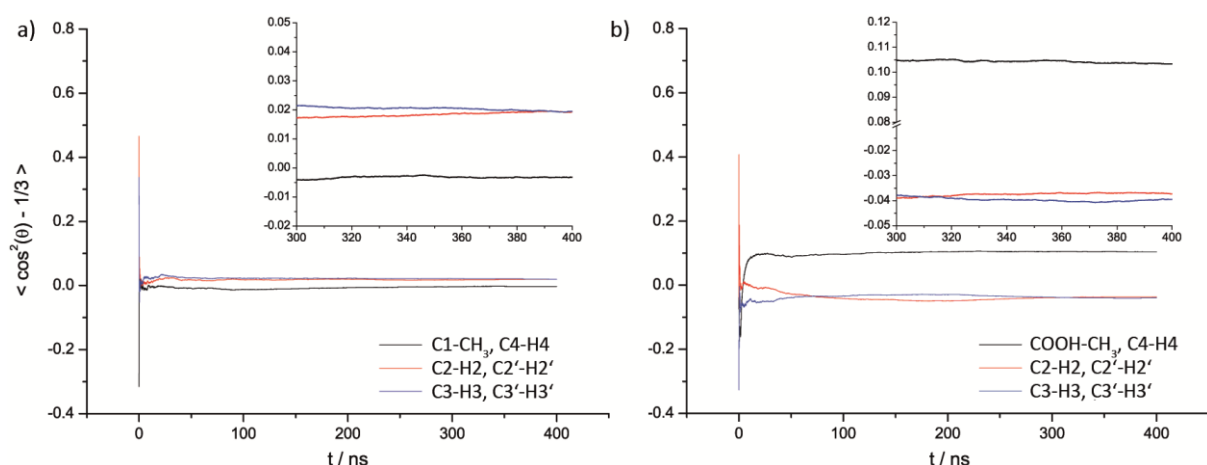


Figure 4-5: Convergence of the averaged angle term $\langle \cos^2\theta - 1/3 \rangle$ between C-H vectors for a) toluene and b) benzoic acid and the vector connecting the fixed methyl groups of the PS polymer strand backbone.

^a Pressure equilibration is not feasible due to the positional restraints. The extremely long duration of the production run covers up short thermal non-equilibrium contributions at the beginning of the run.

Figure 4-5 a) summarizes the averaged orientation of toluene in PS. Orientational equilibration, which is signaled by constant values $\langle \cos^2\theta - 1/3 \rangle$, is reached after about 100 ns. The angle between the polymer strand and C4-H4 (C1-CH₃) shows an overall negative value for $\langle \cos^2\theta - 1/3 \rangle$. For C2-H2 and C3-H3 positive values were obtained. This relationship is inverted for benzoic acid in Figure 4-5 b). Since more positive values for $\langle \cos^2\theta - 1/3 \rangle$ correspond to a more parallel orientation of the corresponding vectors and more negative values indicate predominantly orthogonal orientations (also see Figure 1-11 and corresponding equation (1-3)), these simulations qualitatively reproduce the experimental findings. As MD simulations are based on a set of classical forces that include electrostatic attraction and (steric) repulsion, the rationalization of orientation through shape and electronic effects is further corroborated. However, note that the utilization of standard OPLS charges implies that benzoic acid and toluene are only distinguished by the charges of their substituents while the ring charges are the same.^a Despite this deficiency, divergent experimental results were reproduced. This hints that interactions between the polymer and COOH/CH₃ play a critical role. The analogous treatment of both compounds in PBA yields further insights. Interestingly, these MD simulations first resulted in very weak (benzoic acid) or no (toluene) alignment. After reproducing these observations with different starting geometries, the MD simulation protocol was slightly varied by incorporating 50 DCM molecules.^b This gave pronounced orientations for benzoic acid in PBA as depicted in Figure 4-6.

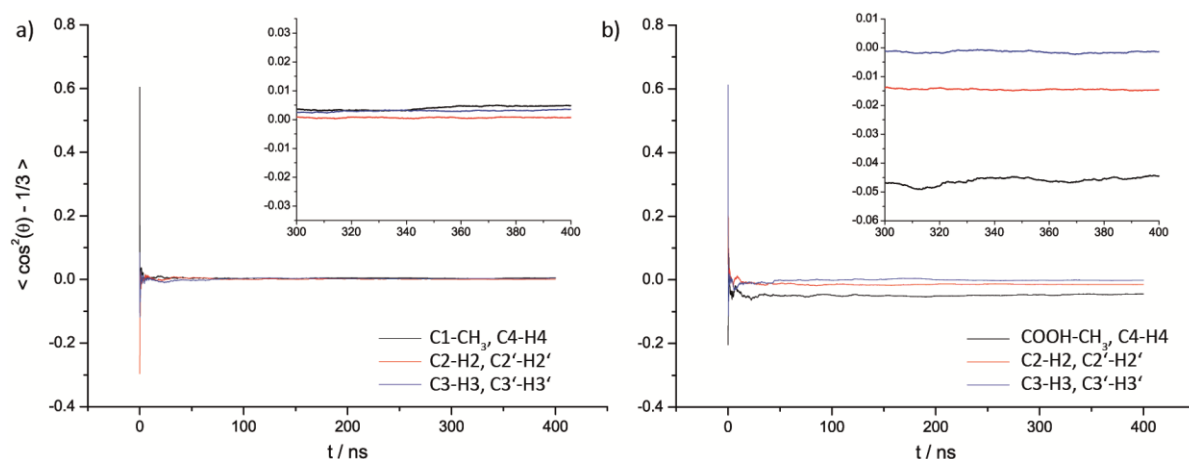


Figure 4-6: Convergence of the averaged angle term $\langle \cos^2\theta - 1/3 \rangle$ between C-H vectors for a) toluene and b) benzoic acid and the vector connecting the fixed methyl group of the PBA polymer strand backbone. The simulation additionally included 50 DCM molecules.

Toluene is still not noticeably ordered within PBA but the orientation of benzoic acid is in accordance with experimental results. The vector C4-H4 exhibits the most negative values for $\langle \cos^2\theta - 1/3 \rangle$ which corresponds to a predominantly orthogonal orientation between the molecule and the polymer strand.

^a Negative carbon and positive hydrogen charges (OPLS standard) should lead to attraction between polymer ring moieties and the solute aromatic rings.

^b Additionally the cut-off scheme for the treatment of molecules in adjacent simulation boxes (periodic boundaries) was changed from *group* to *verlet*. This is a technical modification to improve calculation performance as suggested in the updated user guidelines for GROMACS 5.1.2.^[139]

The MD run involving benzoic acid and PBA/DCM was then further analyzed by spatial distribution functions (SDFs, Figure 4-7).^[142]

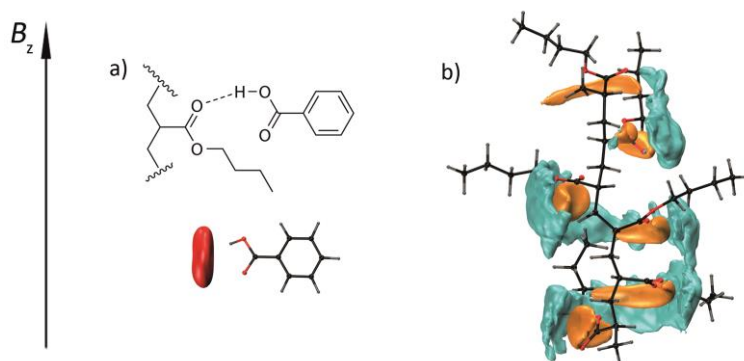


Figure 4-7: a) Top: Scheme of a possible hydrogen bond between a butyl acrylate unit of the polymer and benzoic acid. Bottom: SDF of the six ester oxygen molecules (=O) around the benzoic acid molecule (reference frame) at an isolevel of 3 \AA^{-3} . b) SDF of the center of mass of the benzoic acid (cyan, 2 \AA^{-3}) and the six ester oxygen molecules (=O) (orange, 7 \AA^{-3}) with the simulation cell as reference frame.

These functions are a measure of the occurrence of an observed particle in a chosen reference frame. SDFs can be illustrated by 3D models which resemble the space in which the particle occurs (during the MD run) at least as often as a specified isolevel value indicates. Figure 4-7 a) bottom depicts an SDF (red) in which benzoic acid is taken as reference frame and the spatial distribution of six ester oxygen atoms (=O) is observed. The SDF is shown at an isolevel of 3 \AA^{-3} which shows that during the MD run the carbonyloxygen occurs more than three times per \AA^3 within the red volume (less than three times per \AA^3 outside of the volume). In b) the simulation cell is taken as reference and the six ester oxygen atoms (=O, orange) and the center of mass of the benzoic acid (cyan) are depicted at an isolevel of 7 \AA^{-3} and 2 \AA^{-3} , respectively. The SDFs in Figure 4-7 a) and b) clearly indicate the existence of hydrogen bonds by the prevalent spatial association between PBA carbonyl oxygen and benzoic acid. However, these analyses do not directly correlate hydrogen bonds with the preferred orientation of the molecules. This is achieved through so-called combined distribution functions (CDFs).^[142] In this study a CDF of a radial and an angular distribution function is derived by sampling at a time i) the distance between the carboxylate oxygen atoms of the polymer and the carboxylic hydrogen atom of the benzoic acid and ii) the associated angle between the benzoic acid C1-C4 axis and the polymer strand axis (also see Figure 4-4). The CDF of these two functions is given as a heat map in Figure 4-8.

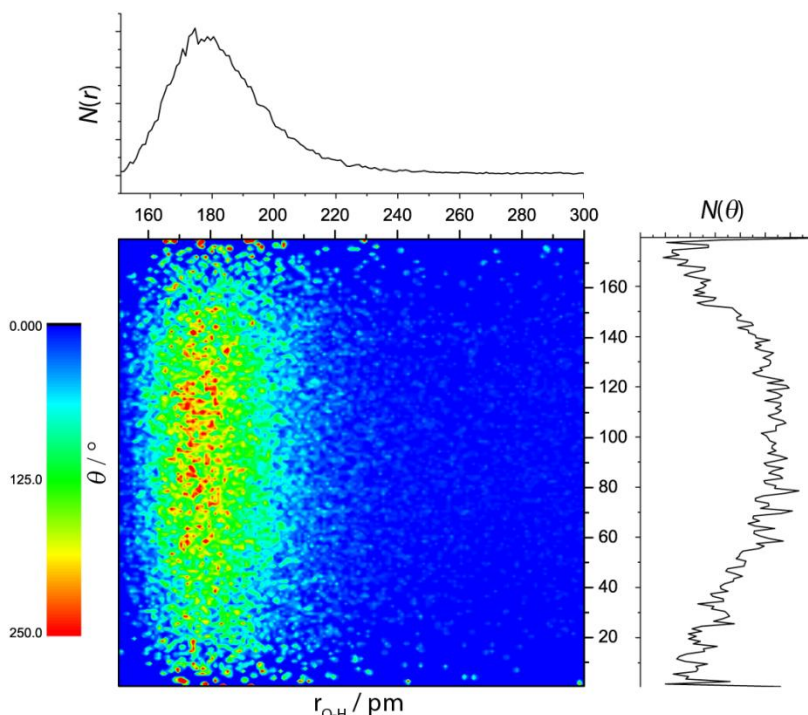


Figure 4-8: CDF for benzoic acid. Horizontal: radial distribution function describing the abundance of a distance between the carboxylic hydrogen of benzoic acid and the carboxylate oxygen of PBA. Vertical: angular distribution function describing the abundance of the angle between the C1-C4 axis (*longitudinal*, quasi C_2) of benzoic acid and the polymer strand. Color-coding is according to the prevalence of the angle-distance combination.

The horizontal scale of the CDF represents the radial distribution and illustrates that short distances between 160 and 200 pm are prevalent. These are typical values for O-H...O hydrogen bonds.^[143] These values are additionally correlated with an increased occurrence of angles around 90° between polymer strand and the C1-C4 benzoic acid axis as shown by the vertical scale which represents the angular distribution. In conclusion, it is very likely, that the orientation of the molecule not only correlates with the existence of hydrogen bonds, but is caused by this directional force.

This highlights the possibility to estimate molecular alignment *a priori* and furthermore illustrates computational methods on how to analyse the alignment phenomena and derive the underlying mechanisms. Admittedly, the potential of MD simulations is not fully exploited in the approaches described above and experimental findings were only qualitatively reproduced. Advanced numerical simulations of alignment in nematic liquid crystals involve large ensembles of up to several thousand molecules (e.g. *Pizzirusso et al.* 2000 molecules with a liquid crystal molecules to solute ratio of 95:5) but which are only covered over a short period of time (e.g. *Pizzirusso et al.* 40 ns). In principle simulation of a smaller ensemble over a longer period of time is equivalent according to the ergodic hypothesis. However, it is likely that the small fixed styrene strand is only a very simplified reproduction of the real situation within the gel and certainly introduces some bias. Furthermore, improvements such as the accurate charge determination of atoms within single polymer units and incorporation of solute molecules are mandatory. In summary, these steps should lead to a drastically improved quantitative agreement between experimental and theoretical alignment. In this case, MD simulations of small

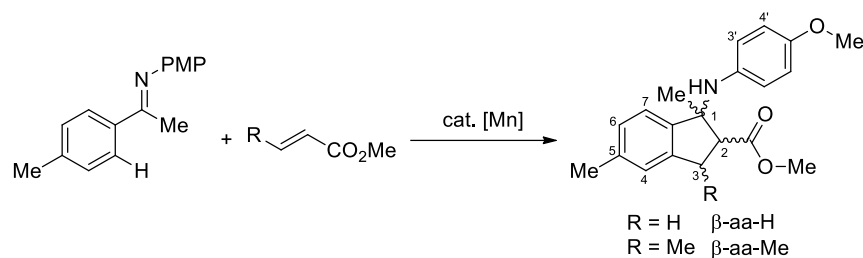
molecules in stretched gels will add valuable information on the alignment and open up tremendous possibilities for further studies. One example is the examination of lithium organic aggregates: in standard solution NMR one cannot distinguish between various species from their chemical shift. However, aggregates align very differently and comparison of experimental and computational results can be a powerful method to evaluate the aggregation motif. Similarly, enantiomers cannot be distinguished in homogeneous solution but two enantiomers in chiral alignment media may give two discernible orientations and allow absolute structure determination through two different sets of RDCs.

4.2 Structure Elucidation of Synthetic Indanes

The increasing variety of small molecules, whether synthesized in the laboratory or found as natural products, requires suitable methods for their structural characterization. While the chemical constitution (i.e. the connectivity of atoms) of a molecule is often known, the determination of the configuration and conformation (i.e. the 3D structure) can become an intricate task. X-ray crystallography has the ability to doubtlessly solve these stereochemical problems but requires well grown crystals and is usually not able to make dynamical structural changes observable.^[144] Alternatively, sophisticated NMR techniques which rely on NOE^[145] as well as the analysis of *J*-couplings^[146] are applied. However, these NMR techniques face several disadvantages since they are limited to locally well-connected spins (through bonds or spatial proximity) or depend on complex parameterized *Karplus* curves.^[59, 147] Additionally, even if NOE and *J*-coupling analysis is possible, these may not resolve ambiguities that arise from unfavourable combinations of configuration and conformation as well as from fast conformational averaging. Insufficient distance and angle information from conventional NMR data can be complemented by incorporating structure verification via RDCs which involves five basic steps.^[66, 70-71, 88]

- i) Experimental determination of RDCs,
- ii) Generation of a comprehensive set of possible molecular structures,
- iii) Determination of alignment tensors from RDCs and structural proposals,
- iv) Back-calculation of RDCs from structural proposals and alignment tensors,
- v) Assessment of the agreement between back-calculated and experimental RDCs.

These steps are further explained in the following analysis of two β -amino acid esters^[148] with previously unknown configuration. β -amino acid esters are frequently employed in the synthesis of non-natural β -peptides, and are versatile intermediates in organic synthesis.^[149] Scheme 4-2 shows the reaction in which a manganese-catalyzed C–H functionalization/alkene annulation of ketimines led to the formation of β -amino acid esters methyl-1-((4-methoxyphenyl)amino)-1,5-dimethyl-2,3-dihydro-1*H*-indene-2-carboxylate (β -aa-H) and methyl-1-((4-methoxyphenyl)amino)-1,3,5-trimethyl-2,3-dihydro-1*H*-indene-2-carboxylate (β -aa-Me) comprising a prominent indane scaffold.^[148]



Scheme 4-2: Formation of $\beta\text{-aa-H}$ and $\beta\text{-aa-Me}$ via a manganese-catalyzed C–H functionalization/ alkene annulation. PMP = *para*-methoxyphenyl.

Note that up to three stereogenic centres are generated in an instant out of two achiral reactants. While the constitutional elucidation of $\beta\text{-aa-H}$ and $\beta\text{-aa-Me}$ from common 1D and 2D NMR procedures was straightforward, the derivation of the relative configuration at C1, C2 and C3 from NOE and J -coupling data was indistinct. Both, hydrogen distances and dihedral angles, rely on the conformation of the 5-membered ring which was likewise not known. For example, a pseudo-*equatorial* substituent on C1 may actually be closer in space to a *trans* compared to a *cis* configured substituent on C3 and vice versa (also see NOE and J -coupling data in the Appendix). The structure verification by RDCs is then performed by five consecutive steps that will be described in detail.

i) Experimental determination of RDCs

First, $\beta\text{-aa-H}$ and $\beta\text{-aa-Me}$ were dissolved in CDCl_3 and isotropic $^1J_{\text{CH}}$ couplings were measured by CLIP-HSQC experiments. This gave 12 $^1J_{\text{CH}}$ couplings for $\beta\text{-aa-H}$ as well as $\beta\text{-aa-Me}$. Alignment of the compounds was achieved by immersing a cylindrical PS stick (crosslinked through 0.3 vol% divinylbenzene) in solutions of $\beta\text{-aa-H}$ and $\beta\text{-aa-Me}$. Swelling/equilibration over four days led to sufficiently homogeneous alignment within the gel. Further CLIP-HSQC spectra were recorded to yield the couplings $^1T_{\text{CH}}$ (Figure 4-9).

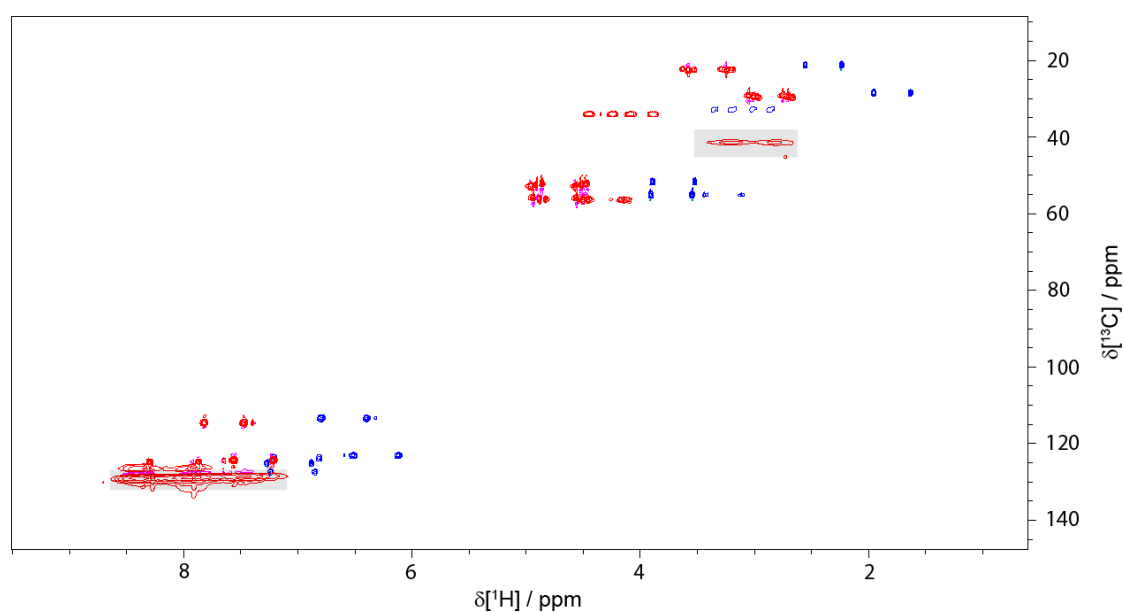


Figure 4-9: CLIP-HSQC spectra of $\beta\text{-aa-H}$. Blue: Spectrum in neat DCM (isotropic). Red: Spectrum in DCM/PBA (anisotropic, shifted in F2 dimension for clarity). Note that residual polymer signal is highlighted grey.

The difference between isotropically and anisotropically measured couplings led to 12 RDCs (Table 4-5).

Table 4-5: RDCs $^1D_{CH}$ in Hz for β -aa-H and β -aa-Me as determined from CLIP-HSQC experiments. Numbering is according to Scheme 4-2. Experimental errors were uniformly estimated as 0.5 Hz.

	$^1D_{CH}$ /Hz for β -aa-H	$^1D_{CH}$ /Hz for β -aa-Me
1-CH ₃	-6.5	-8.9
C2-H2	16.1	1.1
C3-H3	4.5	13.5
C3-H3 ^o	10.2	
3-CH3		-8.8
C4-H4	12.7	-4.1
5-CH3	6.6	10.2
C6-H6	-11.4	12.0
C7-H7 ^a	11.9	-0.5
C3'-H3' ^b	-18.7	-14.5
C4'-H4' ^b	-21.7	-16.6
OCH ₃ ^b	7.3	11.5
COOCH ₃ ^b	7.3	-2.8

^aBonds C7-H7 and C4-H4 are nearly parallel and should therefore exhibit very similar RDCs. Only C4-H4 was further analyzed due to the sharper H4 resonance (no *ortho* couplings).

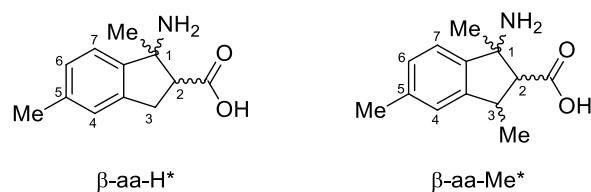
^bNon-rigid RDCs from the PMP subunit were not included in the further analysis.

Further details on sample preparation and CLIP-HSQC parameters are described in the experimental sections 5.1 and 5.2. It is most striking, that the RDCs differ quite strongly e.g. 16.1 vs. 1.1 Hz for C2-H2, 12.7 vs -4.1 Hz for C4-H4 and -11.4 vs. 12.0 Hz for C6-H6. The sole substitution of a single hydrogen in β -aa-H by a methyl group *per se* should not induce this immense change in alignment that becomes obvious through these values. Instead, a more profound change, e.g. in conformation, must have occurred upon this substitution. This is additional knowledge which can be exploited as soon as the conformation of one of either structure is solved by a more detailed RDC analysis.

ii) Generation of a comprehensive set of possible molecular structures

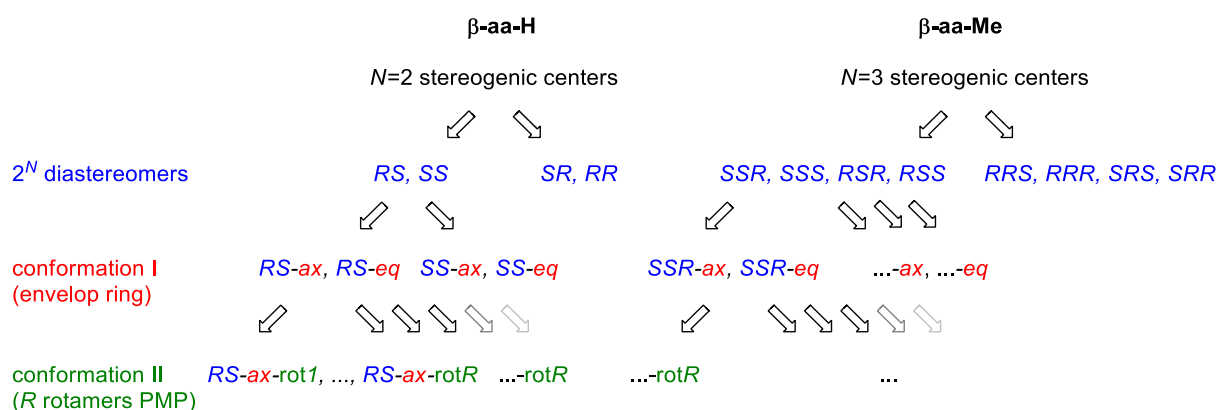
The molecules β -aa-H and β -aa-Me were synthesized as racemic mixtures. The configuration of β -aa-H may be described as follows: C1 is *R* and C2 is *S* configured and the molecule is simply called *RS*. Note that in the following discussion the nomenclature of possible configurations will be analogous to this example. The enantiomer of this molecule is *SR* (ester and amine are *trans* configured in both cases). Additionally, one can propose the configurations *RR* and *SS* (ester and amine are *cis* configured). Since achiral media like PS do not discriminate between enantiomers^[150] one may consider one arbitrary configuration out of these pairings (i.e. *RS* and *SS*). Analogously, the possible configurations for β -aa-Me are *SSR* (C1: *S*, C2: *S*, C3: *R*), *SSS*, *RSR* and *RSS* (implying *RRS*, *RRR*, *SRS* and *SRR*, respectively). Since the configuration of the substituents at the indane backbone is of principal interest, one may focus

on rigid, neutral cut down structures β -aa-H* and β -aa-Me* containing a primary amine and a carboxylic acid substituent.



Scheme 4-3: Cut-down structures of β -aa-H* and β -aa-Me*.

Hence, only the RDCs in the top section of Table 4-5 were exploited. Inclusion of the PMP and ester substituents, which are linked to the indane backbone by several rotatable bonds, would add less essential information (four RDCs that further define the alignment) at the cost of a highly demanding conformational analysis. The indane scaffold of β -aa-H* and β -aa-Me* exhibits two major conformations: the two envelope conformations in which the ester/carboxylic group and its two neighbored *trans* configured substituents change from pseudo-equatorial (“-eq” is added to the stereodescriptors) to pseudo-axial (“-ax”), Scheme 4-4 summarizes the process of preparing the exhaustive blueprints for possible molecular structures of β -aa-H and β -aa-Me.



Scheme 4-4: Illustrative branching of molecules β -aa-H and β -aa-Me into possible diastereomers through different configurations (only one representative enantiomeric set can be examined since achiral PS does not allow distinction) and conformers. The latter rapidly increases the number of possible structures and should therefore be narrowed to a reasonable conformational space e.g. through cut down structures.

In general, the number of possible structures rapidly increases by a too extensive examination of the configurational and conformational space. This highlights the significance of a smart inspection and choice of structural proposals.

After generation of a comprehensive set of possible molecules with likely combinations of configurations and conformations, (non-optimized) 3D structures were created for both compounds from scratch with Chem3D Pro 14.0. This set of structures was optimized by MM2^[151] as well as MMFF94^[152] force fields. This computational level is the common demand on large biomolecules and usually yields sufficient results.^[65c, 65d, 72a, 74a, 153] As required for small molecules, the MM2 optimized set of structures was further geometry optimized on a B3LYP/def2-TZVP level of theory with ORCA 3.0.3^[154]. Stationary points were characterized as minima by computing numerical frequencies. Further

details on the computational methods are described in the experimental section 5.4. The next three steps are usually wrapped up as one routine in RDC analysis software such as MSpin 2.0 Beta^[83] which was also applied in this study.

- iii) **Determination of alignment tensors from RDCs and structural proposals**
- iv) **Back-calculation of RDCs from structural proposals and alignment tensors**
- v) **Assessment of the validity of the structural proposal**

Although MSpin is stable and comprehensive with continuous technical support, it is important to understand these steps and avoid merely applying a black-box. Hence the basics of steps iii) to v) are briefly illustrated: The alignment tensor is derived in a very similar fashion to the procedure described in section 4.1. For an arbitrary molecular alignment of a C_1 symmetric molecule, linear equation systems of (1-7) and (1-4) with five unknown alignment tensor elements have to be solved. Table 4-5 indicates that these linear equation systems are once more overdetermined (seven independent RDCs, five unknown tensor elements). In the case of RDC analysis *Losconzi et al.* established a method called singular value decomposition^a (SVD) to obtain stable solutions for these problems.^[155]

Once the alignment tensors are derived for each structural proposal, equation (1-7) can be straightforwardly applied to back-calculate RDCs for these models. The *correct* structural proposal will then return the best compliance between the experimental RDCs D_{exp} and back-calculated RDCs D_{calc} . This is quantified for all i RDCs through the *Cornilsecu* Q factor.^[156]

$$Q = \sqrt{\frac{\sum (D_{\text{calc},i} - D_{\text{exp},i})^2}{\sum D_{\text{exp},i}}} \quad (4-4)$$

Hence, the smallest Q factor indicates the most suitable structural proposal. These steps were first performed via MSpin for β -aa-H*. Note that for β -aa-H* it is not possible to *a priori* assign the prochiral hydrogen atoms at C3. Instead, a full analysis was performed in which the involved couplings are interchanged. An additional suffix “(i)” signals that the coupling C3-H3 (4.5 Hz) is assigned to the pro-S hydrogen and C3-H3^o (10.2 Hz) to the pro-R hydrogen (for “(ii)” *vice versa*). This results in eight structural proposals from a permutation of two configurations, two conformations and two possible assignments for β -aa-H* (for a detailed exemplary MSpin input see the Appendix). Figure 4-10 shows the resulting Q factors for this set.

^a Iterative least-square methods as applied in section 4.1 are more suitable if previous knowledge of the alignment tensor is available. This is usually the case for symmetric molecules.

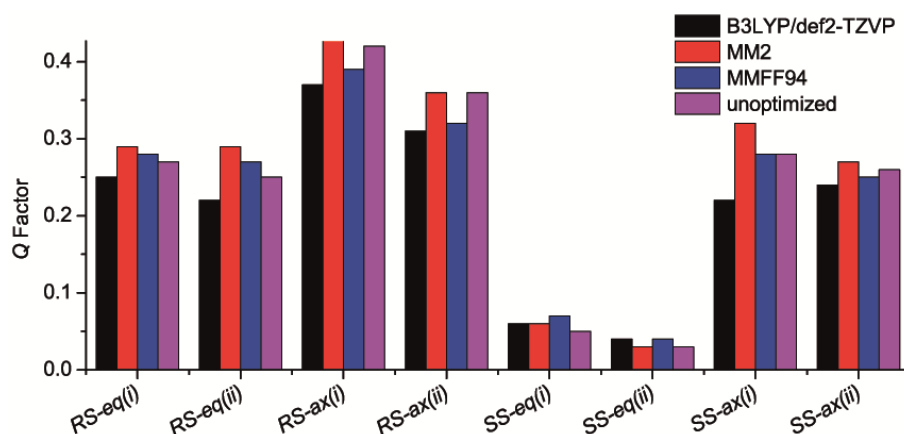
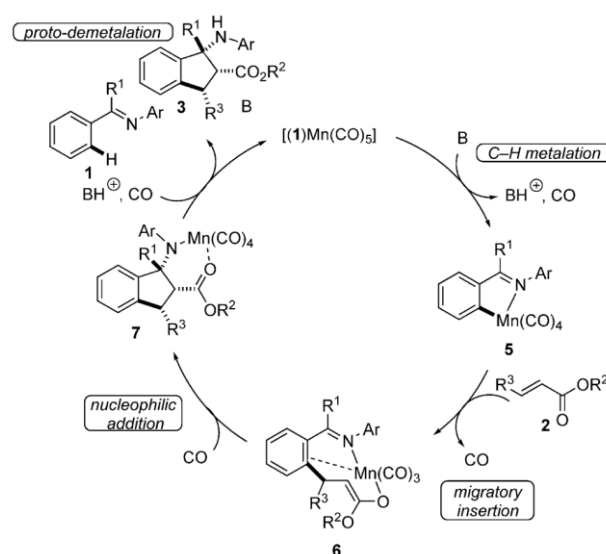


Figure 4-10: Q factors calculated using seven RDCs (Table 4-5) for the set of configurations/conformations/assignments of β -aa- H^* . Each structure was optimized by different theoretical methods (B3LYP/ def2-TZVP, MM2, MMFF94, unoptimized textbook angles). The suffix “(i)” signals that the coupling C3-H3 (4.5 Hz) is assigned to the pro-S hydrogen and C3-H3 $^\circ$ (10.2 Hz) to the pro-R hydrogen (for “(ii)” *vice versa*).

According to Figure 4-10, Q factors for *SS-eq(i)* and *SS-eq(ii)* (~ 0.05) are more than four times smaller than corresponding values for other feasible structures (~ 0.22). Therefore, the substituents in β -aa- H^* are *cis* configured. This is in accordance with the proposed reaction mechanism which suggests chelation of the manganese catalyst by the ester and amine moieties.^[148]



Scheme 4-5: Proposed mechanism for the manganese-catalyzed synthesis of *cis*- β -amino acid esters through organometallic C–H activation of ketimines. Reprint with permission from reference [148]. Copyright 2016 Wiley VCH.

Although the barrier for a ring flip is usually low for indanes, the pseudo-*equatorial* orientation of the ester within the five-membered ring is preferred.^[157] A more advanced fitting of an equilibrium between *RS-eq/RS-ax* and *SS-eq/SS-ax* (DFT structures) with a single tensor approach,^a as implemented in

^a Since alignment tensors for two different conformers are usually different from each other, the analytic solution depends on 1 population parameter and 2×5 alignment tensor elements. The number of RDCs available is less

MSpin, only slightly improved the Q factor at population ratios clearly favouring the equatorial conformer.

Table 4-6: Population fitting of two possible conformers for each diastereomer of β -aa-H* (DFT structures). The single tensor approach as implemented in MSpin 2.0 Beta was applied for this procedure. Note that the reported values for the single conformers result from the adapted alignment tensor and thereby differ from the values reported above.

	<i>Eq</i> -conformer	<i>Ax</i> -conformer	Population	Equilibrium
	Q Factor	Q Factor	(- <i>eq</i> : - <i>ax</i>)	Q Factor
<i>RS(i)</i>	0.22	0.88	1 : 0	0.22
<i>RS(ii)</i>	0.22	0.74	1 : 0	0.22
<i>SS(i)</i>	0.05	0.71	1 : 0	0.05
<i>SS(ii)</i>	0.05	0.54	0.94 : 0.06	0.04

It is striking that DFT calculations for the *SS* configuration gave only a minor energy difference of -0.8 kJ/mol favouring the pseudo-*equatorial* over the *axial* conformation. This discrepancy may stem from the simplified single tensor approach which underestimates the *axial* population. In general, the treatment of highly flexible molecules requires more delicate approaches, yet the indanes are rather rigid.^[158] Hence, it seems more likely that approximations which enable reasonably expensive DFT calculations led to this discrepancy between experiment and theoretical prediction. For instance, the incorporation of the *oPMP* unit in DFT calculations will on one hand change the energy of the system through intramolecular interactions. On the other hand, the *oPMP* unit as well as the aromatic indane scaffold may both (very weakly) interact with the PS strands. This effect is strongly suggested by the results for attractive interaromatic forces between small aromatic molecules and PS as shown in section 4.1. In the case of β -aa-H, a combination of interaromatic forces and generalized repulsive/steric forces certainly adds energy terms which change (supposedly increase) the gap between the pseudo-*equatorial* and the *axial* conformation. This has a delicate consequence on structure elucidation via RDCs: Interaction between alignment medium and aligned molecule may, contrary to results from former studies,^[159] bias the conformation of the observed molecule (configurations are still doubtlessly derived). The conformation of a molecule plays an important role for its reactivity, stereoselectivity etc. Hence, wrong assumptions from biased experimental data must be avoided by supportive high-level theoretical methods.

Although modern RDC studies on small molecules usually employ at least DFT level structure optimization, remarkably accurate results can also be obtained at very low computational cost. As shown by Figure 4-10 MD simulations and most basic textbook distances and angles (unoptimized) can be employed to obtain the correct result. Even the minute distinctions between *SS-eq(i)* and *SS-eq(ii)*

than this. Therefore, approximations are made in which atoms of two conformations are superimposed according to least-squares methods. This yields a shared coordinate system with a single alignment tensor. An equation in the form of population weighted sums of (1-7) is then solved. Detailed procedures are described in the MSpin reference manual.

are likewise reproduced by all four methods and allow a conclusive assignment of the prochiral hydrogen atoms at C3 in β -aa-H*. Although only marginally better (smaller) Q values were found for DFT structures of *SS-eq(ii)* over *SS-eq(i)*, it is safe to assume that *SS-eq(ii)* is in fact the true assignment. With the knowledge that ester and amine groups are *cis* configured and the ester is predominantly in a pseudo-equatorial position the correct assignment *SS-eq(ii)* can be confirmed by the NOE contacts. E.g. the integral for the cross-peak between H4 (aromatic ring plane) and pseudo-axial H3^{pro-R} (-0.44) is much smaller than the integral for the cross-peak between H4 and pseudo-equatorial H3^{pro-S} (-0.75) indicating closer contacts for the latter pair of protons. Further experimental information and data on the supplementary NOESY experiment is given in the Appendix B (Extended NMR Spectroscopy of Synthetic Indanes).

Successful structure verification of β -aa-H* facilitates the subsequent analysis of β -aa-Me*. Since both compounds result from analogous syntheses, the *cis* configuration between ester and amine substituents will not be altered.^a Consequently, four initially possible configurations are reduced to two stereoisomers that differ in the relative configuration of the methyl group at C3: *SSR* and *SSS* (see Figure 4-11). Yet, the full set of permutations was examined via MSpin (4 configurations \times 2 conformations for β -aa-Me*).

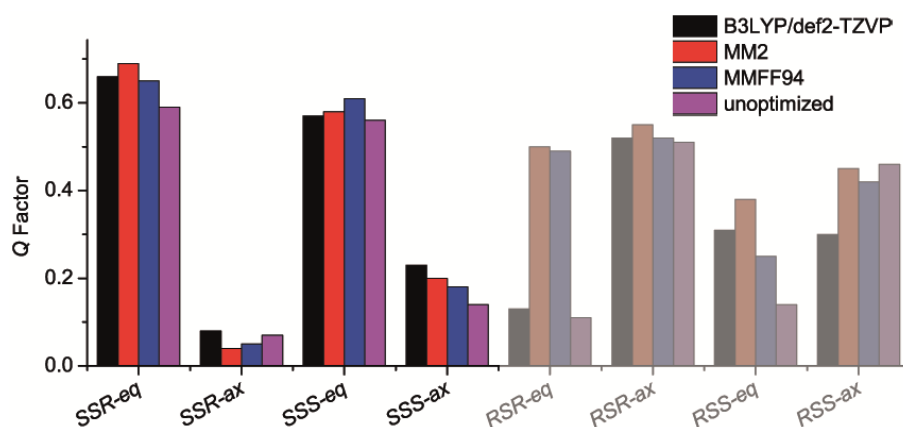


Figure 4-11: Q factors calculated using seven RDCs (Table 4-5) for the set of configurations/conformations of β -aa-Me*, each optimized by different theoretical methods. Four structures are greyed out which differ in the substituent at C1. These can be ignored due to results for the structurally analogous compound β -aa-H*.

Here, *SSR-ax* provides the smallest Q values regardless of the optimization method. Population fitting (DFT structures) gave some remarkable results for the full set.

^a Chelating of the manganese catalyst by the amine and carboxylic function as shown in Scheme 4-2 will naturally lead to *cis* configured ester and amine substituents.

Table 4-7: Population fitting of two possible conformers for each diastereomer of β -aa-Me* (DFT structures). The single tensor approach as implemented in MSpin 2.0 Beta was applied for this procedure. Note that the reported values for the single conformers result from the adapted alignment tensor and thereby differ from the values reported above.

	<i>Eq</i> -conformer	<i>Ax</i> -conformer	Population	Equilibrium
	Q Factor	Q Factor	(- <i>eq</i> : - <i>ax</i>)	Q Factor
SSR	1.09	0.07	0 : 1	0.07
SSS	1.02	0.23	0 : 1	0.23
RSR ^a	0.33	2.65	0.89 : 0.11	0.10
RSS ^a	0.83	1.27	0.60 : 0.40	0.07

^aStructures RSR and RSS can be omitted due to the results from β -aa-H.

The possible configurations SSR and SSS showed only pseudo-*axial* conformers and therefore no improvement of the Q factors. RSR and RSS can be excluded as possible configurations but exhibit remarkably good Q values. If no preceding analysis of β -aa-H had been made, these structures would impose some additional efforts on the structure validation of β -aa-Me.

The single conformer solution SSR-*ax* with no pseudo-*equatorial* contribution is also supported by DFT calculations which gave an energy difference of -9.3 kJ/mol in favour of the pseudo-*axial* over the *equatorial* conformer. Compared to the next best solution which is SSS-*ax*, Q factors found for SSR-*ax* are smaller by a factor of more than two. A direct comparison between measured and back-calculated RDCs for SSR-*ax* and SSS-*ax* is shown in Figure 4-12.

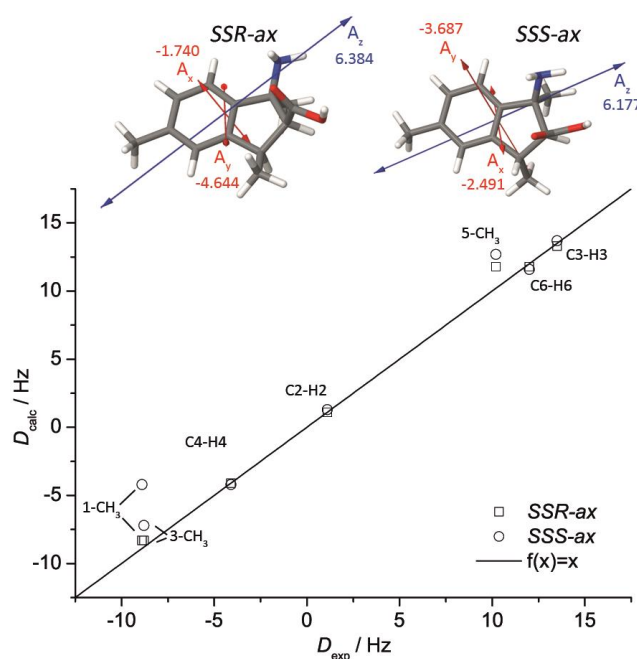


Figure 4-12: Top: B3LYP/def2-TZVP structures for SSR-*ax* and SSS-*ax* with vectors representing the principal axes of the alignment tensor (alignment tensor elements are given in units of 10⁻⁴). Bottom: Back-calculated (D_{calc}) vs. measured (D_{exp}) RDCs for SSR-*ax* and SSS-*ax*. Notably 1-CH₃ and 3-CH₃ exhibit stronger deviations between experimental and back-calculated RDCs.

Due to similar alignment tensors, nearly identical back-calculated RDCs result for C2-H2, C3-H3, C4-H4, C6-H6 and C7-H7. However, both isomers are clearly distinguished by the back-calculated RDCs for 1-CH₃, 3-CH₃ and 5-CH₃ which in *SSR-ax* agree much better with the experimental RDCs. These RDCs in particular affect the bonds (configurations) that are under discussion and hence, leave no doubt about the correct structure being *SSR-ax*.

The *axial* preference of the ester substituent in β -aa-Me (as opposed to an *equatorial* preference in β -aa-H) is highly remarkable. This is most likely due to steric repulsion between the large amine substituent at C1 and the methyl group at C3 (both *pseudo-axial*) in *SSR-eq*. This was already anticipated by the large deviations between the RDCs of both compounds (see Table 4-5). In summary, from the lowest absolute Q factors and the best agreement between measured and back-calculated RDCs for the crucial bonds it is possible to select *SSR-ax* as the likely correct structure of β -aa-Me. However, this conclusion is strongly facilitated by the previous analysis of β -aa-H which excluded the configurations *RSR* and *SSR*. In fact, DFT structures of *SSR-ax* and *RSR-eq* give very similar Q factors. Hence, a fast and reliable method to further exclude *RSR-eq* is desirable. Fortunately, results on different geometry optimization methods for *RSR-eq* already conveyed a strong dependence of the Q factor on (minor) structural changes. This led to the implementation of a further criterion on structure validity: Q factor stability against addition of structural noise.

4.2.1 Structural Error Analysis

MSpin allows an estimate of confidence for Q factors through RDC errors (uniformly set to 0.5 Hz for the study above as indicated in Table 4-5) that are statistically analyzed.^[83] Furthermore, within MSpin the stability (validity) of the SVD technique is indicated by so-called condition numbers (<30 is suggested to be reliable). Herein, *SSR-ax* and *RSR-eq* give similar standard deviations for Q factors (0.012 and 0.013, respectively) and inconspicuous SVD condition numbers (8.9 and 17.4, respectively). From these measures, the DFT-optimized structure *RSR-eq* is a valid structural proposal. However, results for different optimization methods of *RSR-eq* have shown that the low Q factor is not stable, but can rapidly increase upon minor structural changes.

In order to obtain an estimate of the reliability of the Q factors, structural noise was added to four B3LYP/def2-TZVP optimized structures that cover a range of Q factors: *SSR-ax* with Q = 0.08 (most probable), *RSR-eq* with Q = 0.13 (highly probable), *SSS-ax* with Q = 0.23 (fairly probable) and *SSR-eq* with Q = 0.66 (highly improbable). All these structures possess SVD conditions number <30 which is suggested to be reliable. By using a simple Python script an ensemble of 1000 molecular models was created for each structure in which the atomic coordinates of the substituents on C1, C2 and C3 were randomly varied within individual spheres of a maximum radius Δr_{\max} of 0.01 Å (the corresponding script can be found in the Appendix B).

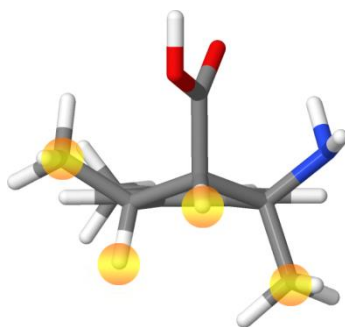


Figure 4-13: Structure *SSR-ax* optimized by MM2 force fields with illustrations of radially exaggerated spheres in which the corresponding substituent is randomly displaced.

Subsequently, this procedure was repeated for all structures with maximum random displacements of 0.02, 0.03 and 0.04 Å. Finally, Q values were determined for all these ensemble structures via MSpin. In order to quantify the stability of the Q factors, standard deviations σ were determined. The results of this procedure are summarized in Figure 4-14 a).

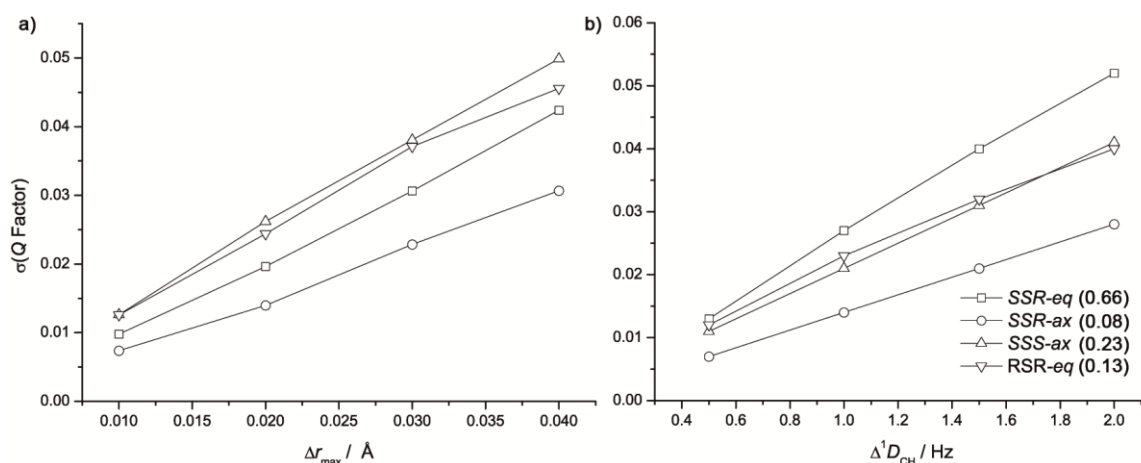


Figure 4-14: a) Standard deviations σ of the Q factors for the four exemplary structures *SSR-ax*, *RSR-eq*, *SSS-ax* and *SSR-eq* against the maximum random displacement Δr_{max} of the coordinates of the substituents on C1, C2, and C3. b) Standard deviation of the Q factors σ for the four exemplary structures against the maximum statistical error for the RDCs $\Delta^1 D_{\text{CH}}$ of the internuclear vectors attached to C1, C2, and C3.

Figure 4-14 a) shows that structure *SSR-ax* with the lowest Q factor is also least vulnerable against random displacement of significant atoms. Even though Q factors for the other structures are quite diverse, their standard deviations are agglomerated and about twice as high as for *SSR-ax*. The only other plausible structure which is *SSS-ax* (considering the results for β -aa-H and the analogous synthesis of β -aa-Me), actually showed the highest standard deviation. Hence, Q factor stability against addition of structural noise is indeed a valuable criterion in structure verification/falsification.

Note that besides this structural error analysis, an elaborate RDC error analysis may yield similar results. To evaluate this complementary approach the experimental errors of the RDC for β -aa-Me* $\Delta^1 D_{\text{CH}}$ was set to zero, except for the RDCs associated to C1, C2, C3 and their substituents which were set to 0.5 Hz. Next, the structures were analysed again with MSpin and standard deviations σ were directly obtained as MSpin output. This procedure was repeated for $\Delta^1 D_{\text{CH}}$ 1.0, 1.5 and 2.0 Hz. The results of this procedure are summarized in Figure 4-14 b). Obviously, this experimental error analysis

is also a sophisticated criterion in structure verification/falsification, as again the correct structure *SSR-ax* yields the least deviations compared to the other structures. However, the order of validity for a possible configuration differs between structural and experimental noises (validity structural noise: $SSR-ax > SSR-eq > RSR-eq \approx SSS-ax$, experimental noise: $SSR-ax > RSR-eq \approx SSS-ax > SSR-eq$). Therefore, both methods may be used complementary.

MSpin directly yields the standard deviation of Q factors for the experimental noise without further possibilities to analyse the resulting Q factor distributions. In contrast, the results obtained from random noise can be further examined by binning the values for the 1000 generated structures within steps of 0.005. Results are depicted in Figure 4-15.

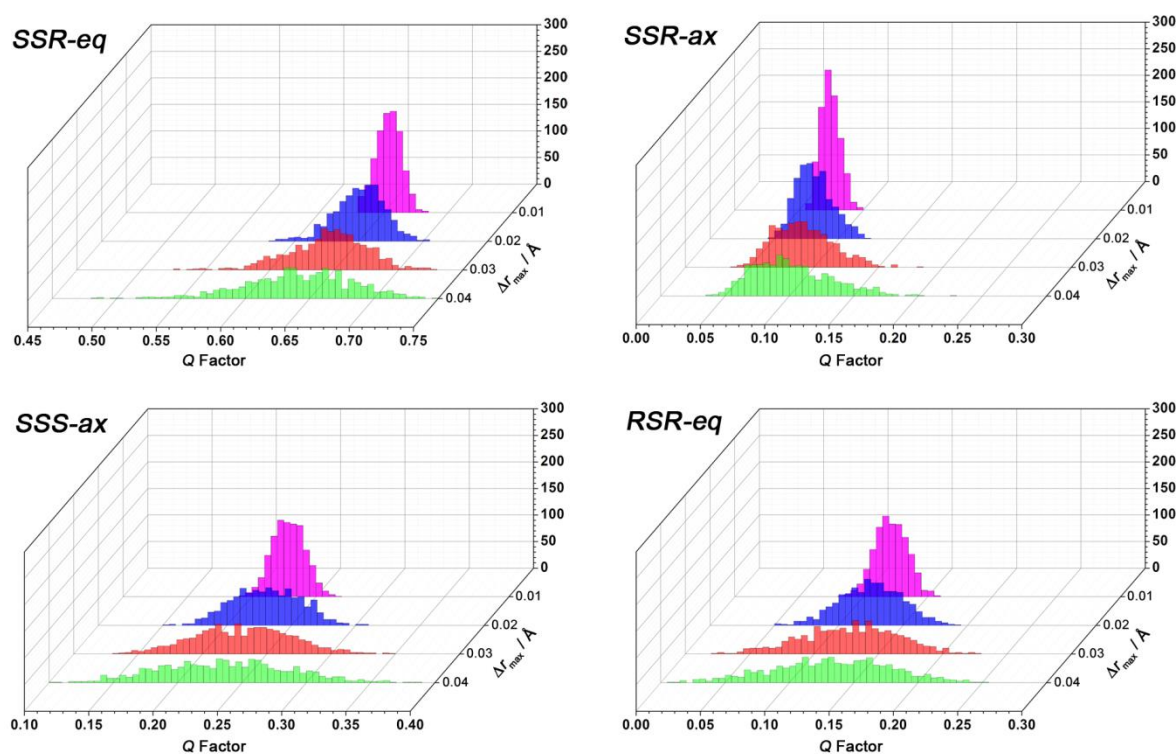


Figure 4-15: Distributions of Q factors for the four exemplary structures *SSR-ax*, *RSR-eq*, *SSS-ax* and *SSR-eq* resulting from the maximum random displacement Δr_{\max} of the coordinates of the substituents on C1, C2, and C3. Values are binned within steps of $\Delta Q = 0.005$.

The shapes of these distributions are most noticeable: the three incorrect structures exhibit typical *Gaussian*-like symmetric distributions centred at the Q value of the *parent* structure with equally broad tails towards low and high Q. For instance, *RSR-eq* gives astonishingly small Q factors. In contrast, the distribution for *SSR-ax* is generally sharper and has a higher skewness, in particular, it exhibits a steep decline towards the low Q value cut-off close to the *parent* Q. Hence, the optimized structure for *SSR-ax* represents the ensemble minimum and the best possible agreement with the experimental RDCs. These observations allow a qualitative discrimination. A simple quantitative discrimination is

achieved by fitting *Gaussian* probability functions and inspection of the goodness of these fits e.g. through the reduced chi-squared χ_v^2 (Figure 4-16).^a

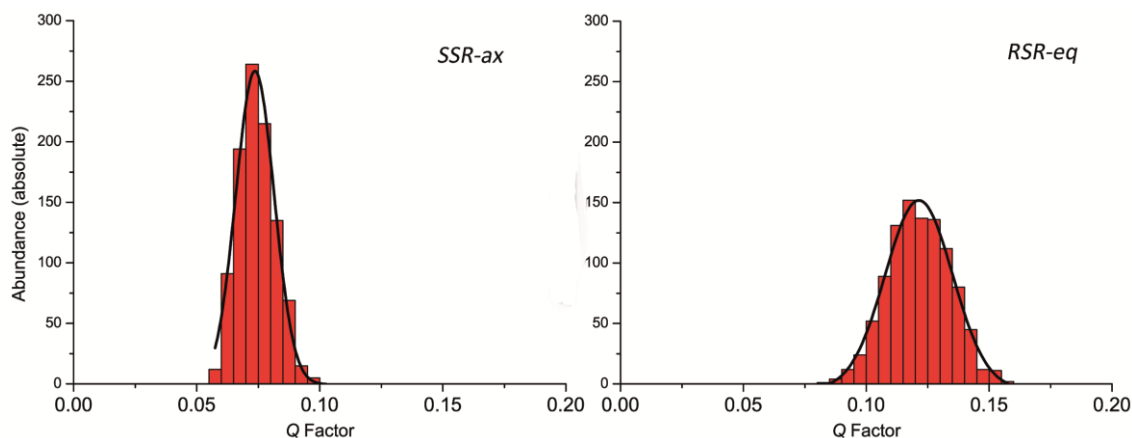


Figure 4-16: Distribution of Q factors and *Gaussian* fit for *SSR-ax* and *RSR-eq*, *SSR-ax* ($Q=0.08$, $\chi_v^2=156.8$) and *RSR-eq* ($Q = 0.13$, $\chi_v^2 = 31.8$) at a maximum random displacement $\Delta r_{\max}=0.01 \text{ \AA}$.

A summary for Q factors, standard deviations σ ($\Delta r_{\max}=0.01 \text{ \AA}$) and reduced chi-squared χ_v^2 ($\Delta r_{\max}=0.01 \text{ \AA}$) is given in Table 4-8.

Table 4-8: Q values, standard deviations $\sigma(Q)$ for $\Delta R_{\max}=0.01 \text{ \AA}$ and reduced chi-squared χ_v^2 for $\Delta R_{\max}=0.01 \text{ \AA}$ for structures *SSR-eq*, *SSR-ax*, *SSS-ax* and *RSR-eq*. The latter criterion shows the most distinguished differences between wrong and true structures.

	<i>SSR-eq</i>	<i>SSR-ax</i>	<i>SSS-ax</i>	<i>RSR-eq</i>
Q (MSpin)	0.66	0.08	0.23	0.13
$\sigma(Q)$ $\Delta r_{\max} = 0.01 \text{ \AA}$	0.020	0.014	0.026	0.024
$\chi_v^2(\text{Gauss})$ $\Delta r_{\max} = 0.01 \text{ \AA}$	37.8	156.8	21.9	31.8

This additional criterion obviously enables a very pronounced discrimination compared to Q factors alone. The peculiar features of Q factor distributions upon addition of structural noise may be explained i) by a structure that gives a large or *accidentally* low Q factor both of which behave stochastically and give *Gaussian* (normal) distributions and ii) by the correct structure with a good Q factor that remains largely unchanged upon addition of noise (structural uncertainty is inherent in experimental RDCs due to vibrations etc.) or can only be biased to higher Q factors (resembles a Q factor minimum).

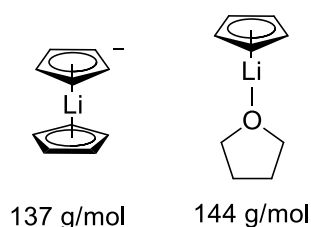
Thus, RDC structure elucidation may significantly benefit from this criterion as it may resemble an absolute criterion, i.e. no full set of structures has to be examined and compared, but as soon as the non-*Gaussian* solution is found the corresponding structure is verified. Further (empirical) studies are

^a A range of specifically designed methods exists to prove normal distributions of populations and base statistically valid descriptions on this (e.g. *Shapiro-Wilk* test^[160]). The application of these methods was beyond the scope of this work, instead the bell-curve fitting and goodness-of-fit approach was which is largely determined by the skewness of the distribution.

needed to put this hypothesis to the test and (possibly) show its limits according to the number of RDCs.

4.3 Quantitative Analysis of Residual Quadrupolar Couplings

Lithium organic compounds show interesting variations in coordination number, ligand identity and molecular symmetry as shown in section 1.1. The clarification of these structural features in solution enables the derivation of structure reactivity relationships which in turn allow the smart *design* of lithiation reactions. DOSY molecular weight determination is a useful method to distinguish between aggregates of different sizes/*MWs* (e.g. monomers, dimers, trimers or coordination by one, two, three solvent molecules).^[33-34, 36] Nevertheless, hypothetical aggregates may exhibit very similar *MWs* (Scheme 4-6) which cannot be unambiguously assigned to experimental findings from DOSY.



Scheme 4-6: Two hypothetical aggregates of CpLi in THF. Both molecules show molecular weights around 140 g/mol. These minute differences are not resolvable with the established MW determination by DOSY.

Traditionally, solid state NMR spectroscopy has been invested in the examination of quadrupolar couplings to obtain information about the EFG and thereby clarify structures (also see subsection 1.3.2).^[61a, 161] For example, the aggregates in Scheme 4-6 will exhibit considerably different EFGs around the lithium nuclei (highly symmetric, D_{5d} sandwich complexes exhibit only minute charge imbalances around the central nucleus i.e. small EFGs^[162]). In solid state NMR, χ and η are obtained with ease from lineshape-fitting of the quadrupolar coupling as given by equation (1-11).^[85-86] Main group elements^[163] but also heavy transition metals^[164] represent popular objectives for these heteronuclear NMR experiments. Yet, in isotropic solution this valuable structural information gets lost.

Fortunately, NMR spectroscopy in uniaxial media enables an elegant recovery of quadrupolar couplings. These RQCs are straightforwardly correlated to the EFG (V_{ij}) by formula (1-13).

$$\nu_{\text{rq}} = \frac{3eQ}{2I(2I-1)h} \sum_{i,j=x,y,z} V_{ij}A_{ij} \quad (1-13)$$

Additionally, first principles prediction of EFG parameters by theoretical means can nowadays be achieved with great success.^[162b, 163a, 163e, 163f, 163h, 164d] Hence, within this work a new method for structure determination of lithium organic aggregates is outlined that involves six basic steps:

- i) Experimental determination of RDCs and RQCs
- ii) Generation of a comprehensive set of possible molecular structures
- iii) Determination of alignment tensors from RDCs and structural proposals

- iv) Determination of EFGs from *complementary techniques*
- v) Calculation of theoretical RQCs from alignment tensor and EFG
- vi) Assessment of the agreement between calculated and experimental RQCs

These steps are further explained in the following case study of lithium cyclopentadienide in THF.^[17d, 165]

Experimental determination of RDCs and RQCs

First, CpLi was dissolved in THF and a single isotropic $^1J_{\text{CH}}$ coupling was measured by CLIP-HSQC. Alignment of the compound was achieved by immersing a cylindrical PS stick (crosslinked through 0.4 vol% divinylbenzene) in a solution of CpLi in THF. Swelling/equilibration over eight days led to sufficiently homogeneous alignment within the gels. The magnitude of the RQC $|v_{\text{rq,exp}}|$ was determined from the chemical shift differences between the central transition and the satellite peaks from a ^7Li spectrum (Figure 4-17).

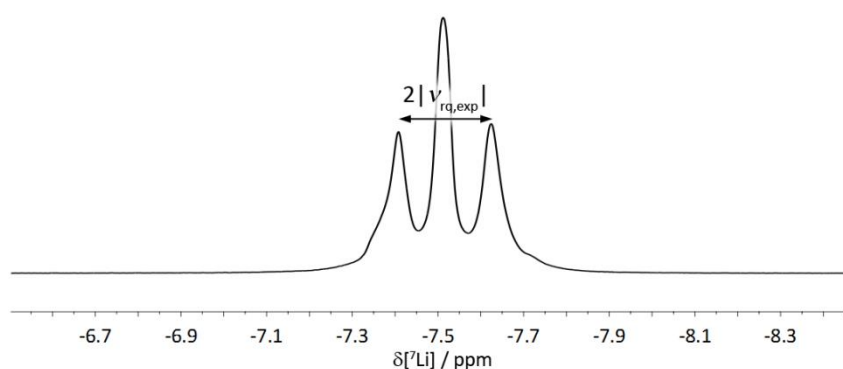


Figure 4-17: ^7Li spectrum of CpLi in THF/PS (0.4% DVB) eight days after sample preparation.

A CLIP-HSQC spectrum was recorded to yield the coupling $^1T_{\text{CH}}$. The difference between isotropically and anisotropically measured couplings led to the RDC listed in Table 4-5.

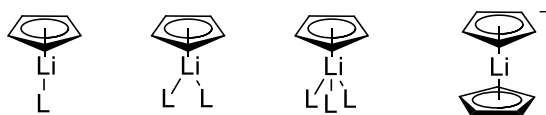
Table 4-9: RDC $^1D_{\text{CH}}$ in Hz for CpLi in THF/PS (4% DVB) as determined from CLIP-HSQC experiments. Experimental errors were determined as standard deviations in repeating manual peak picking attempts.

$ v_{\text{rq,exp}} $ / Hz	$^1D_{\text{CH}}$ / Hz
16.8 ± 1.2	11.4 ± 0.3

Details on sample preparation, ^7Li and CLIP-HSQC parameters are described in the experimental sections 5.1 and 5.2.

Generation of a comprehensive set of possible molecular structures

Common structural motifs in metal (M) cyclopentadienide chemistry as found from crystal structures are heteroleptic CpML_x (L = Ligand) complexes with $x = 1-3$ as well as homoleptic Cp_2M complexes with varying charges.^[17d, 165a-e] Further motifs such as solvent separated ion pairs may exist, but in the case of CpLi in THF the reduced set of structures is reasonable. Corresponding *Lewis* formulas are depicted in Scheme 4-7.

Scheme 4-7: Lewis formulas of possible ligand CpLi complexes and $[\text{Cp}_2\text{Li}]^-$.

Starting geometries were built from existing crystal structures of Cp derivatives. To simplify the geometry optimization of these complexes THF molecules were substituted by Me_2O . This structural cut-down decreases the conformational space while the crucial electronic surrounding of the lithium stays largely unchanged.

The set of structures was geometry optimized by DFT calculations at the B3LYP/def2-TZVP level of theory with ORCA 3.0.3.^[154] Stationary points were characterized as minima by computing numerical frequencies.

i) Determination of alignment tensors from RDCs and structural proposals

The rapid rotation of the Cp ring moiety yields a single RDC. Similar to a methyl group this averaged RDC $\langle D \rangle$ (attributed to five C–H bonds) can be converted into an RDC D_{\parallel} of a virtual bond pointing along the fivefold-symmetry axis of the Cp ring by applying formula (4-1) and an idealized angle of 90° . This is equivalent to a scaling factor of -2 to obtain the virtual RDC. Since only a single RDC can be measured, the alignment tensor must be treated as axially symmetric. According to Table 4-2 this simplification is valid for the symmetries of $[\text{CpLi}(\text{OMe}_2)]$, $[\text{CpLi}(\text{OMe}_2)_3]$ and $[\text{Cp}_2\text{Li}]^-$ but requires a more careful analysis of the quasi C_{2v} symmetric $[\text{CpLi}(\text{OMe}_2)_2]$ as will be pointed out when appropriate in the following analysis.^a

Similar to the alignment tensor determination in section 4.1, the fivefold axis of Cp, i.e. the virtual C–H vector (22.8 Hz), is set in a coordinate system pointing along the z axis. This allows straightforward calculation of values for the alignment tensor as summarized in Table 4-10.

Table 4-10: Alignment tensor elements $A_{\bar{x}}$, $A_{\bar{y}}$ and $A_{\bar{z}}$ for CpLi in swollen THF/PS (0.4% DVB) as determined from its RDC. Additionally, the magnitude of these values and the orientation within the molecular frame are represented by arrows with corresponding length, their sign is represented by their color (blue: positive, red: negative).

$A_{\bar{x}}/10^{-4}$	-1.6	
$A_{\bar{y}}/10^{-4}$	-1.6	
$A_{\bar{z}}/10^{-4}$	3.2	

Note that the alignment tensor elements are again ordered by increasing magnitude according to conventions. A graphical representation of \mathbf{A} in its PAS is given to illustrate the alignment. The

^a An axially symmetric alignment tensor may be valid for $[\text{CpLi}(\text{OMe}_2)_2]$ if dissociation and coordination of L (THF) occurs faster than the molecular rotations.

orientation of the magnetic field along the Li–Cp bond is in accordance with steric considerations (contrary to the small aromatic molecules discussed in section 4.1).^a

Determination of EFGs from complementary techniques

Absolute values of the EFG can be determined from solid state NMR spectra of quadrupolar nuclei. Subsequent molecular structure determination from X-ray crystallography enables the correlation of EFG values with structural features. In 1998 *Johnels, Boman and Edlund* conducted a study in which several cyclopentadienide lithium complexes were studied by ⁷Li solid state NMR.^[165b] Table 4-11 summarizes some of their results:

Table 4-11: CpLi derivatives, added ligand and number of coordinating atoms in parentheses, quadrupolar coupling constants χ and derived $|V_z|$ and asymmetry parameter η . Note that the gradient elements are reported in units of a.u. (1 a.u. = $9.71736 \cdot 10^{21}$ Vm⁻²). Values taken from reference [165b].

Cp Derivative	Ligand	χ /kHz	$ V_z $ /a.u.	η
Tris-TMS-CpLi	THF (1O)	125	0.0133	0.1
Tris-TMS-CpLi	Diglyme (2O)	94	0.0100	0.1
CpLi	12-crown-4 (4O)	21	0.0022	0.5
CpLi	PPh ₄ ⁺ ^a	45	0.0048	0.3
CpLi	Diglyme ^a	26	0.0028	0.3

^a In these cases sandwich complexes are formed.

Most obviously, χ decreases as the number of oxygen atoms around lithium increases. Also, in the case of sandwich complexes only weak coupling constants are measured (the authors additionally report a “very low” χ for polymeric CpLi strands). In general, sandwich complexes give only very small quadrupolar coupling constants as the electric field is assumed to be mostly isotropic in complexes with increased hapticity.^[162] The EFGs are largely axially symmetric ($\eta \approx 0$) as can already be anticipated from molecular geometry (slightly higher values for η are again usually correlated with lower χ which indicates increasing spherically symmetric charge distributions). Comparison with crystal structures led to the conclusion that quadrupolar coupling constants “depend on the number and type of donor atoms” rather than on distances and angles.^[165b]

This final notion is particularly important as it signals the possibility to transfer results to further Cp derivatives and different ligands although these are encompassed by (small) structural changes. Hence, the examined system CpLi in THF and the associated structural proposals may show slightly different angles and distances but can nevertheless be directly correlated to the values given in Table 4-11. This correlation of the structural proposals to EFGs with idealized axial symmetry is summarized in Table 4-12.

^a Note, that a more detailed MD simulation of the alignment of possible aggregates (with different numbers of L) would already allow their discrimination by subsequent comparison with experimental alignment. This approach was not further followed here, since, so far, the results in subsection 4.1.1 were not quantitatively matching the experimental benchmark test studies and lithium organics in general have not been very well adapted to molecular dynamics simulations.

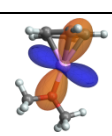
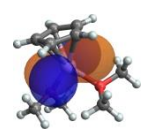
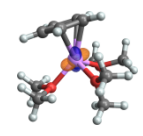
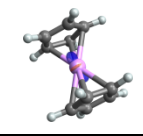
Table 4-12: Estimated EFG element V_z and asymmetry parameter η for $[\text{CpLi}(\text{OMe}_2)]$, $[\text{CpLi}(\text{OMe}_2)_2]$, $[\text{CpLi}(\text{OMe}_2)_4]$ and $[\text{Cp}_2\text{Li}]^-$.

Cp Derivative	$ V_z $ /a.u.	η
$[\text{CpLi}(\text{OMe}_2)]$	0.013	0
$[\text{CpLi}(\text{OMe}_2)_2]$	0.010	0
$[\text{CpLi}(\text{OMe}_2)_4]$	0.002	0
$[\text{Cp}_2\text{Li}]^-$	0.004	0

This list lacks the value for a trifold Me_2O coordination, but it seems reasonable that $[\text{CpLi}(\text{OMe}_2)_3]$ exhibits a quadrupolar coupling constant ranging between $[\text{CpLi}(\text{OMe}_2)_2]$ and $[\text{CpLi}(\text{OMe}_2)_4]$. This dependence on the availability of solid state EFG data is an obvious drawback. Furthermore, information about the orientation of the EFG within the molecular framework cannot be provided by solid state NMR. Yet, this knowledge can be obtained by first principle DFT calculations of the EFG tensor.

ORCA^[154] (for gas phase calculations) or CASTEP^[166] (for solid state calculations) enable a straightforward determination of NMR properties. The B3LYP density functional and *Gaussian* type basis sets such as 6-311++G** have proven to yield reliable results for the calculation of the EFG of aluminium compounds. In this work the theoretical quadrupolar coupling χ of trimesityl aluminum ($\text{Al}(\text{Mes})_3$) was reproduced utilizing ORCA 3.0.3^[154] and the *fgrad* keyword (for EFG calculation) at the B3LYP/6-311++G** level^[167] of theory. This procedure was then repeated for the structural proposals $[\text{CpLi}(\text{OMe}_2)]$, $[\text{CpLi}(\text{OMe}_2)_2]$, $[\text{CpLi}(\text{OMe}_2)_3]$ and $[\text{Cp}_2\text{Li}]^-$. Resulting values are summarized in Table 4-13. Additionally, the table includes graphical tensor representations that illustrate the EFG orientation and asymmetry.

Table 4-13: Theoretical values for $|V_z|$ and η and their graphical representation by spherical polar plots^[162b] for $\text{Al}(\text{Mes})_3$, $[\text{CpLi}(\text{OMe}_2)]$, $[\text{CpLi}(\text{OMe}_2)_2]$, $[\text{CpLi}(\text{OMe}_2)_3]$ and $[\text{Cp}_2\text{Li}]^-$.

	$ V_z $	η	
$[\text{CpLi}(\text{OMe}_2)]$	0.015	0.5	
$[\text{CpLi}(\text{OMe}_2)_2]$	0.018	0.9	
$[\text{CpLi}(\text{OMe}_2)_3]$	0.005	0.4	
$[\text{Cp}_2\text{Li}]^-$	0.005	0.2	

The computational values $|V_z|$ for $[\text{CpLi}(\text{OMe}_2)]$ and $[\text{Cp}_2\text{Li}]^-$ agree well with experimental data provided by *Edlund et al.* but the asymmetry is larger than anticipated. In both cases the axial

component is largely oriented along the Li–Cp bond. [CpLi(OMe₂)₃] gives a reasonable, i.e. small, $|V_z|$ similar to [Cp₂Li][−]. This is in accordance with the trend provided by experimental data in Table 4-12 and, therefore, this value can, with reservations, be further evaluated. This is different for [CpLi(OMe₂)₂] which additionally exhibits EFG tensor elements that are strongly tilted away from the Li–Cp bond. The calculated values for [CpLi(OMe₂)₂] largely exceed the experimental $|V_z|$ and η . This may be explained from the effects of crystal packing and therefore one should very carefully interpret proceed with the following discussions related to this case..

Calculation of theoretical RQCs from alignment tensor and EFG and assessment of the agreement between calculated and experimental RQCs

Three major simplifications are made to obtain preliminary results:

- i) Principal axes belonging to the axial components of **A** and **V** (A_z and V_z) are collinear and lie within the Li–Cp bond axis.
- ii) **V** is axially symmetric
- iii) **A** is axially symmetric

According to these assumptions formula (1-13) can be simplified to equation (4-5) for ⁷Li RQCs.

$$\nu_{\text{rq}} = \frac{3eQ}{4h} A_z V_z \quad (4-5)$$

The resulting theoretical RQCs $|\nu_{\text{rq,calc}}|$ for the hypothetical CpLi structures are summarized in Table 4-14.

Table 4-14: Alignment tensor element A_z , EFG tensor element $|V_z|$ and theoretical quadrupolar couplings $|\nu_{\text{rq,calc}}|$ for [CpLi(OMe₂)], [CpLi(OMe₂)₂] [CpLi(OMe₂)₃] and [Cp₂Li][−] and experimental value $|\nu_{\text{rq,exp}}|$.

Cp Derivative	$A_z / 10^{-4}$	$ V_z / \text{a.u.}$	$ \nu_{\text{rq,calc}} / \text{Hz}$	$ \nu_{\text{rq,exp}} / \text{Hz}$
[CpLi(OMe ₂)]		0.013	29.4	
[CpLi(OMe ₂) ₂]	3.2	0.010	22.6	
[CpLi(OMe ₂) ₃] ^a		0.005	11.3	16.8
[Cp ₂ Li] [−]		0.004	9.00	

^a Value for $|V_z|$ is taken from DFT calculations.

The difference between values $|\nu_{\text{rq,calc}}|$ and $|\nu_{\text{rq,exp}}|$ for [CpLi(OMe₂)], [CpLi(OMe₂)₂] [CpLi(OMe₂)₃] and [Cp₂Li][−] is additionally illustrated in Figure 4-18.

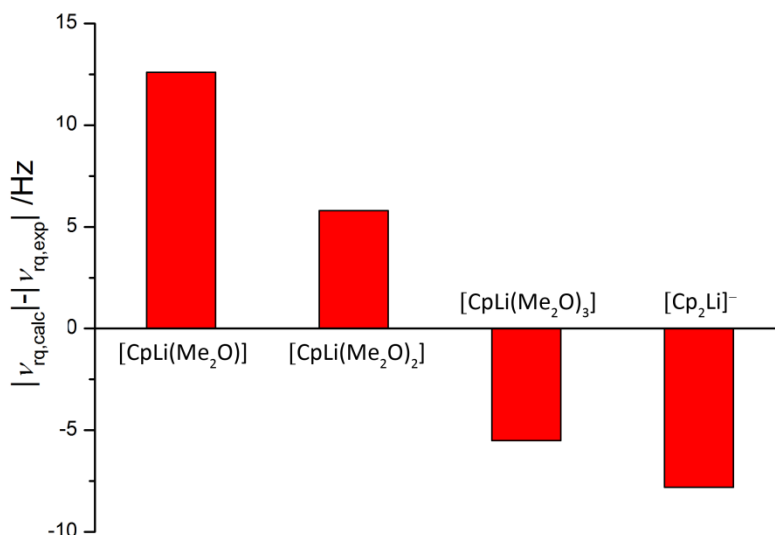


Figure 4-18: Difference between values $|v_{rq,calc}|$ and $|v_{rq,exp}|$ for [CpLi(OMe₂)], [CpLi(OMe₂)₂], [CpLi(OMe₂)₃] and [Cp₂Li]⁻.

It must be emphasized that differentiation between structures must, so far, be based on matching absolute values, since only the magnitude of the experimental RQC can be measured. This has to be addressed in future attempts to develop the RQC structure verification method. However, a quantitative match is a strong indicator and supports complementary methods such as DOSY *MW* determination. Indeed, [CpLi(THF)₂] has been confirmed to be the predominant species by DOSY.^[165a] This result can be confirmed in the following discussion:

[CpLi(OMe₂)] and [Cp₂Li]⁻ show large deviations from the experimental coupling. Additionally, the assumptions made about axial symmetry of **A** and **V** are most valid for both structures. Hence, they are far less probable than [CpLi(OMe₂)₂] or [CpLi(OMe₂)₃]. In turn, one cannot unambiguously differentiate between the latter proposals as both deviate by ~6 Hz from the experimental value. A more nuanced inspection of the initial approximations is therefore required for [CpLi(OMe₂)₂] and [CpLi(OMe₂)₃].

Especially, assumption i) (collinear tensor elements A_z and V_z) is insufficient since results from DFT calculations (Table 4-13) indicate an orthogonal orientation of the principal axes belonging to the axial components (Figure 4-19).

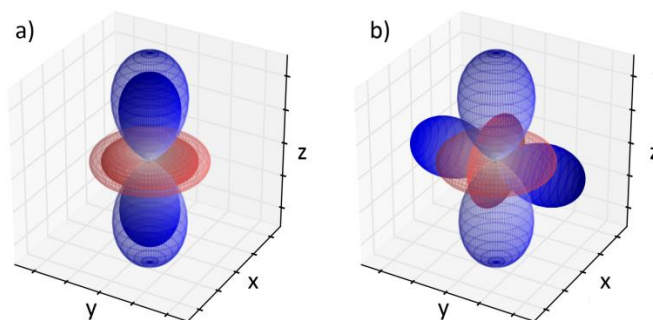


Figure 4-19: Two possible orientations of the alignment tensor (bigger, lighter) and the EFG tensor (smaller, darker) represented by spherical plots. a) Initial assumption that axial components are parallel to each other. b) Adapted assumption that axial components are orthogonal to each other.

According to the original formula (1-13), the tilt towards orthogonal orientation will result in smaller sums and the theoretical $|v_{\text{rq,cal}}|$ will decrease. This leads to a bigger disagreement between theory and experiment for a hypothetical aggregate $[\text{CpLi}(\text{OMe}_2)_3]$ but the hypothetical RQC for $[\text{CpLi}(\text{OMe}_2)_2]$ will approach the experimental value. Retaining assumptions ii) and iii) (axially symmetric tensors **A** and **V**), the RQCs can be calculated for the adapted EFG tensor orientation (Figure 4-19 b). The appropriate Cartesian tensor elements are summarized in Table 4-15.

Table 4-15: Adapted **V** and **A** values in a shared Cartesian coordinate system in which the Li–Cp bond is aligned along z and the O–Li–O bonds are set within the x-z plane.

i=j	$V_i [\text{CpLi}(\text{OMe}_2)_2] / \text{a.u.}^a$	$V_i [\text{CpLi}(\text{OMe}_2)_3] / \text{a.u.}^a$	$A_i / 10^{-4}$
x	-0.005	-0.0025	-1.6
y	0.010	0.0050	-1.6
z	-0.005	-0.0025	3.2

^a Note the distinction from the ordering (naming) convention of the EFG tensor.

In this coordinate system, the Li–Cp bond is aligned along z and the O–Li–O bonds are set within the x-z plane. This leads to adapted theoretical $|v_{\text{rq,cal}}|$ as presented in Table 4-16.

Table 4-16: Values $|v_{\text{rq,cal}}|$ as obtained from an orthogonal orientation of axially symmetric **A** and **V** for $[\text{CpLi}(\text{OMe}_2)_2]$ and $[\text{CpLi}(\text{OMe}_2)_3]$.

	$ v_{\text{rq,cal}} / \text{Hz}$	$ v_{\text{rq,exp}} / \text{Hz}$
$[\text{CpLi}(\text{OMe}_2)_2]$	11.3 Hz	16.8
$[\text{CpLi}(\text{OMe}_2)_3]$	5.6 Hz	

In accordance with the qualitatively anticipated decrease of theoretical $|v_{\text{rq,cal}}|$, the result for $[\text{CpLi}(\text{OMe}_2)_2]$ is closer to the experimental value, yet, it is not a particularly good match. Application of the EFG tensor values from DFT calculations (Table 4-13) does not improve the agreement ($|v_{\text{rq,cal}}| = 2.7 \text{ Hz}$). Simplification iii) can be further examined by decreasing the symmetry of **A**. In any case, A_z is constant ($3.2 \cdot 10^{-4}$), but A_x can range between 0 and $-A_z$ while A_y changes dependently according to $A_y = -A_z - A_x$ (Table 4-17).

Table 4-17: Adapted **V** and changing **A** values in a shared Cartesian system, in which Li–Cp bond is aligned along z and the O–Li–O bonds are set within the x-z plane.

i=j	$V_i [\text{CpLi}(\text{OMe}_2)_2] / \text{a.u.}^a$	$A_i / 10^{-4}$
x	-0.005	$0 \rightarrow -3.2$
y	0.010	$-3.2 \rightarrow 0$
z	-0.005	3.2

^a Note the distinction from the ordering (naming) convention of the EFG tensor.

This leads to new values $|v_{\text{rq,cal}}|$. The corresponding differences $\left||v_{\text{rq,cal}}|-|v_{\text{rq,exp}}|\right|$ for varying A_x is illustrated in Figure 4-20.

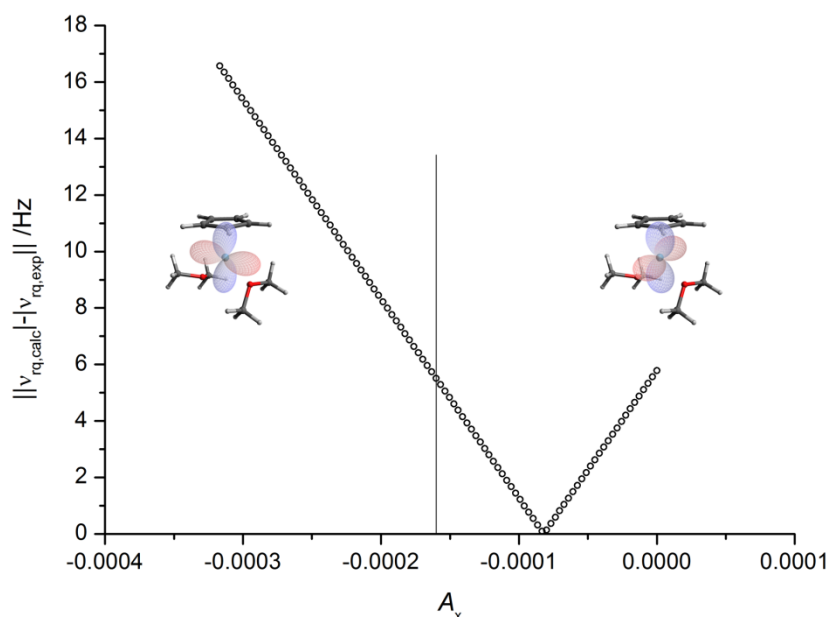


Figure 4-20: Resulting values $\left||v_{\text{rq,calc}}|-|v_{\text{rq,exp}}|\right|$ for $[\text{CpLi}(\text{OMe}_2)_2]$ and varying A_x . The vertical line marks the axially symmetric alignment tensor with $A_y = A_x = -1.6 \cdot 10^{-4}$. Additionally, the *boundary* alignment tensors for $A_x = -A_z$, $A_y = 0$ (left) and $A_x = 0$, $A_y = -A_z$ (right) are illustrated by polar plots.

Recall, that a more negative alignment tensor element indicates a less preferred orientation of the related axis along the magnetic field/elongated pore. According to the obvious spatial dimensions of the molecule, the axis that is perpendicular to the O–Li–O plane is least likely aligned with the magnetic field vector. This corresponds to the right side of Figure 4-20. Hence, under more realistic assumptions, the theoretical $|v_{\text{rq,cal}}|$ approaches the experimental value.

In summary, the results suggest that $[\text{CpLi}(\text{THF})_2]$ is the correct structure of CpLi in solutions of THF. Furthermore, Figure 4-18 shows, that RQCs can unambiguously resolve issues such as exemplified in Scheme 4-6: The two structures $[\text{CpLi}(\text{THF})]$ and $[\text{Cp}_2\text{Li}]^-$ have very similar weight and thus cannot be distinguished by DOSY but give very different RQCs which can be quantitatively matched to structural proposals. Moreover, a detailed estimation of the EFG by first principles enables a more detailed application of equation (1-13) taking account of the relative orientation of the tensors. As such, structure verification by RQCs is a valuable tool for the examination of molecules bearing quadrupolar nuclei.

Finally, note that according to formula (1-13) measurement of an arbitrary molecule in five different alignment media could provide five different alignment tensors which do not linearly depend on each other (this number is decreased if the position of the quadrupolar atom coincides with symmetry elements). In this case, a linear equation system would provide the full set of EFG tensor elements in the chosen coordinate frame (typically the alignment tensor PAS). The full EFG and information on the asymmetry parameter η would *per se* provide subtle structural information.

4.4 Summary and Outlook Anisotropic NMR Spectroscopy

The studies described above show that anisotropic NMR spectroscopy is an effective tool to

- observe and understand alignment phenomena in uniaxial environments,
- determine the configuration and conformation of small, rigid molecules,
- determine the aggregation state of organometallic compounds.

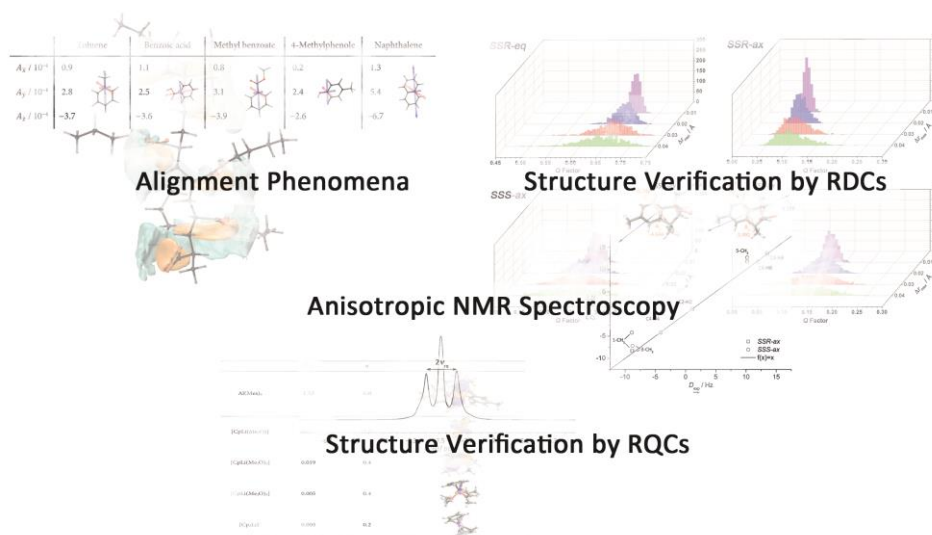


Figure 4-21: Anisotropic NMR spectroscopy and its applications: i) Alignment of small molecules in stretched polymer gels. ii) Structure determination of small rigid molecules by RDCs. iii) Structure determination of organometallic compounds by RQCs.

Alignment phenomena were observed by RDC measurements of small aromatic molecules in PS and PBA. These results showed surprising alignment phenomena which are characterized by a mismatched orientation of molecules considering their spatial extensions and the pore geometry. An *orthogonal* orientation (between *longest* pore and *longest* molecule axis) was found for toluene, 4-methylphenol and naphthalene in PS (Figure 4-22) and for benzoic acid and 4-methylphenol in PBA.

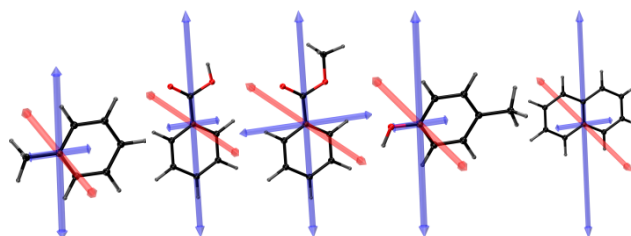


Figure 4-22: Alignment tensor representation of small aromatic molecules in PS. From left to right: toluene, benzoic acid, methyl benzoate, 4-methylphenol and naphthalene.

This provides evidence that the *steric* orientation (*parallel* orientation of longest axes) which is the basic assumption for a range of first principle alignment tensor prediction methods (Tramite^[77], PALES^[81]) must always be questioned for small molecules.

The alignment phenomena were rationalized by an interplay of repulsive and attractive interactions between the alignment medium and the aligned solute: The repulsive interaction induces an, intuitive, parallel orientation since the elongated polymer strains function as barriers. The attractive interaction can be further described as a rather unspecific interaromatic interaction in the case of PS and as a hydrogen bridge bonds in the case of PBA.

These alignment phenomena were qualitatively reproduced by MD simulations, which will become an increasingly important tool for alignment tensor prediction, considering the failure of traditional (sterical) prediction methods in cases like the ones described above.^[80, 132] In contrast to traditional prediction methods, MD simulations can additionally account for dynamic features such as conformational changes. Finally, the results on small molecule alignment showed that the solvent effects of anisotropic media cannot be neglected. Unlike in a former study, the observations clearly showed that anisotropic media can have profound impact on solutes.^[159] E.g. interactions leading to the orthogonal and parallel orientation may both be present within different rotatable moieties of a molecule and hence change the conformational landscape of the observed substance within uniaxial environments. This has to be considered in studies concerning small molecules in polymers and may again be best addressed through comparative MD simulations (alignment) as well as high level theoretical calculations (population of conformers).

Structure verification by RDCs has been successfully applied to two synthetic β -amino acid esters both resulting from the same catalytic C–H bond activation cycle. In accordance with the proposed reaction mechanism, the *cis* configuration of the amine and ester were unambiguously determined as the true structure. Furthermore, RDC-based conformation analysis showed mono-populated conformations with a switch from the pseudo-*equatorially* oriented ester to a pseudo-*axial* conformation upon the substitution of a hydrogen atom by a methyl group at the C3 ring position. This change is remarkable and may play a vital role in the discussion of these amino acid esters as β -peptide building blocks with varying secondary structures.

Furthermore, the traditional Q factor for the assessment of structural proposal was shown to be insufficient in special cases, but the addition of structural noise emerged as a useful further criterion that complements the traditional approach: The addition of structural noise to critical C–H and C–Me bonds generated a thermally accessible ensemble of molecules and enables the generation of Q factor distributions. The shape of the distribution for the true structure (asymmetric, steep) differs from the distributions of wrong structures (symmetric, broad). In the case of only few RDCs, it is possible that wrong structural proposals *accidentally* give a sufficiently good Q factor. These wrong structures behave stochastically and hence are well described by *Gaussian* distributions. In contrast, true structures show a Q factor minimum i.e. an asymmetric, narrower distribution. The validation process by structural noise can be easily implemented and adds a straightforward and very reliable figure to amend the traditional Q factor, which will become especially useful in automated structure elucidation by algorithms.^[168]

Structure verification by RQCs was finally presented as a novel approach to derive structural information on lithium organic compounds. This technique fundamentally relies on the alignment tensor determination from RDCs and the *a priori* determination of the EFG either by first principles (DFT methods)^[162b] or solid state NMR spectroscopy^[61c]. While this method is far from being fully developed

i.e. various technical issues, in particular the sign determination of experimental RQCs, have to be addressed in future studies, its potential was shown in a case study on CpLi in THF. The combination of RQC/RDC and theoretical calculations of the EFG led to $[\text{CpLi}(\text{THF})_2]$ as the solution structure of CpLi in THF. This is in accordance with DOSY *MW* methods.

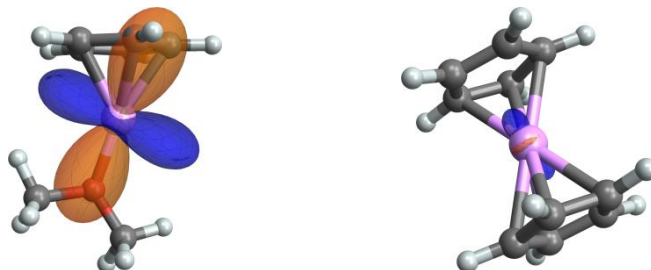


Figure 4-23: Spherical polar plot EFG tensor representation around the lithium atom in hypothetical $[\text{CpLi}(\text{OMe}_2)]$ (left) and $[\text{Cp}_2\text{Li}]^-$ (right).

It is important to mention that RQCs are, in contrast to DOSY, also able to clarify the distinction between molecular aggregates with small variations in *MW*. For example the hypothetical $[\text{CpLi}(\text{THF})]$ and $[\text{Cp}_2\text{Li}]^-$ have very similar *MW* but can be clearly distinguished by the size of the reproduction of RQCs through theoretical EFGs and experimental RDCs.

5 EXPERIMENTAL PART

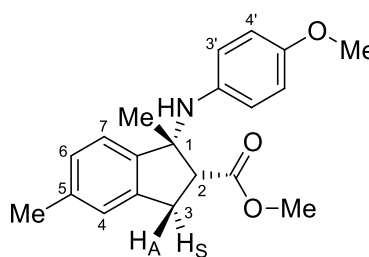
5.1 Sample Preparation

Synthesis and Handling of Air- and Moisture-Sensitive Compounds

All experimental standard *Schlenk* procedures took place under inert gas atmosphere (purified, dry argon) in a glove box. Glassware was dried at 130°C, assembled hot and cooled down under vacuum. All solvents were dried and distilled over sodium-potassium alloy and finally degassed (three *pump-freeze* cycles) prior to use. Chemicals were purchased as commercially available products or synthesized according to known literature. *n*BuLi was kindly donated by Rockwood Lithium.

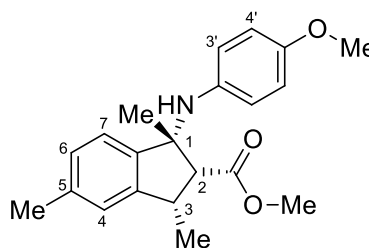
Methyl-1-((4-methoxyphenyl)amino)-1,5-dimethyl-2,3-dihydro-1*H*-indene-2-carboxylate

Formula: $C_{20}H_{23}NO_3$
Molecular weight: 339.2 g/mol



Methyl-1-((4-methoxyphenyl)amino)-1,3,5-trimethyl-2,3-dihydro-1*H*-indene-2-carboxylate

Formula: $C_{21}H_{25}NO_3$
Molecular weight: 367.2 g/mol



The indane compounds were synthesized by *Liu* and *Zell* from the *Ackermann* group and then employed as obtained. Experimental characterization has already been performed in reference [148].

Cyclopentadienide-lithium

Formula: C_5H_5Li
Molecular weight: 72.0 g/mol



CpLi was synthesized by *Bachmann* from the *Stalke* group and then employed as obtained. Experimental characterization has already been performed in reference [165a].

Preparation of Polymer Sticks

Styrene, DVB (industrial grade, ~80% *m*-DVB), butyl acrylate and 1,4-butanediol diacrylate were separated from inhibitors by flash column chromatography through neutral aluminium oxide (50 g). Mixtures of DVB in styrene and 1,4-butanediol diacrylate in butyl acrylate were prepared according to the desired volume percentage (0.2%-0.5%). Under argon atmosphere, these monomer/linker solutions were degassed (three *pump-freeze* cycles) and filled into glass tubes sealed at one end (10 cm length, 3.8 mm inner diameter). These tubes were then heated to 115°C over three days to accomplish the polymerization.

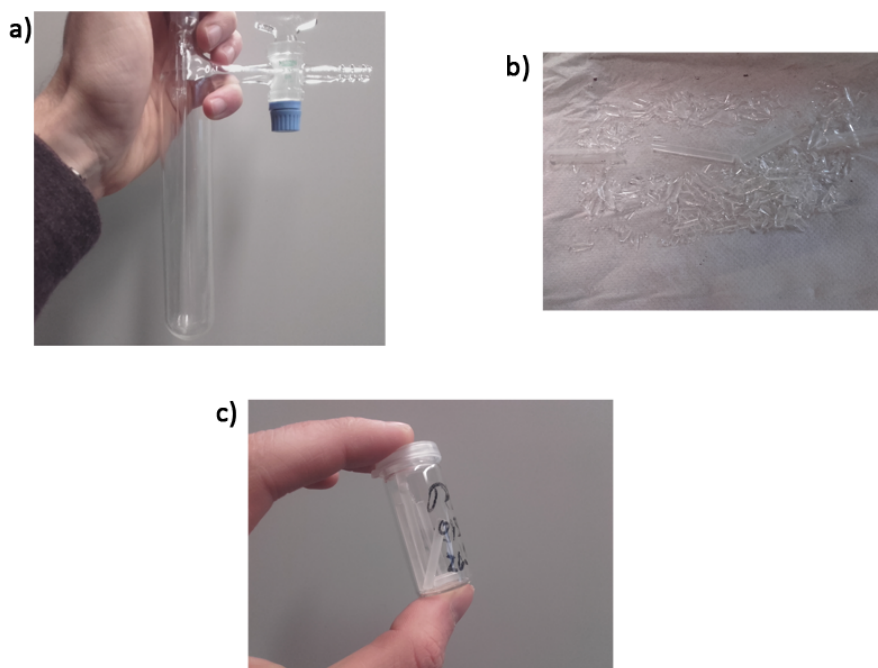
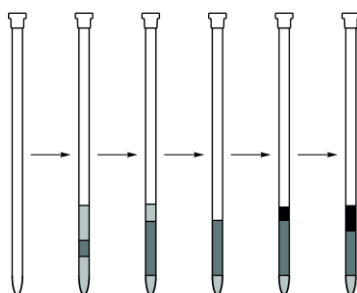


Figure 5-1: Preparation of the polymer sticks: a) Filling of glass tubes with mixtures of PS and DVB under inert gas atmosphere within a *Schlenk* tube and subsequent heating at 115°C over three days. b) Careful breaking of the glass tubes gives the ready-to-use polymer sticks (c).

Note that swollen samples of polymers were usually shimmed *on the FID* i.e. the FID area was maximized by manual shim routines.

Observation of Reactions



Scheme 5-1: Basic steps of sample preparation: 1. A crosslinked polymer is immersed in a solution which contains the reactive (lithium organic) species. 2. The polymer is swollen/equilibrated over a sufficiently long time period. This can be evaluated through homogeneous signal intensities via SSE as shown by Pöppler *et al.* 3. The supernatant solution above the polymer is removed. 4. A solution containing a further substrate is added on top of the polymer and diffuses into the gel. 5. This is then followed by SSE over an appropriate time period.

Due to the high reactivity of organolithium reagents and to prevent rapid convection and/or diffusion out of the active sample volume, *n*BuLi was first diffused into a PS gel matrix, and all steps of the preparation were carried out strictly under argon atmosphere. For this purpose, a solution of *n*BuLi in *n*-hexane (0.4 ml, 1.9 M, 0.74 mmol) was concentrated *in vacuo* and toluene-*d*₈ (0.40 ml) was added. This solution was then transferred into a 5 mm NMR tube, and a cylindrical PS stick (10 × 3.8 mm, crosslinked through 0.2 vol% divinylbenzene) was immersed in the solution approximately 1 cm above the bottom of the tube. After 7 days (the PS stick had swollen to a length of ~3 cm) the supernatant *n*BuLi/toluene solution above the polymer was removed and replaced by a solution of PMDTA (0.12 ml, 0.60 mmol) in toluene-*d*₈ (0.15 ml). Slice-selective ¹H and ⁷Li NMR measurements started after ~3 h and went on for three days.

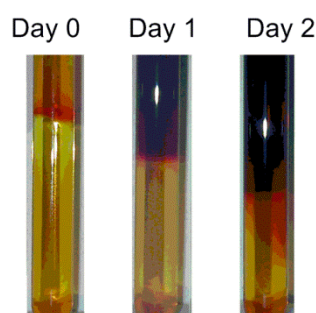
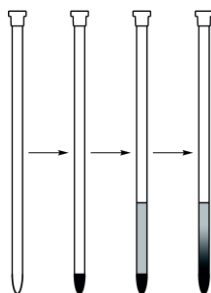


Figure 5-2: Visual tracing of the PMDTA diffusing into the *n*BuLi/PS gel over the course of three days.

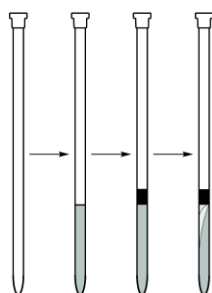
Single-Shot Titration



Scheme 5-2: Basic steps of sample preparation: 1. The first compound (neat) is filled into the NMR tube. 2. The second compound, immersed in deuterated solvent, is carefully added on top of the first compound. The compound at the bottom then diffuses into the solution above. 3. This can be readily observed through SSE.

12-crown-4 (50 μ L, 3.1 mmol, m.p. 16 $^{\circ}$ C) was filled into a standard 5 mm NMR tube and cooled to 5 $^{\circ}$ C. Then a solution of LiClO₄ (24 mg, 2.3 mmol) in acetonitrile-d₃ (0.45 ml) was layered on top of the solid ether. This procedure prevents initial mixing of both components prior to the measurements. Inside the NMR magnet (25 $^{\circ}$ C, standing tube) the ether melts and slowly diffuses into the LiClO₄ solution resulting in a smooth concentration gradient along the tube axis. Approximately 3, 6 and 9 h after sample preparation slice-selective ¹H and ⁷Li NMR spectra were recorded.

Qualitative Chromatography



Scheme 5-3: Basic steps of sample preparation: 1. The stationary phase is prepared within the NMR tube. 2. A mixture of compounds to separate is added on top of the stationary phase. 3. The compounds start diffusing into the mixture at different rates. 4. The diffusion progress can be readily observed through SSE.

As a case study the separation of amino acids L-Asp acid, L-His, L-Tyr and L-Val within agar hydrogel was examined. This gel matrix has lately been described as a beneficial diffusion controlling medium in an elegant titration experiment by *Mitrev et al.* Agar (7.8 mg, 2wt%) was immersed in D₂O (0.35 ml) in an NMR tube, and the suspension was heated until a homogeneous gel resulted. A mixture of L-aspartic acid (asp, 7.3 mg), L-His (his, 8.3 mg), L-Tyr (tyr, 8.0 mg) and L-Val (val, 6.5 mg) in slightly alkaline D₂O (0.15 ml) was added on top of the gel sample and SSE spectra (19 slices, 1 mm thickness each) were recorded immediately afterwards as well as 2, 6 and 10 h later. This procedure was repeated for gel samples with varying amounts of agar weight-percentages (3wt% and 5wt%).

Anisotropic NMR Spectroscopy

In general, about 20 mg of the target compound were dissolved in an appropriate deuterated solvent (small molecules: CD_2Cl_2 , indanes: CDCl_3 , lithium-compounds: $\text{THF-}d_8$). Alignment of the compounds was achieved by strain-induced alignment in a gel (SAG)^[75b, 91, 95, 169]: Either a cylindrical PS stick (10×3.8 mm, crosslinked through 0.3 vol% divinylbenzene) or a cylindrical PBA stick (10×3.8 mm, crosslinked through 0.6 vol% 1,4-butanediol diacrylate) was immersed in the isotropic solutions. Swelling/equilibration over several days (small molecules: three days for PBA and four days for PS, indanes: four days, lithium compounds: eight days) led to sufficiently homogeneous alignment within the gel. The polymer homogeneity was evaluated by eye as well as the ^2H quadrupolar splitting of solvent signals in deuterium SSE experiments.

5.2 NMR Experiments

Measurements were performed at 25°C unless stated otherwise. Measurements were always performed on non-spinning samples. Spectra were recorded and processed with TopSpin 2.1 (Bruker Biospin, Rheinstetten, Germany). Further processing was done with MestreNova 10.0 (Mestrelab Research, Santiago de Compostela, Spain). Chemical shifts are reported in ppm relative to TMS (^1H , ^2H) and LiCl (^7Li). J -couplings, RDCs and residual quadrupolar couplings are given in units of Hz.

Special care was taken for the shimming of heterogeneous samples. Automated gradient shimming is usually not possible due to low signal-to-noise ratios or split deuterium signals. Likewise manual shimming *on the lock signal* proved impracticable due to low signal-to-noise ratios as well. Hence, shimming *on the FID* involved continuous acquisition of FIDs without saving but with parallel integration of the FID area. Homogeneously shimmed samples display larger integrals thereby enabling a quantitative observable to estimate and improve the field homogeneity.

Imaging Experiment

Imaging experiments were performed on a Bruker Avance III 400 spectrometer with a standard BBI probe with z gradient. The pulse program *echo_image.gau* written in Bruker Avance syntax is given in the Appendix. After a 90° hard pulse the magnetization develops under a gradient applied at 20% maximum strength for 10 ms. This is followed by a 180° pulse and signal acquisition under a gradient applied at 20% maximum strength.

SSE

SSE NMR spectra were measured on a Bruker Avance III 400 spectrometer with a standard BBI probe with z gradient. The pulse program *zg_slice.gau* written in Bruker Avance syntax is given in the Appendix.

Slice-selective ^1H experiments were performed at 400.13 MHz with a gradient strength of 25.6 G/cm. *Gauss-Cascade G4* shaped pulses were used to excite a bandwidth of ~ 11000 Hz which corresponds to ~ 1 mm thick slices. Frequency offsets for these pulses ranged between 98209 Hz and -98209 Hz with steps of 10912 Hz. This resulted in 19 slices with 1 mm distance between each slice.

Slice-selective ^2H experiments were performed at 61.4 MHz with a gradient strength of 51.1 G/cm. *G4* shaped pulses were used to excite a bandwidth of ~ 1700 Hz which corresponds to ~ 0.5 mm thick slices. Frequency offsets for these pulses ranged between 23560 Hz and -23560 Hz with steps of 3366 Hz. This resulted in 15 slices with 1 mm distance between each slice.

Slice-selective ^7Li experiments were performed at 155.5 MHz with a gradient strength of 51.1 G/cm. *G4* shaped pulses were used to excite a bandwidth of ~ 8500 Hz which corresponds to ~ 1 mm thick slices. Frequency offsets for these pulses ranged between 76723 Hz and -76723 Hz with steps of 8525 Hz. This resulted in 19 slices with 1 mm distance between each slice.

The Bruker Topspin automation program *poptau* offers a simple method to vary different parameters, among them are frequency offsets. Additionally, results obtained by *poptau* can be saved as pseudo-2D spectra. Further customization of this AU can easily achieve a standardized and fast acquisition procedure.

CPMG

CPMG NMR spectra were measured on a Bruker Avance III 400 spectrometer with a standard BBI probe with z gradient. The echo time before and after the π pulse (19.08 μs) was 2 ms.

Transverse relaxation was monitored by echo trains of 1, 2, 5, 10, 20, 40, 80, 150, 250, 450 and 700 cycles (corresponding to a maximum of 1.82 s between π /pulse and signal acquisition).

DOSY

DOSY measurements were performed on a Bruker Avance III 400 spectrometer equipped with a standard BBI probe with z gradient. The pulse program *dstebpgp3s* was applied in the Bruker Topspin automation program *xau dosy*.

Herein, the pulsed gradient strength was incremented in 16 steps from 2% to 98%. The diffusion time (50 ms) and the gradient (1.38 ms) were optimized beforehand to achieve a $\sim 100:5$ signal attenuation between maximum and minimum gradient strength.

CLIP-HSQC

CLIP-HSQC spectra were recorded using an INEPT delay of $1/(4 \times {}^1J_{\text{CH}})$ with ${}^1J_{\text{CH}} = 145$ Hz. For the small molecules in the F1 dimension a total of 128 data points were sampled over a spectral width of 200 ppm. For the indanes a total of 256 data points were sampled over a spectral width of 200 ppm. In the F2 dimension 2048 data points were recorded with a spectral width of 15 ppm. The number of scans was 4 for the isotropic as well as anisotropic sample.

NOESY

NOESY spectra were recorded on a Bruker Avance III 500 MHz instrument, using a relaxation delay of 2 s. In the F1 dimension a total of 1024 data points were sampled over a spectral width of 8 ppm. In the F2 dimension 2048 data points were recorded with a spectral width of 8 ppm. Number of scans was 16.

HMBC

Conventional HMBC spectra were recorded on a Bruker Avance III 500 MHz instrument, using a relaxation delay of 2 s. In the F1 dimension a total of 512 data points were sampled over a spectral width of 180 ppm. In the F2 dimension 4096 data points were recorded with a spectral width of 8 ppm, corresponding to an acquisition time of 0.5 s. Number of scans was 8. The J -couplings were determined from the peak multiplets along F2 with an accuracy of ~ 1 Hz.

5.3 Processing of NMR Data

Homogeneous Reference Samples

Homogeneous reference samples for ^7Li were prepared by dissolving LiClO_4 in toluene- d_8 to obtain three different solutions with known concentration (0.49, 0.58, 0.73 mol/L) were prepared and measured via SSE. The signals were integrated and corrected according to the sensitivity profile. Homogeneous reference samples for ^1H were prepared by solving acetone in toluene- d_8 to obtain three different solutions with known concentration (1.97, 2.19, 2.46 mol/L) were prepared and measured via SSE. The signals were integrated and corrected according to the sensitivity profile. Concentration c can then be obtained from intensities I through $c = m \cdot I$ with $m(\text{H}) = 8.4 \cdot 10^{-4}$ mol/L and $m(\text{Li}) = 1.71$ mol/L. The obtained calibration curves are depicted in Figure 5-3.

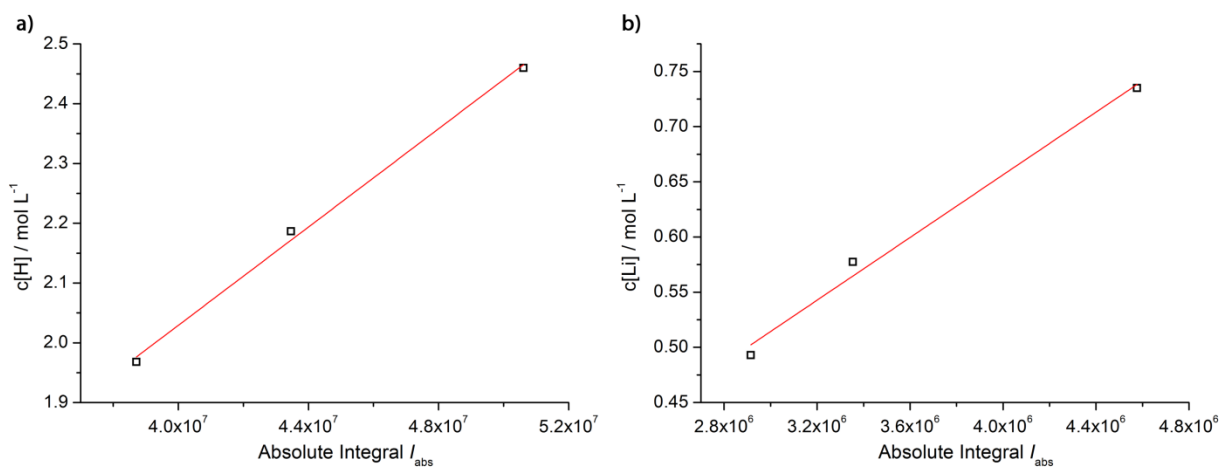


Figure 5-3: Calibration curves for the concentration determination of a) hydrogen-containing molecules and b) lithium-containing molecules.

RDC Data Analysis

Isotropic measurements yielded $^1J_{\text{CH}}$ couplings by CLIP-HSQC.^[68] Further CLIP-HSQC spectra of anisotropic samples were recorded to yield couplings $^1T_{\text{CH}}$ that are the sums of scalar and RDCs J and D , respectively: $^1T_{\text{CH}} = ^1J_{\text{CH}} + ^1D_{\text{CH}}$.

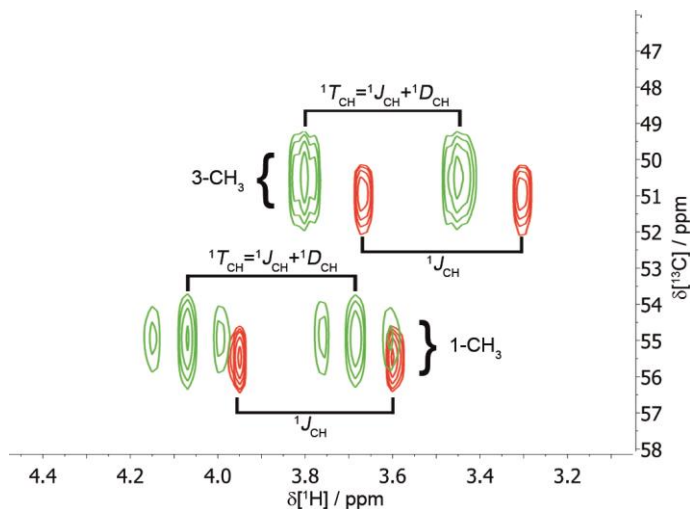


Figure 5-4: Excerpt of a CLIP-HSQC spectrum of β -aa-Me. Red: Isotropic $^1J_{\text{CH}}$ couplings. Green: Anisotropic $^1T_{\text{CH}}$ couplings. Note that for methyl groups an additional triplet appears through $^2D_{\text{HH}}$ couplings.

RDCs were saved as *.lab file. These contain the atom numbers of the C–H pairs and their RDC values in Hz. In the case of prochiral H atoms with unclear assignments a swapping of these RDCs can be achieved by the command “permutations”. An exemplary *.lab file is included in the Appendix.

Within MSpin errors were computed with an estimated standard deviation of 0.5 Hz. Furthermore methyl was treated as freely rotating and otherwise the defaults were applied. MSpin computes values for back-calculated RDCs and additionally gives a more detailed output concerning the alignment of the molecules. An exemplary output file is included in Appendix B Exemplary MSpin Input and Output.

DOSY Data Analysis

Spectra as acquired via the *xau dosy* routine were processed with the help of the Topspin T_1/T_2 relaxation module (Figure 5-5).

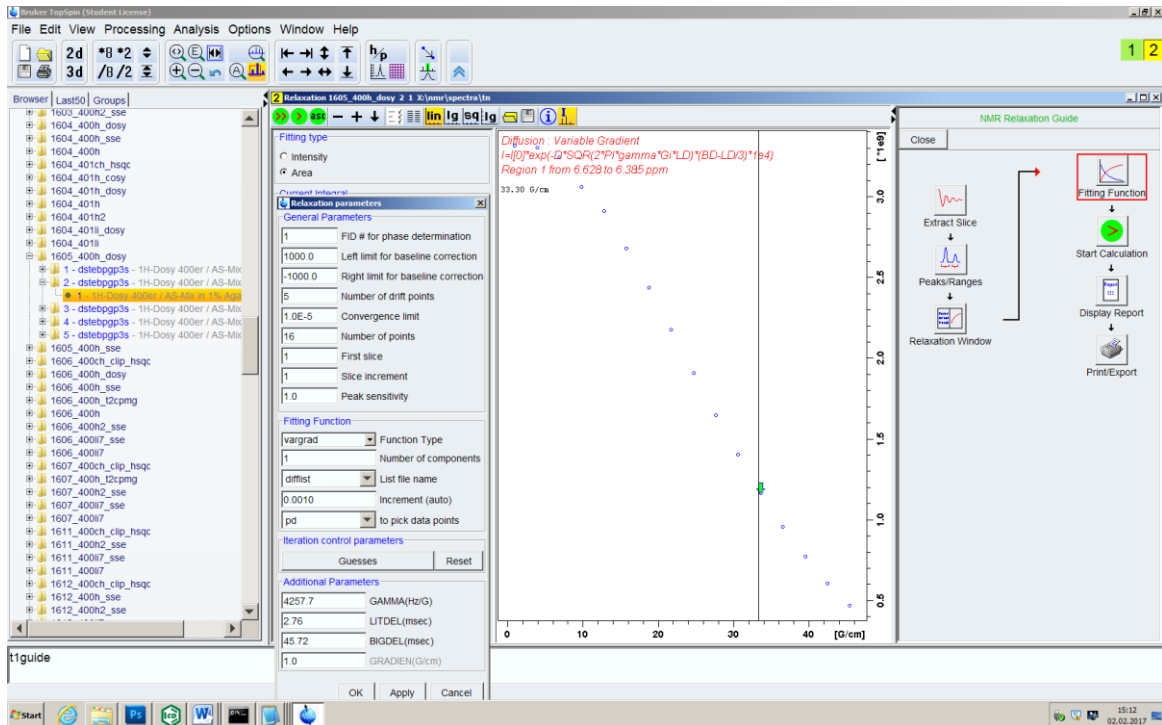


Figure 5-5: Screenshot of DOSY data processing for the extraction of diffusion coefficients from *Gaussian* fitting of the signal intensity or area decay.

This module reads either intensity or integral data for chosen signals and directly yields diffusion coefficients from *Gaussian* decay fitting.

5.4 Computational Techniques

MD Simulations and Analysis

MD Simulations were performed using GROMACS 5.1.2^[139] and the established OPLS-AA force field.^[140,170] Topologies were built according to GROMACS guidelines and with OPLS standard charges. Prior to the production run, systems were energy optimized. The corresponding production run file is given in Appendix B Exemplary MD Simulation Run File. MD production runs were performed in small cubic boxes with periodic boundary conditions (PS and PBA without DCM: 1.95 nm, PS and PBA with DCM: 2.05 nm, Figure 5-6).

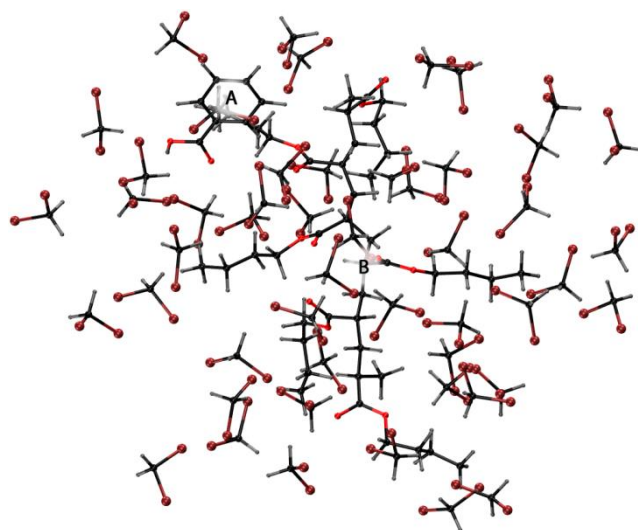


Figure 5-6: Screenshots of MD simulation box containing benzoic acid (A), PBA (B) and DCM molecules.

To mimic solvent effects in cases where no DCM molecules were added, Langevin dynamics were set on to add friction terms. During 400 ns production runs the Berendsen thermostat^[141] was set to a target temperature of 300 K with no preceding thermal equilibration. MD runs were analyzed with the GROMACS built-in utility gmx angle, the VMD^[171] molecular dynamics visualizing software and *M. Brehm's* Travis^[138] as a further tool to analyze trajectories from MD runs. In particular, Travis allowed the determination of spatial distribution functions as well as combined distribution functions.

DFT Calculations of Optimized Structures and EFG Analysis

DFT calculations were performed using ORCA 3.0.1.^[154a] Exemplary input files are provided in Appendix B Exemplary ORCA input files geometry optimization and numerical frequencies. Optimized structures were evaluated as minima by numerically calculating vibrational frequencies which in turn must not include (large) imaginary wavenumbers. Electrostatic potentials were calculated through the CHELPG^[138] routine as implemented in ORCA. EFG tensors were calculated using the fgrad keyword. Graphical representation of the EFG tensors by spherical polar plots was achieved by a *Mathematica*^[172] script provided publically^[173] by *J. Autschbach* (Figure 5-7).

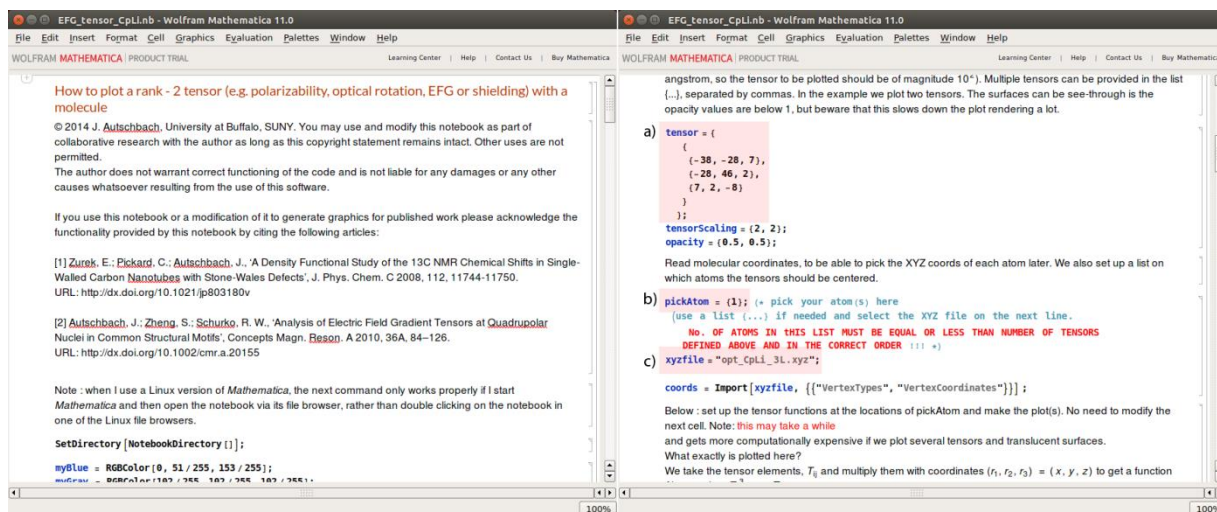


Figure 5-7: Screenshots of *Mathematica* script for tensor representation. Right window: a) Multiple of EFG as determined from ORCA. b) Number of atom with EFG in xyz file. c) xyz file correlated to ORCA EFG calculation.

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A Applications of Slice-Selective NMR Spectroscopy

Preparation of SSE Experiment Data

Imaging Pulse Sequence

Pulse program *echo_image.gau* in Bruker Avance syntax.

```

;gradcalib
#include <Avance.incl>
#include <Grad.incl>
1 ze
2 d1 groff
p1 ph1
3u UNBLKGRAD
10m gron1
3u groff
p2 ph2
6u gron1
go=2 ph31
30m mc #0 to 2 F0(zd)
3u groff
3u BLKGRAD
exit
ph1=0
ph2=1 3
ph31=0
;p10 : 0W
;p11 : f1 channel - power level for pulse (default)
;sp1: f1 channel - shaped pulse
;p11: f1 channel - 90 (or 270) degree shaped pulse

```

SSE Pulse Sequence

Pulse program *zg_slice.gau* in Bruker Avance syntax.

```

;zg-slice.gau
;avance-version
;selective excitation using a shaped pulse
;$CLASS=HighRes
;$DIM=1D
#include <Avance.incl>
#include <Grad.incl>
1 ze
2 30m
d1
3u UNBLKGRAD
3u gron1
(p11:sp1 ph1):f1
3u groff
3u BLKGRAD
go=2 ph31
30m mc #0 to 2 F0(zd)
exit
ph1=0 2 2 0 1 3 3 1
ph31=0 2 2 0 1 3 3 1
;sp1: f1 channel - shaped pulse
;p11: f1 channel - pulse length

```

;d1 : relaxation delay; 1-5 * T1
;ns: 1 * n, total number of scans: NS * TD0
;choose p11 according to desired selectivity i.e. slice-width the flip-angle is determined by the amplitude
;choose frequency offsets according to desired slice distance for a given gradient strength

Observation of Reactions Data

Exemplary set of spectra:

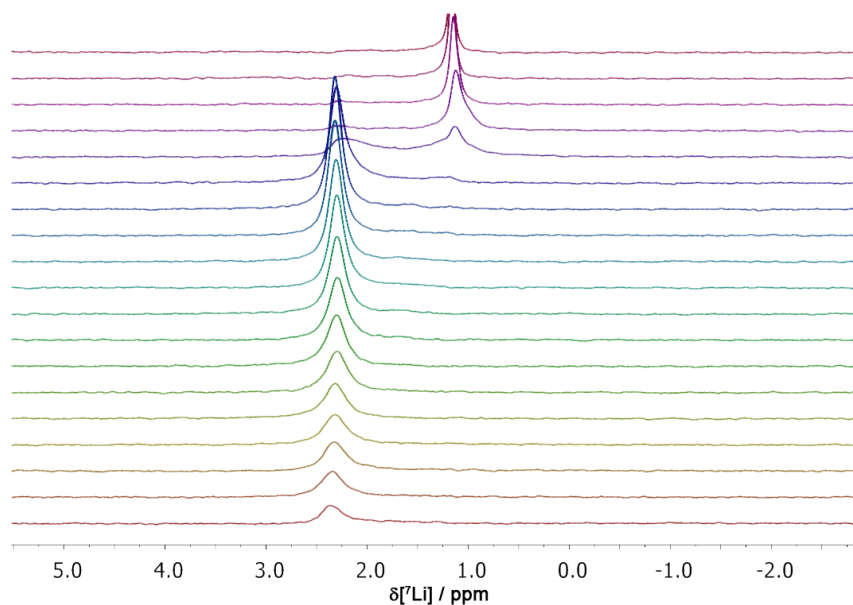


Figure A-1: SSE ⁷Li spectra for the lithiation of PMDTA by *n*BuLi. Spectra were taken approximately 3 h after addition of PMDTA.

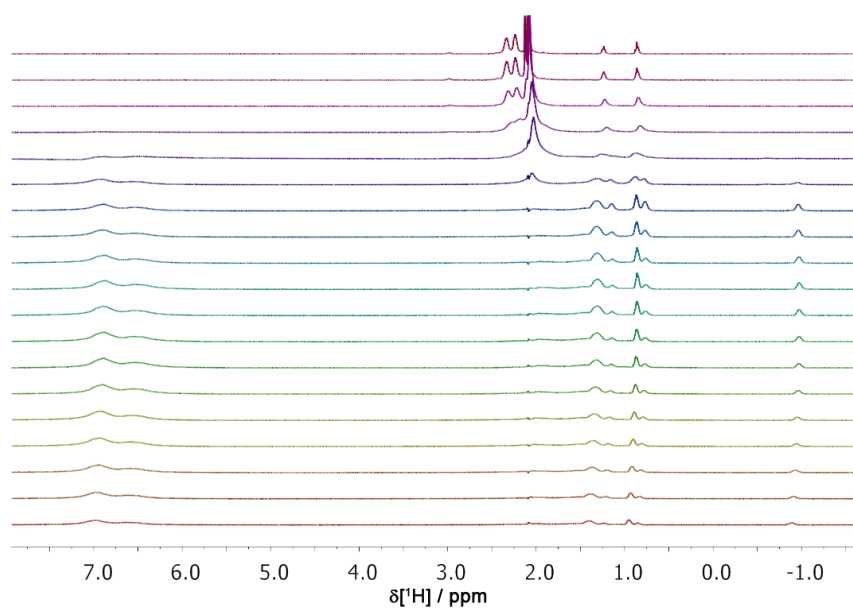


Figure A-2: SSE ¹H spectra for the lithiation of PMDTA by *n*BuLi. Spectra were taken approximately 3 h after addition of PMDTA.

The full set of spectra is stored at the Stalke group server and can be obtained upon request.

Single-Shot Titration Data

Exemplary set of spectra:

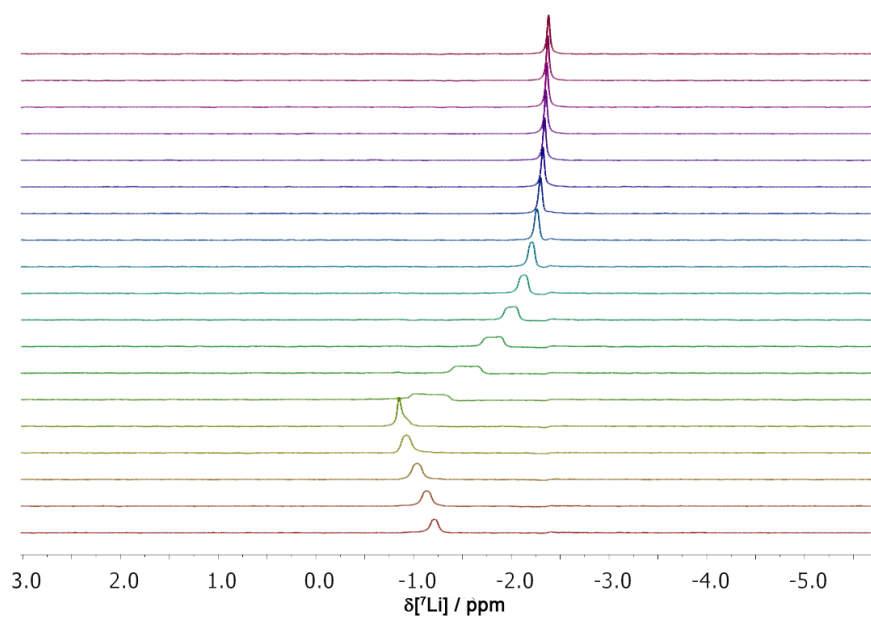


Figure A-3: SSE ${}^7\text{Li}$ spectra for the titration of Li^+ by 12-crown-4. Spectra were taken approximately 3 h after preparation of the sample.

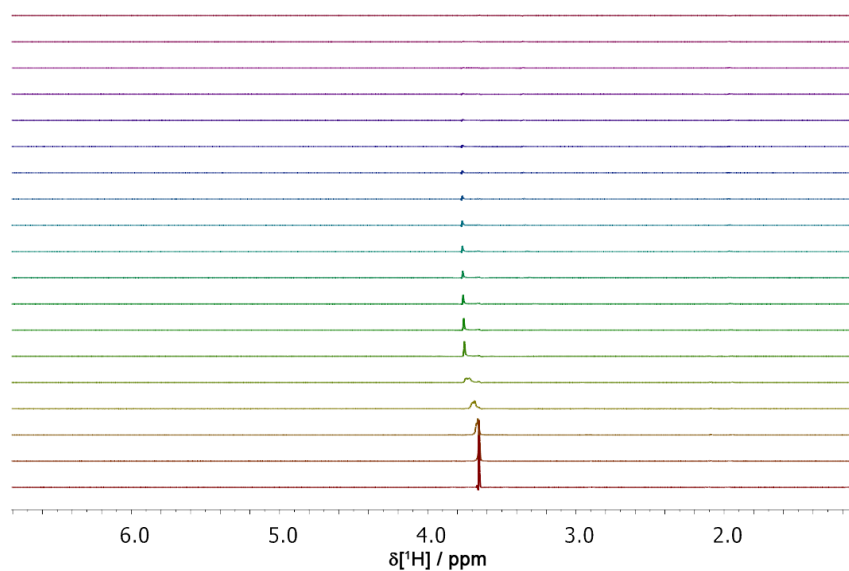


Figure A-4: SSE ${}^1\text{H}$ spectra for the titration of Li^+ by 12-crown-4. Spectra were taken approximately 3 h after preparation of the sample.

The full set of spectra is stored at the Stalke group server and can be obtained upon request.

Qualitative Chromatography Data

The full set of spectra is stored at the Stalke group server and can be obtained upon request. Exemplary output file from the T_1/T_2 relaxation module and analysis of the signal of L-His at the low field limit of the spectrum.

AREA fit : Diffusion : Variable Gradient :

$$I=I[0]*\exp(-D*\text{SQR}(2*\text{PI}*\text{gamma}*Gi*LD)*(BD-LD/3)*1e4)$$

16 points for Integral 1, Integral Region from 7.681 to 7.465 ppm

Converged after 57 iterations!

Results Comp. 1

I[0] = 1.888e-002
 Diff Con. = 4.967e-010 m²/s
 Gamma = 4.258e+003 Hz/G
 Little Delta = 2.760m
 Big Delta = 45.720m

RSS = 2.015e-006

SD = 3.549e-004

Point	Gradient	Expt	Calc	Difference
1	9.270e-001	1.799e-002	1.886e-002	8.780e-004
2	3.893e+000	1.886e-002	1.854e-002	-3.232e-004
3	6.860e+000	1.810e-002	1.784e-002	-2.625e-004
4	9.826e+000	1.702e-002	1.680e-002	-2.232e-004
5	1.279e+001	1.552e-002	1.548e-002	-3.748e-005
6	1.576e+001	1.410e-002	1.397e-002	-1.260e-004
7	1.873e+001	1.258e-002	1.234e-002	-2.332e-004
8	2.169e+001	1.089e-002	1.067e-002	-2.202e-004
9	2.466e+001	8.344e-003	9.032e-003	6.885e-004
10	2.762e+001	7.615e-003	7.484e-003	-1.308e-004
11	3.059e+001	5.888e-003	6.069e-003	1.811e-004
12	3.356e+001	5.017e-003	4.818e-003	-1.990e-004
13	3.652e+001	3.633e-003	3.745e-003	1.119e-004
14	3.949e+001	3.270e-003	2.848e-003	-4.212e-004
15	4.246e+001	1.992e-003	2.121e-003	1.286e-004
16	4.542e+001	1.184e-003	1.546e-003	3.616e-004

Relaxation Data

Table A-1: Relaxation data from CPMG experiments.

t/s	2wt% Agarose	3wt% Agarose	5wt% Agarose
0.00804	1	1	1
0.01206	0.9294	0.91439	0.9074
0.02412	0.86402	0.82709	0.768
0.04422	0.75438	0.69152	0.58333
0.08442	0.56542	0.4788	0.33846
0.16481	0.32781	0.23659	0.11695
0.3256	0.1129	0.05845	0.01509
0.60699	0.01809	0.00549	6.58E-04
1.00897	0.00148	2.40E-04	5.94E-05
1.81293	2.00E-05	9.45E-06	2.37E-05

B Anisotropic NMR Spectroscopy in Polymeric Gels

Alignment of Small Molecules Data

Least-Square Solution Script

```

import random
import os
import numpy as np
import scipy as sp
from scipy import special
rdc_list_name = raw_input("Enter name of file containing RDCs (one column) ")
#rdc_list_name = "d_toluene"
angle_list_name = raw_input("Enter name of file containing geometric information (solved sin2cos2,... coefficients) ")
#angle_list_name = "a_toluene"
rdc_list = np.genfromtxt(str(os.path.realpath("__file__")[:-8])+str(rdc_list_name))
angle_list = np.genfromtxt(str(os.path.realpath("__file__")[:-8])+str(angle_list_name))
rdc_array = np.asarray(rdc_list)
no_of_rdc = rdc_array.shape[0]
rdc_array = rdc_array/-69959.8
rdc_array.shape = (no_of_rdc, 1)
angle_array = np.asarray(angle_list)
angle_array.shape = (no_of_rdc,3)
print str(angle_list_name)+"="
print angle_array
print str(rdc_list_name)+"="
print rdc_array
alignment = np.linalg.lstsq(angle_array, rdc_array)
alignment_tensor_solution=open(str(os.path.realpath("__file__")[:-8])+str("alignment_"+str(angle_list_name)+"-"+str(rdc_list_name)), "w")
alignment_tensor_solution.write(str(alignment[0]))
alignment_tensor_solution.close()

```

MD Energy Optimization Run File

```

; minim.mdp - used as input into grompp to generate em.tpr
integrator      = steep          ; Algorithm (steep = steepest descent minimization)
emtol          = 50.0           ; Stop minimization when the maximum force < 1000.0 kJ/mol/nm
emstep        = 0.01           ; Energy step size
nsteps         = 50000         ; Maximum number of (minimization) steps to perform

; Parameters describing how to find the neighbors of each atom and how to calculate the interactions
nstlist        = 1             ; Frequency to update the neighbor list and long range forces
cutoff-scheme  = Verlet
ns_type        = grid          ; Method to determine neighbor list (simple, grid)
coulombtype    = PME           ; Treatment of long range electrostatic interactions
rcoulomb       = 1.0           ; Short-range electrostatic cut-off
rvdw           = 1.0           ; Short-range Van der Waals cut-off
pbc            = xyz           ; Periodic Boundary Conditions (yes/no)

```

Exemplary MD Simulation Run File

Adapted MD simulation script with comments:

```

; run.mdp - used as input into grompp to generate run.tpr
; Run parameters
integrator      = md           ; leap-frog integrator

```

```

nsteps          = 200000000      ; 2 * 200000000 = 400000 ps (400 ns)
dt              = 0.002          ; 2 fs
; Output control
nstxout         = 5000            ; save coordinates every 10.0 ps
nstvout         = 5000            ; save velocities every 10.0 ps
nstenergy       = 5000            ; save energies every 10.0 ps
nstlog          = 5000            ; update log file every 10.0 ps
nstxout-compressed = 5000      ; save compressed coordinates every 10.0 ps
                    ; nstxout-compressed replaces nstxtcout
compressed-x-grps = System      ; replaces xtc-grps
; Bond parameters
continuation    = yes            ; Restarting after NPT
constraint_algorithm = lincs      ; holonomic constraints
constraints     = all-bonds      ; all bonds (even heavy atom-H bonds) constrained
lincs_iter      = 1              ; accuracy of LINCS
lincs_order     = 4              ; also related to accuracy
; Neighborsearching
cutoff-scheme   = Verlet
ns_type         = grid           ; search neighboring grid cells
nstlist         = 10             ; 20 fs, largely irrelevant with Verlet scheme
rcoulomb        = 1.0           ; short-range electrostatic cutoff (in nm)
rvdw            = 1.0           ; short-range van der Waals cutoff (in nm)
; Electrostatics
coulombtype     = PME            ; Particle Mesh Ewald for long-range electrostatics
pme_order       = 4              ; cubic interpolation
fourierspacing  = 0.16         ; grid spacing for FFT
; Temperature coupling is on
tcoupl         = v-rescale       ; modified Berendsen thermostat
tc_grps        = System         ; One group
tau_t          = 0.1            ; time constant, in ps
ref_t          = 300            ; reference temperature, one for each group, in K
; Pressure coupling is on
pcoupl         = no             ; no pressure coupling
; Periodic boundary conditions
pbc            = xyz            ; 3-D PBC
; Dispersion correction
DispCorr       = EnerPres       ; account for cut-off vdW scheme
; Velocity generation
gen_vel        = no            ; Velocity generation is off

```

Exemplary OPLS Topologies

Polybutylacrylate Topology

See reference [132] for original simulation and topology on PS. Note that below only shortened topologies of PBA, toluene and benzoic acid are shown. These do not include entries for bonds, pairs and proper dihedrals. The full topologies are stored at the Stalke group server and can be obtained upon request.

Polybutylacrylate

```

[ moleculetype ]
; Name          nrexcl
pba             3

[ atoms ]
; nr   type resnr residue atom  cgnr  charge  mass typeB  chargeB  massB
; residue 1 P1S rtp P1S q 0.0
1  opls_135  1  P1S  C1   1  -0.18  12.011 ; qtot -0.18
2  opls_140  1  P1S  H11  1   0.06   1.008 ; qtot -0.12
3  opls_140  1  P1S  H12  1   0.06   1.008 ; qtot -0.06
4  opls_140  1  P1S  H13  1   0.06   1.008 ; qtot 0
5  opls_135  1  P1S  C2   1  -0.06  12.011 ; qtot -0.06
6  opls_140  1  P1S  H21  1   0.06   1.008 ; qtot 0

```

```
7  opls_465  1  P1S  C3  2  0.51  12.011 ; qtot 0.51
8  opls_466  1  P1S  O1  2  -0.43  15.9994 ; qtot 0.08
9  opls_467  1  P1S  O2  2  -0.33  15.9994 ; qtot -0.25
10 opls_490  1  P1S  C4  2  0.19  12.011 ; qtot -0.06
11 opls_469  1  P1S  H41  2  0.03  1.008 ; qtot -0.03
12 opls_469  1  P1S  H42  2  0.03  1.008 ; qtot 0
13 opls_135  1  P1S  C5  3  -0.12  12.011 ; qtot -0.12
14 opls_140  1  P1S  H51  3  0.06  1.008 ; qtot -0.06
15 opls_140  1  P1S  H52  3  0.06  1.008 ; qtot 0
16 opls_135  1  P1S  C6  3  -0.12  12.011 ; qtot -0.12
17 opls_140  1  P1S  H61  3  0.06  1.008 ; qtot -0.06
18 opls_140  1  P1S  H62  3  0.06  1.008 ; qtot 0
19 opls_135  1  P1S  C7  3  -0.18  12.011 ; qtot -0.18
20 opls_140  1  P1S  H71  3  0.06  1.008 ; qtot -0.12
21 opls_140  1  P1S  H72  3  0.06  1.008 ; qtot -0.06
22 opls_140  1  P1S  H73  3  0.06  1.008 ; qtot 0
; residue 2 P1 rtp P1 q 0.0
23 opls_135  2  P1  C1  4  -0.12  12.011 ; qtot -0.12
24 opls_140  2  P1  H11  4  0.06  1.008 ; qtot -0.06
25 opls_140  2  P1  H12  4  0.06  1.008 ; qtot 0
26 opls_135  2  P1  C2  4  -0.06  12.011 ; qtot -0.06
27 opls_140  2  P1  H21  4  0.06  1.008 ; qtot 0
28 opls_465  2  P1  C3  5  0.51  12.011 ; qtot 0.51
29 opls_466  2  P1  O1  5  -0.43  15.9994 ; qtot 0.08
30 opls_467  2  P1  O2  5  -0.33  15.9994 ; qtot -0.25
31 opls_490  2  P1  C4  5  0.19  12.011 ; qtot -0.06
32 opls_469  2  P1  H41  5  0.03  1.008 ; qtot -0.03
33 opls_469  2  P1  H42  5  0.03  1.008 ; qtot 0
34 opls_135  2  P1  C5  6  -0.12  12.011 ; qtot -0.12
35 opls_140  2  P1  H51  6  0.06  1.008 ; qtot -0.06
36 opls_140  2  P1  H52  6  0.06  1.008 ; qtot 0
37 opls_135  2  P1  C6  6  -0.12  12.011 ; qtot -0.12
38 opls_140  2  P1  H61  6  0.06  1.008 ; qtot -0.06
39 opls_140  2  P1  H62  6  0.06  1.008 ; qtot 0
40 opls_135  2  P1  C7  6  -0.18  12.011 ; qtot -0.18
41 opls_140  2  P1  H71  6  0.06  1.008 ; qtot -0.12
42 opls_140  2  P1  H72  6  0.06  1.008 ; qtot -0.06
43 opls_140  2  P1  H73  6  0.06  1.008 ; qtot 0
; residue 3 P1 rtp P1 q 0.0
44 opls_135  3  P1  C1  7  -0.12  12.011 ; qtot -0.12
45 opls_140  3  P1  H11  7  0.06  1.008 ; qtot -0.06
46 opls_140  3  P1  H12  7  0.06  1.008 ; qtot 0
47 opls_135  3  P1  C2  7  -0.06  12.011 ; qtot -0.06
48 opls_140  3  P1  H21  7  0.06  1.008 ; qtot 0
49 opls_465  3  P1  C3  8  0.51  12.011 ; qtot 0.51
50 opls_466  3  P1  O1  8  -0.43  15.9994 ; qtot 0.08
51 opls_467  3  P1  O2  8  -0.33  15.9994 ; qtot -0.25
52 opls_490  3  P1  C4  8  0.19  12.011 ; qtot -0.06
53 opls_469  3  P1  H41  8  0.03  1.008 ; qtot -0.03
54 opls_469  3  P1  H42  8  0.03  1.008 ; qtot 0
55 opls_135  3  P1  C5  9  -0.12  12.011 ; qtot -0.12
56 opls_140  3  P1  H51  9  0.06  1.008 ; qtot -0.06
57 opls_140  3  P1  H52  9  0.06  1.008 ; qtot 0
58 opls_135  3  P1  C6  9  -0.12  12.011 ; qtot -0.12
59 opls_140  3  P1  H61  9  0.06  1.008 ; qtot -0.06
60 opls_140  3  P1  H62  9  0.06  1.008 ; qtot 0
61 opls_135  3  P1  C7  9  -0.18  12.011 ; qtot -0.18
62 opls_140  3  P1  H71  9  0.06  1.008 ; qtot -0.12
63 opls_140  3  P1  H72  9  0.06  1.008 ; qtot -0.06
64 opls_140  3  P1  H73  9  0.06  1.008 ; qtot 0
; residue 4 P1 rtp P1 q 0.0
65 opls_135  4  P1  C1  10 -0.12  12.011 ; qtot -0.12
66 opls_140  4  P1  H11  10  0.06  1.008 ; qtot -0.06
67 opls_140  4  P1  H12  10  0.06  1.008 ; qtot 0
68 opls_135  4  P1  C2  10 -0.06  12.011 ; qtot -0.06
69 opls_140  4  P1  H21  10  0.06  1.008 ; qtot 0
```

```
70 opls_465 4 P1 C3 11 0.51 12.011 ; qtot 0.51
71 opls_466 4 P1 O1 11 -0.43 15.9994 ; qtot 0.08
72 opls_467 4 P1 O2 11 -0.33 15.9994 ; qtot -0.25
73 opls_490 4 P1 C4 11 0.19 12.011 ; qtot -0.06
74 opls_469 4 P1 H41 11 0.03 1.008 ; qtot -0.03
75 opls_469 4 P1 H42 11 0.03 1.008 ; qtot 0
76 opls_135 4 P1 C5 12 -0.12 12.011 ; qtot -0.12
77 opls_140 4 P1 H51 12 0.06 1.008 ; qtot -0.06
78 opls_140 4 P1 H52 12 0.06 1.008 ; qtot 0
79 opls_135 4 P1 C6 12 -0.12 12.011 ; qtot -0.12
80 opls_140 4 P1 H61 12 0.06 1.008 ; qtot -0.06
81 opls_140 4 P1 H62 12 0.06 1.008 ; qtot 0
82 opls_135 4 P1 C7 12 -0.18 12.011 ; qtot -0.18
83 opls_140 4 P1 H71 12 0.06 1.008 ; qtot -0.12
84 opls_140 4 P1 H72 12 0.06 1.008 ; qtot -0.06
85 opls_140 4 P1 H73 12 0.06 1.008 ; qtot 0
; residue 5 P1 rtp P1 q 0.0
86 opls_135 5 P1 C1 13 -0.12 12.011 ; qtot -0.12
87 opls_140 5 P1 H11 13 0.06 1.008 ; qtot -0.06
88 opls_140 5 P1 H12 13 0.06 1.008 ; qtot 0
89 opls_135 5 P1 C2 13 -0.06 12.011 ; qtot -0.06
90 opls_140 5 P1 H21 13 0.06 1.008 ; qtot 0
91 opls_465 5 P1 C3 14 0.51 12.011 ; qtot 0.51
92 opls_466 5 P1 O1 14 -0.43 15.9994 ; qtot 0.08
93 opls_467 5 P1 O2 14 -0.33 15.9994 ; qtot -0.25
94 opls_490 5 P1 C4 14 0.19 12.011 ; qtot -0.06
95 opls_469 5 P1 H41 14 0.03 1.008 ; qtot -0.03
96 opls_469 5 P1 H42 14 0.03 1.008 ; qtot 0
97 opls_135 5 P1 C5 15 -0.12 12.011 ; qtot -0.12
98 opls_140 5 P1 H51 15 0.06 1.008 ; qtot -0.06
99 opls_140 5 P1 H52 15 0.06 1.008 ; qtot 0
100 opls_135 5 P1 C6 15 -0.12 12.011 ; qtot -0.12
101 opls_140 5 P1 H61 15 0.06 1.008 ; qtot -0.06
102 opls_140 5 P1 H62 15 0.06 1.008 ; qtot 0
103 opls_135 5 P1 C7 15 -0.18 12.011 ; qtot -0.18
104 opls_140 5 P1 H71 15 0.06 1.008 ; qtot -0.12
105 opls_140 5 P1 H72 15 0.06 1.008 ; qtot -0.06
106 opls_140 5 P1 H73 15 0.06 1.008 ; qtot 0
; residue 6 P1E rtp P1E q 0.0
107 opls_135 6 P1E C1 16 -0.12 12.011 ; qtot -0.12
108 opls_140 6 P1E H11 16 0.06 1.008 ; qtot -0.06
109 opls_140 6 P1E H12 16 0.06 1.008 ; qtot 0
110 opls_135 6 P1E C2 16 -0.06 12.011 ; qtot -0.06
111 opls_140 6 P1E H21 16 0.06 1.008 ; qtot 0
112 opls_465 6 P1E C3 17 0.51 12.011 ; qtot 0.51
113 opls_466 6 P1E O1 17 -0.43 15.9994 ; qtot 0.08
114 opls_467 6 P1E O2 17 -0.33 15.9994 ; qtot -0.25
115 opls_490 6 P1E C4 17 0.19 12.011 ; qtot -0.06
116 opls_469 6 P1E H41 17 0.03 1.008 ; qtot -0.03
117 opls_469 6 P1E H42 17 0.03 1.008 ; qtot 0
118 opls_135 6 P1E C5 18 -0.12 12.011 ; qtot -0.12
119 opls_140 6 P1E H51 18 0.06 1.008 ; qtot -0.06
120 opls_140 6 P1E H52 18 0.06 1.008 ; qtot 0
121 opls_135 6 P1E C6 18 -0.12 12.011 ; qtot -0.12
122 opls_140 6 P1E H61 18 0.06 1.008 ; qtot -0.06
123 opls_140 6 P1E H62 18 0.06 1.008 ; qtot 0
124 opls_135 6 P1E C7 18 -0.18 12.011 ; qtot -0.18
125 opls_140 6 P1E H71 18 0.06 1.008 ; qtot -0.12
126 opls_140 6 P1E H72 18 0.06 1.008 ; qtot -0.06
127 opls_140 6 P1E H73 18 0.06 1.008 ; qtot 0
128 opls_135 6 P1E C8 19 -0.18 12.011 ; qtot -0.18
129 opls_140 6 P1E H81 19 0.06 1.008 ; qtot -0.12
130 opls_140 6 P1E H82 19 0.06 1.008 ; qtot -0.06
131 opls_140 6 P1E H83 19 0.06 1.008 ; qtot 0
```

[bonds]


```

; ai aj funct c0 c1 c2 c3
  1 2 1

[ pairs ]
; ai aj funct c0 c1 c2 c3
  1 8 1

[ angles ]
; ai aj ak funct c0 c1 c2 c3
  2 1 3 1

[ dihedrals ]
; ai aj ak al funct c0 c1 c2 c3 c4 c5
  2 1 5 6 3

[ dihedrals ]
; ai aj ak al funct c0 c1 c2 c3
  5 8 7 9 1 improper_O_C_X_Y
 26 29 28 30 1 improper_O_C_X_Y
 47 50 49 51 1 improper_O_C_X_Y
 68 71 70 72 1 improper_O_C_X_Y
 89 92 91 93 1 improper_O_C_X_Y
110 113 112 114 1 improper_O_C_X_Y

[ position_restraints ]
; atom type fx fy fz
  1 1 10000 10000 10000
128 1 10000 10000 10000

```

Toluene

```

[ moleculetype ]
; Name nrexcl
tol 3

[ atoms ]
; nr type resnr residue atom cgnr charge mass typeB chargeB massB
; residue 1 TOL rtp TOL q 0.0
  1 opls_148 1 TOL C1 1 -0.065 12.011 ;qtot -0.065
  2 opls_140 1 TOL H11 1 0.06 1.008 ;qtot -0.005
  3 opls_140 1 TOL H12 1 0.06 1.008 ;qtot 0.055
  4 opls_140 1 TOL H13 1 0.06 1.008 ;qtot 0.115
  5 opls_145 1 TOL C2 2 -0.115 12.011 ;qtot 0
  6 opls_145 1 TOL C3 3 -0.115 12.011 ;qtot -0.115
  7 opls_146 1 TOL H31 3 0.115 1.008 ;qtot 0
  8 opls_145 1 TOL C4 4 -0.115 12.011 ;qtot -0.115
  9 opls_146 1 TOL H41 4 0.115 1.008 ;qtot 0
 10 opls_145 1 TOL C5 5 -0.115 12.011 ;qtot -0.115
 11 opls_146 1 TOL H51 5 0.115 1.008 ;qtot 0
 12 opls_145 1 TOL C6 6 -0.115 12.011 ;qtot -0.115
 13 opls_146 1 TOL H61 6 0.115 1.008 ;qtot 0
 14 opls_145 1 TOL C7 7 -0.115 12.011 ;qtot -0.115
 15 opls_146 1 TOL H71 7 0.115 1.008 ;qtot 0

[ bonds ]
; ai aj funct c0 c1 c2 c3
  1 2 1

[ pairs ]
; ai aj funct c0 c1 c2 c3
  1 7 1

[ angles ]
; ai aj ak funct c0 c1 c2 c3
  2 1 3 1

```

```
[ dihedrals ]
; ai aj ak al funct      c0      c1      c2      c3      c4      c5
  2  1  5  6  3
```

```
[ dihedrals ]
; ai aj ak al funct      c0      c1      c2      c3
  1  5  6  7  1 improper_Z_CA_X_Y
  1  5 14 12  1 improper_Z_CA_X_Y
  5  6  8  9  1 improper_Z_CA_X_Y
  6  8 10 11  1 improper_Z_CA_X_Y
  8 10 12 13  1 improper_Z_CA_X_Y
 10 12 14 15  1 improper_Z_CA_X_Y
```

Benzoic Acid

```
[ moleculetype ]
; Name      nrexcl
bza        3
```

```
[ atoms ]
; nr  type resnr residue atom  cgnr  charge  mass typeB  chargeB  massB
; residue 1 BZA rtp BZA q 0.0
  1  opls_470  1  BZA  C1  1  0.635  12.011 ;qtot 0.635
  2  opls_268  1  BZA  O1  1  -0.53  15.9994 ;qtot 0.105
  3  opls_269  1  BZA  O2  1  -0.44  15.9994 ;qtot -0.335
  4  opls_270  1  BZA  H1  1  0.45  1.008 ;qtot 0.115
  5  opls_145  1  BZA  C2  2  -0.115  12.011 ;qtot 0
  6  opls_145  1  BZA  C3  3  -0.115  12.011 ;qtot -0.115
  7  opls_146  1  BZA  H31  3  0.115  1.008 ;qtot 0
  8  opls_145  1  BZA  C4  4  -0.115  12.011 ;qtot -0.115
  9  opls_146  1  BZA  H41  4  0.115  1.008 ;qtot 0
 10  opls_145  1  BZA  C5  5  -0.115  12.011 ;qtot -0.115
 11  opls_146  1  BZA  H51  5  0.115  1.008 ;qtot 0
 12  opls_145  1  BZA  C6  6  -0.115  12.011 ;qtot -0.115
 13  opls_146  1  BZA  H61  6  0.115  1.008 ;qtot 0
 14  opls_145  1  BZA  C7  7  -0.115  12.011 ;qtot -0.115
 15  opls_146  1  BZA  H71  7  0.115  1.008 ;qtot 0
```

```
[ bonds ]
; ai aj funct      c0      c1      c2      c3
  1  2  1
```

```
[ pairs ]
; ai aj funct      c0      c1      c2      c3
  1  7  1
```

```
[ angles ]
; ai aj ak funct      c0      c1      c2      c3
  2  1  3  1
```

```
[ dihedrals ]
; ai aj ak al funct      c0      c1      c2      c3      c4      c5
  3  1  2  4  3
```

```
[ dihedrals ]
; ai aj ak al funct      c0      c1      c2      c3
  1  5  6  7  1 improper_Z_CA_X_Y
  1  5 14 12  1 improper_Z_CA_X_Y
  2  5  1  3  1 improper_O_C_X_Y
  5  6  8  9  1 improper_Z_CA_X_Y
  6  8 10 11  1 improper_Z_CA_X_Y
  8 10 12 13  1 improper_Z_CA_X_Y
 10 12 14 15  1 improper_Z_CA_X_Y
```

Structural Elucidation of Synthetic Indanes Data

Exemplary ORCA input files geometry optimization and numerical frequencies

Optimization:

```
# ORCA JOB on forcefield optimized structure
#-----
! RKS B3LYP D3 def2-TZVP def2-TZVP/J RIJCOSX SmallPrint TightSCF Grid4 Cosmo(solvens) Opt
%maxcore X
%pal nprocs Y
#choose X and Y according to hardware
*xyzfile NAME.xyz A B
#coordinate file forcefield optimized structure; choose A and B according to charge and multiplicity
```

Numerical calculation of vibrational frequencies:

```
# ORCA JOB on ORCA optimized structure
#-----
! RKS B3LYP D3 def2-TZVP def2-TZVP/J RIJCOSX SmallPrint TightSCF Grid4 Cosmo(solvens)
%maxcore X
%pal nprocs Y
#choose X and Y according to hardware
! MORead
    %moinp "NAME.gbwn"
! NumFreq
*xyzfile NAME.xyz A B
#coordinate and gbwn file ORCA optimized structure; choose A and B according to charge and multiplicity
```

Structure Data β -aa-H*RS-*eq*

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -2.9251 0.3532 -0.2972	C -2.8223 0.3442 -0.3608	C -2.9109 0.3597 -0.3481	C -2.8602 0.4367 -0.1465
C -2.9267 -1.0442 -0.3460	C -2.8382 -1.0030 -0.3722	C -2.8864 -1.0480 -0.3925	C -2.8603 -0.9673 -0.1465
C -1.7014 -1.7187 -0.3393	C -1.6704 -1.6698 -0.3095	C -1.6593 -1.7217 -0.3461	C -1.6589 -1.6794 -0.1402
C -0.5158 -1.0018 -0.2859	C -0.5216 -0.9845 -0.2380	C -0.4908 -0.9773 -0.2914	C -0.4730 -0.9652 -0.1465
C -0.5289 0.3891 -0.2372	C -0.5194 0.3531 -0.2348	C -0.5165 0.4224 -0.2572	C -0.4729 0.4341 -0.1465
C -1.7336 1.0751 -0.2413	C -1.6715 1.0334 -0.2912	C -1.7275 1.1004 -0.2773	C -1.6586 1.1486 -0.1401
C -4.2181 -1.8149 -0.4316	C -4.1548 -1.7396 -0.4516	C -4.1765 -1.8147 -0.4318	C -4.1610 -1.7083 -0.1533
C 0.8981 -1.5181 -0.2607	C 0.8707 -1.5475 -0.1631	C 0.9121 -1.4931 -0.2009	C 0.9250 -1.4975 -0.1036
C 1.7266 -0.2773 -0.6327	C 1.7107 -0.3163 -0.5798	C 1.7134 -0.2457 -0.6187	C 1.7268 -0.2657 -0.5483
C 0.8812 0.9648 -0.1534	C 0.8746 0.9168 -0.1477	C 0.8878 0.9855 -0.1501	C 0.9252 0.9662 -0.1035
N 1.1538 2.1025 -1.0286	N 1.0448 2.0296 -1.0937	N 1.0862 2.1291 -1.0467	N 1.0987 2.0645 -1.0153
C 3.1249 -0.2974 -0.0835	C 3.1442 -0.3238 -0.0849	C 3.1122 -0.2984 -0.0700	C 3.0858 -0.2659 0.1075
H 0.7043 2.9389 -0.6712	H 0.4885 2.8453 -0.7855	H 0.4830 2.8964 -0.7579	H 0.4339 2.8038 -0.7874
H 0.7701 1.9348 -1.9532	H 0.6512 1.7715 -2.0151	H 0.7870 1.8731 -1.9874	H 0.8394 1.7692 -1.9566
H 4.8656 0.3451 -0.4389	H 4.7987 0.5410 -0.3754	H 4.8807 0.1042 -0.4476	H 4.7471 0.5245 0.3755
O 3.9947 0.3621 -0.8706	O 3.8940 0.5959 -0.7247	O 4.0335 0.1677 -0.9352	O 3.9501 0.7241 -0.1438
O 3.4611 -0.8247 0.9527	O 3.6198 -1.0451 0.7578	O 3.4559 -0.7182 1.0232	O 3.3935 -1.1598 0.8595
C 1.1860 1.3831 1.2896	C 1.1706 1.3868 1.2833	C 1.1589 1.4278 1.2943	C 1.3397 1.5264 1.2507
H -3.8700 0.8840 -0.3059	H -3.7752 0.8976 -0.4084	H -3.8640 0.8853 -0.3604	H -3.8158 0.9814 -0.1516
H -1.6838 -2.8025 -0.3832	H -1.6426 -2.7705 -0.3144	H -1.6185 -2.8073 -0.3510	H -1.6551 -2.7793 -0.1305
H -1.7578 2.1585 -0.2075	H -1.6728 2.1346 -0.2793	H -1.7679 2.1837 -0.2289	H -1.6546 2.2485 -0.1302
H -4.1936 -2.6962 0.2119	H -4.0283 -2.8453 -0.4495	H -4.0126 -2.8535 -0.7369	H -3.9645 -2.8038 -0.1523
H -5.0670 -1.1982 -0.1357	H -4.7987 -1.4751 0.4180	H -4.6432 -1.8131 0.5580	H -4.7471 -1.4362 0.7529
H -4.3992 -2.1678 -1.4514	H -4.6944 -1.4664 -1.3869	H -4.8660 -1.3665 -1.1548	H -4.7381 -1.4360 -1.0652
H 1.0679 -2.3446 -0.9515	H 1.0207 -2.4162 -0.8428	H 1.0869 -2.3227 -0.8932	H 1.0591 -2.3475 -0.8095
H 1.1709 -1.8686 0.7378	H 1.0660 -1.8532 0.8903	H 1.1274 -1.8212 0.8214	H 1.2278 -1.9070 0.8860
H 1.7934 -0.1927 -1.7180	H 1.7551 -0.3322 -1.6992	H 1.7770 -0.2654 -1.7184	H 1.8829 -0.2657 -1.6503
H 2.1932 1.7949 1.3735	H 2.1997 1.8028 1.3675	H 2.1715 1.8313 1.4091	H 1.6452 2.5903 1.1349
H 0.4742 2.1499 1.6024	H 0.4644 2.1907 1.5926	H 0.4593 2.2110 1.6117	H 0.4815 1.4621 1.9566
H 1.0953 0.5393 1.9723	H 1.0769 0.5528 2.0151	H 1.0475 0.5966 2.0008	H 2.1947 0.9363 1.6501

RS-*ax*

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -1.9869 0.9857 0.7836	C -1.8701 0.9603 0.7311	C -1.9420 0.9935 0.7360	C -1.7309 1.0962 0.7185
C -1.8547 -0.4062 0.7849	C -1.7482 -0.3812 0.7533	C -1.8236 -0.4089 0.7896	C -1.7309 -0.3078 0.7185
C -0.6028 -1.0160 0.7264	C -0.5438 -0.9764 0.7483	C -0.5704 -1.0287 0.7686	C -0.5295 -1.0198 0.7248
C 0.5344 -0.2227 0.6624	C 0.5552 -0.2115 0.7156	C 0.5647 -0.2347 0.6926	C 0.6564 -0.3056 0.7184
C 0.4150 1.1641 0.6539	C 0.4457 1.1220 0.7054	C 0.4502 1.1616 0.6512	C 0.6564 1.0938 0.7185
C -0.8299 1.7691 0.7192	C -0.7555 1.7144 0.7071	C -0.7872 1.7857 0.6826	C -0.5294 1.8081 0.7249
C 1.9949 -0.6549 0.6751	C 1.9954 -0.6534 0.7178	C 2.0189 -0.6527 0.7052	C 2.0544 -0.8377 0.7613
C 2.7260 0.6678 0.2141	C 2.6909 0.6337 0.1939	C 2.7053 0.6416 0.1821	C 2.8562 0.3941 0.3166
C 1.7682 1.8204 0.5898	C 1.7878 1.7973 0.6705	C 1.7993 1.8162 0.6089	C 2.0544 1.6260 0.7614
C 2.9838 0.6543 -1.2704	C 2.8042 0.6459 -1.3261	C 2.8833 0.6554 -1.3223	C 3.0192 0.3942 -1.1835
O 4.2318 0.2390 -1.5620	O 4.0564 0.3640 -1.7387	O 4.0817 0.1633 -1.7015	O 3.6900 1.3842 -1.7835
O 2.1881 0.9631 -2.1268	O 1.9213 0.8735 -2.1179	O 2.0995 1.0523 -2.1673	O 2.5441 -0.4997 -1.8428
N 2.4752 -0.9631 2.0404	N 2.3941 -0.8904 2.1159	N 2.4239 -0.9099 2.1021	N 2.4194 -1.2522 2.0890
C 2.2713 -1.8775 -0.1911	C 2.2713 -1.9125 -0.1128	C 2.3108 -1.9150 -0.1106	C 2.2794 -2.0742 -0.0990
C -3.3448 1.6346 0.8227	C -3.2260 1.6262 0.7339	C -3.2934 1.6440 0.7895	C -3.0315 1.8374 0.7115
H -2.7454 -1.0221 0.8306	H -2.6510 -1.0138 0.7784	H -2.7165 -1.0277 0.8501	H -2.6866 -0.8524 0.7134
H -0.5305 -2.0974 0.7322	H -0.4636 -2.0742 0.7741	H -0.4996 -2.1104 0.8179	H -0.5256 -2.1197 0.7345
H -0.9136 2.8505 0.7108	H -0.8214 2.8136 0.6885	H -0.8573 2.8693 0.6580	H -0.5254 2.9081 0.7348

Appendix

H 3.6850 0.7311 0.7223	H 3.7183 0.7430 0.6154	H 3.7048 0.7867 0.6143	H 3.8737 0.3940 0.7676
H 1.8051 2.6324 -0.1360	H 1.8149 2.6787 -0.0077	H 1.8226 2.6557 -0.0932	H 2.1887 2.4761 0.0557
H 2.0467 2.2419 1.5594	H 2.0552 2.1186 1.7042	H 2.0633 2.1708 1.6112	H 2.3573 2.0353 1.7511
H 4.3261 0.2184 -2.5295	H 4.0645 0.3799 -2.7098	H 4.0678 0.2320 -2.6796	H 3.6713 1.1848 -2.7347
H 1.9596 -1.7613 2.3998	H 1.9094 -1.7268 2.4838	H 1.8854 -1.7003 2.4577	H 1.7631 -1.9611 2.4164
H 2.2599 -0.1920 2.6658	H 2.0646 -0.1098 2.7098	H 2.1456 -0.1178 2.6827	H 2.2959 -0.4723 2.7347
H 1.7365 -2.7406 0.2102	H 1.8083 -2.8136 0.3488	H 1.8332 -2.7998 0.3284	H 2.6493 -2.9081 0.5385
H 1.9300 -1.7259 -1.2153	H 1.8703 -1.8284 -1.1476	H 1.9310 -1.8365 -1.1355	H 1.3199 -2.3717 -0.5782
H 3.3357 -2.1130 -0.1948	H 3.3665 -2.1041 -0.1778	H 3.3871 -2.1158 -0.1586	H 3.0330 -1.8464 -0.8857
H -4.1038 0.9422 1.1883	H -4.0645 0.8945 0.7331	H -4.0823 0.9442 0.4944	H -3.8737 1.1097 0.7070
H -3.3440 2.5138 1.4691	H -3.3395 2.2652 1.6391	H -3.5000 1.9945 1.8054	H -3.0987 2.4762 1.6204
H -3.6492 1.9677 -0.1741	H -3.3429 2.2686 -0.1685	H -3.3384 2.4945 0.1011	H -3.0891 2.4768 -0.1976

SS-eg

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -2.7815 0.3272 -0.0226	C -2.7737 0.2922 -0.1036	C -2.8450 0.3121 -0.0849	C -2.8602 0.3204 -0.2065
C -2.8209 -1.0664 -0.1415	C -2.8091 -1.0542 -0.0920	C -2.8563 -1.0959 -0.1188	C -2.8603 -1.0836 -0.2065
C -1.6151 -1.7705 -0.1897	C -1.6480 -1.7333 -0.0492	C -1.6452 -1.7991 -0.0879	C -1.6589 -1.7957 -0.2002
C -0.4084 -1.0864 -0.1355	C -0.4891 -1.0615 -0.0228	C -0.4578 -1.0833 -0.0627	C -0.4730 -1.0815 -0.2065
C -0.3848 0.3000 -0.0393	C -0.4630 0.2756 -0.0456	C -0.4479 0.3166 -0.0448	C -0.4729 0.3178 -0.2065
C -1.5724 1.0144 0.0350	C -1.6115 0.9652 -0.0786	C -1.6421 1.0229 -0.0442	C -1.6586 1.0322 -0.2001
C -4.1350 -1.7991 -0.2074	C -4.1366 -1.7745 -0.1197	C -4.1645 -1.8318 -0.1214	C -4.1610 -1.8247 -0.2133
C 0.9853 -1.6511 -0.1587	C 0.8904 -1.6501 0.0303	C 0.9323 -1.6296 0.0050	C 0.9250 -1.6138 -0.1636
C 1.8630 -0.4307 -0.4773	C 1.7616 -0.4433 -0.3930	C 1.7585 -0.4089 -0.4453	C 1.7268 -0.3821 -0.6083
C 1.0285 0.8549 -0.0174	C 0.9445 0.8225 -0.0082	C 0.9699 0.8514 0.0311	C 0.9252 0.8499 -0.1635
N 1.3257 1.3458 1.3310	N 1.2594 1.2443 1.3675	N 1.2351 1.2823 1.4083	N 1.2902 1.2641 1.1643
C 3.2143 -0.4732 0.1704	C 3.1847 -0.5330 0.1294	C 3.1809 -0.5453 0.0294	C 3.0858 -0.3822 0.0475
O 4.1184 0.2831 -0.4868	O 3.8493 0.6309 0.0034	O 3.8608 0.6139 0.0891	O 3.9501 0.6078 -0.2038
O 3.4888 -1.0666 1.1883	O 3.7217 -1.5084 0.5965	O 3.7588 -1.5929 0.2748	O 3.3935 -1.2762 0.7995
C 1.2236 2.0061 -0.9998	C 1.1481 1.9904 -0.9848	C 1.1899 2.0413 -0.9120	C 1.1503 2.0864 -1.0236
H -3.7120 0.8814 0.0180	H -3.7199 0.8582 -0.1320	H -3.7835 0.8630 -0.0785	H -3.8158 0.8651 -0.2116
H -1.6263 -2.8511 -0.2825	H -1.6340 -2.8342 -0.0324	H -1.6317 -2.8854 -0.0823	H -1.6551 -2.8957 -0.1904
H -1.5639 2.0933 0.1299	H -1.6068 2.0661 -0.0820	H -1.6490 2.1074 0.0059	H -1.6546 2.1322 -0.1902
H -4.9396 -1.1393 -0.5336	H -4.0246 -2.8811 -0.0833	H -4.0396 -2.8602 -0.4762	H -3.9645 -2.9202 -0.2123
H -4.0816 -2.6422 -0.8979	H -4.7591 -1.4729 0.7533	H -4.5802 -1.8596 0.8905	H -4.7471 -1.5526 0.6929
H -4.4102 -2.2017 0.7722	H -4.6914 -1.5242 -1.0526	H -4.8801 -1.3412 -0.7894	H -4.7381 -1.5523 -1.1252
H 1.1140 -2.4445 -0.8960	H 1.0107 -2.5240 -0.6489	H 1.0721 -2.4738 -0.6775	H 1.0591 -2.4639 -0.8695
H 1.2507 -2.0717 0.8143	H 1.0928 -1.9558 1.0831	H 1.1668 -1.9432 1.0284	H 1.2278 -2.0233 0.8260
H 2.0050 -0.3492 -1.5546	H 1.8439 -0.4749 -1.5089	H 1.8031 -0.4139 -1.5454	H 1.8829 -0.3820 -1.7103
H 2.2568 1.7474 1.3787	H 2.2515 1.5245 1.4529	H 2.2189 1.5364 1.4961	H 2.2897 1.4652 1.1950
H 1.2834 0.5967 2.0135	H 1.1322 0.4514 2.0202	H 1.0641 0.5129 2.0534	H 1.1666 0.4841 1.8098
H 4.9551 0.2537 0.0071	H 4.7591 0.4974 0.3155	H 4.7722 0.3435 0.3341	H 4.7471 0.4082 0.3155
H 0.6104 2.8574 -0.7034	H 0.5512 2.8811 -0.6861	H 0.5797 2.9059 -0.6241	H 1.5203 2.9202 -0.3860
H 2.2681 2.3258 -1.0169	H 2.2108 2.3166 -1.0170	H 2.2347 2.3689 -0.9158	H 1.9045 1.8585 -1.8098
H 0.9375 1.7018 -2.0064	H 0.8471 1.7121 -2.0202	H 0.9174 1.7884 -1.9437	H 0.1912 2.3840 -1.5035

SS-ax

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -1.8240 0.8352 0.7568	C -1.7900 0.6978 0.6095	C -1.8494 0.7196 0.5821	C -1.7308 0.8734 0.6749
C -1.7578 -0.5612 0.7448	C -1.7616 -0.6487 0.6485	C -1.8330 -0.6882 0.6367	C -1.7309 -0.5307 0.6749
C -0.5355 -1.2282 0.6831	C -0.6011 -1.3254 0.6585	C -0.6275 -1.3953 0.6353	C -0.5296 -1.2427 0.6812
C 0.6341 -0.4854 0.6310	C 0.5462 -0.6355 0.6260	C 0.5612 -0.6841 0.5788	C 0.6563 -0.5286 0.6748
C 0.5841 0.9039 0.6455	C 0.5316 0.7016 0.5952	C 0.5495 0.7157 0.5322	C 0.6564 0.8708 0.6749
C -0.6318 1.5654 0.7084	C -0.6265 1.3744 0.5825	C -0.6398 1.4275 0.5472	C -0.5293 1.5852 0.6813
C 2.0594 -1.0015 0.6005	C 1.9490 -1.1847 0.6353	C 1.9782 -1.2076 0.6102	C 2.0543 -1.0608 0.7177

C 2.8729 0.3131 0.1961	C 2.7398 0.0458 0.1097	C 2.7577 0.0317 0.0620	C 2.8561 0.1709 0.2730
C 1.9667 1.4971 0.5945	C 1.9192 1.2780 0.5631	C 1.9446 1.2650 0.5181	C 2.0545 1.4028 0.7179
C 3.1390 0.3024 -1.2823	C 2.8748 0.0798 -1.4085	C 2.9106 0.0767 -1.4475	C 3.0192 0.1710 -1.2271
O 4.3799 -0.1397 -1.5704	O 3.9479 0.8020 -1.7893	O 3.5882 1.1708 -1.8602	O 3.6900 1.1610 -1.8271
O 2.3489 0.6138 -2.1436	O 2.1568 -0.4475 -2.2233	O 2.5730 -0.7625 -2.2657	O 2.5440 -0.7228 -1.8863
N 2.1812 -2.1476 -0.2924	N 2.0729 -2.3600 -0.2422	N 2.1380 -2.4617 -0.1254	N 2.2276 -2.1591 -0.1942
C 2.5005 -1.4400 2.0023	C 2.3695 -1.5612 2.0645	C 2.3836 -1.4938 2.0690	C 2.4687 -1.6212 2.0719
C -3.1487 1.5475 0.8287	C -3.0974 1.4549 0.5960	C -3.1528 1.4623 0.6248	C -3.0314 1.6146 0.6679
H -2.6770 -1.1346 0.7773	H -2.7057 -1.2176 0.6737	H -2.7675 -1.2432 0.6845	H -2.6867 -1.0752 0.6699
H -0.4986 -2.3105 0.6648	H -0.5890 -2.4262 0.6935	H -0.6212 -2.4804 0.6819	H -0.5258 -2.3427 0.6909
H -0.6650 2.6496 0.7064	H -0.6173 2.4752 0.5504	H -0.6325 2.5134 0.5302	H -0.5252 2.6852 0.6912
H 3.8308 0.3274 0.7112	H 3.7719 0.0569 0.5331	H 3.7794 0.0546 0.4631	H 3.8737 0.1708 0.7240
H 2.0412 2.3270 -0.1084	H 2.0068 2.1431 -0.1311	H 2.0249 2.1231 -0.1558	H 2.1888 2.2530 0.0121
H 2.2631 1.8780 1.5754	H 2.2045 1.5989 1.5917	H 2.2296 1.5710 1.5304	H 2.3574 1.8122 1.7075
H 4.4744 -0.1664 -2.5379	H 3.9839 0.8055 -2.7598	H 3.6368 1.0647 -2.8343	H 3.6714 0.9616 -2.7782
H 3.1445 -2.4525 -0.3751	H 3.0425 -2.4752 -0.5834	H 3.1197 -2.7309 -0.1474	H 3.2186 -2.3907 -0.2631
H 1.8339 -1.9356 -1.2210	H 1.5000 -2.2173 -1.0921	H 1.8932 -2.2908 -1.1085	H 1.9684 -1.8636 -1.1354
H 3.5279 -1.8130 1.9825	H 3.4217 -1.9258 2.0840	H 3.4339 -1.8020 2.1337	H 2.7741 -2.6852 1.9560
H 2.4502 -0.6063 2.7039	H 2.2960 -0.6948 2.7598	H 2.2602 -0.6147 2.7122	H 3.3242 -1.0315 2.4707
H 1.8484 -2.2363 2.3613	H 1.7243 -2.3742 2.4691	H 1.7728 -2.2940 2.5059	H 1.6110 -1.5564 2.7782
H -3.9713 0.8784 0.5746	H -3.9839 0.7822 0.5907	H -3.9734 0.8407 0.2514	H -3.8737 0.8871 0.6634
H -3.3315 1.9376 1.8346	H -3.1749 2.1051 1.4970	H -3.3789 1.7621 1.6527	H -3.0986 2.2535 1.5768
H -3.1762 2.3985 0.1458	H -3.1621 2.0988 -0.3106	H -3.1093 2.3562 -0.0063	H -3.0889 2.2541 -0.2412

Structure Data β -aa-Me*

SSS-*eq*

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -2.7399 0.2892 -0.4703	C -2.6730 0.2478 -0.4685	C -2.7521 0.2246 -0.5278	C -2.3324 1.0591 0.2337
C -2.7775 -0.6789 0.5424	C -2.7095 -0.6576 0.5273	C -2.7632 -0.7376 0.4997	C -1.9413 1.5725 -1.0132
C -1.5852 -1.0300 1.1777	C -1.5636 -0.9755 1.1575	C -1.5759 -1.0399 1.1805	C -0.7332 1.1838 -1.5960
C -0.3898 -0.4269 0.8020	C -0.4132 -0.3976 0.7836	C -0.3995 -0.4104 0.7962	C 0.0704 0.2915 -0.9071
C -0.3699 0.5264 -0.2079	C -0.3889 0.4972 -0.2110	C -0.3931 0.5377 -0.2401	C -0.3194 -0.2202 0.3356
C -1.5471 0.8957 -0.8461	C -1.5229 0.8319 -0.8416	C -1.5690 0.8679 -0.8991	C -1.5209 0.1496 0.9154
C 0.9902 -0.6551 1.3738	C 0.9535 -0.6196 1.3746	C 0.9762 -0.5537 1.3943	C 1.3773 -0.2962 -1.3384
C 1.8998 -0.0193 0.3035	C 1.8708 0.0028 0.2902	C 1.8352 0.0027 0.2342	C 1.9147 -0.8383 -0.0060
C 1.0251 1.0684 -0.4600	C 0.9961 1.0396 -0.4684	C 0.9921 1.1256 -0.4366	C 0.6910 -1.1972 0.8495
C 1.3639 1.1421 -1.9461	C 1.3248 1.1208 -1.9670	C 1.3696 1.2980 -1.9116	C 0.9663 -0.9865 2.3325
C 3.1839 0.5366 0.8486	C 3.1864 0.5246 0.8433	C 3.1846 0.4677 0.7055	C 2.7584 -2.0660 -0.2470
C 1.3196 -2.1135 1.6827	C 1.2844 -2.0858 1.6753	C 1.3588 -1.9823 1.7646	C 2.3006 0.7654 -1.9215
O 4.0254 0.9211 -0.1345	O 3.9015 1.1792 -0.0911	O 4.1385 0.3326 -0.2365	O 3.3267 -2.7003 0.7850
O 3.4695 0.6590 2.0170	O 3.6086 0.3915 1.9665	O 3.4619 0.9372 1.7964	O 2.9206 -2.4762 -1.3716
N 1.1096 2.4278 0.0876	N 1.1161 2.3698 0.1532	N 1.0358 2.4389 0.2147	N 0.2849 -2.5732 0.7511
H -3.6600 0.5684 -0.9709	H -3.6072 0.5219 -0.9869	H -3.6771 0.4872 -1.0374	C -2.8281 2.5461 -1.7250
C -4.0830 -1.3248 0.9304	C -4.0232 -1.2894 0.9238	C -4.0564 -1.3796 0.9108	H -3.2874 1.3763 0.6780
H -1.5989 -1.7780 1.9618	H -1.5658 -1.7032 1.9830	H -1.5795 -1.7530 1.9999	H -0.4280 1.5765 -2.5772
H -1.5412 1.6426 -1.6301	H -1.5173 1.5780 -1.6512	H -1.5795 1.6263 -1.6759	H -1.8285 -0.2623 1.8879
H 1.0890 -0.0821 2.3009	H 0.9954 -0.0139 2.3124	H 1.0417 0.0841 2.2844	H 1.2911 -1.0618 -2.1418
H 2.1687 -0.7828 -0.4308	H 2.1361 -0.8117 -0.4285	H 2.0033 -0.8142 -0.4857	H 2.5535 -0.0837 0.5051
H 0.6770 1.8245 -2.4461	H 0.6880 1.8722 -2.4853	H 0.7258 2.0304 -2.4133	H 0.7764 -1.9344 2.8840
H 2.3792 1.5143 -2.0958	H 2.3799 1.4273 -2.1388	H 2.4041 1.6425 -2.0207	H 2.0278 -0.6825 2.4727
H 1.2773 0.1594 -2.4096	H 1.1700 0.1399 -2.4710	H 1.2710 0.3556 -2.4630	H 0.2962 -0.1889 2.7243
H 0.6378 -2.5153 2.4337	H 0.6391 -2.4944 2.4853	H 0.7301 -2.3581 2.5790	H 1.8811 1.1347 -2.8840
H 2.3341 -2.2110 2.0746	H 2.3403 -2.1917 2.0132	H 2.3999 -2.0313 2.1001	H 3.3063 0.3236 -2.1010

Appendix

H 1.2311 -2.7282 0.7840	H 1.1479 -2.7254 0.7740	H 1.2422 -2.6657 0.9162	H 2.3893 1.6133 -1.2060
H 4.8229 1.2921 0.2791	H 4.7444 1.4524 0.3063	H 4.9572 0.6591 0.1937	H 3.8151 -3.4498 0.4048
H 0.8928 2.4440 1.0783	H 0.9057 2.3099 1.1647	H 0.7590 2.3476 1.1935	H 0.0795 -2.7983 -0.2223
H 2.0397 2.8149 -0.0345	H 2.0854 2.7254 0.0862	H 1.9858 2.8066 0.2246	H 1.0703 -3.1775 0.9928
H -3.9386 -2.0673 1.7151	H -3.9149 -2.0170 1.7589	H -3.8809 -2.3608 1.3641	H -2.3577 2.8405 -2.6898
H -4.7956 -0.5821 1.2972	H -4.7444 -0.5070 1.2531	H -4.5755 -0.7445 1.6350	H -3.8151 2.0716 -1.9236
H -4.5442 -1.8251 0.0751	H -4.4650 -1.8359 0.0596	H -4.7057 -1.5346 0.0428	H -2.9708 3.4498 -1.0911

SSS-ax

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -2.3364 0.0138 -0.8948	C -2.2270 -0.0209 -0.9261	C -2.2821 -0.0070 -0.9699	C -2.0829 1.5811 0.8278
C -2.1862 -0.9273 0.1278	C -2.1270 -0.8900 0.0982	C -2.2040 -0.9371 0.0826	C -1.6918 2.0945 -0.4191
C -0.9288 -1.2885 0.6001	C -0.9349 -1.2152 0.6247	C -0.9711 -1.2961 0.6325	C -0.4837 1.7058 -1.0019
C 0.1932 -0.6991 0.0373	C 0.1725 -0.6621 0.1127	C 0.1855 -0.7222 0.1218	C 0.3199 0.8135 -0.3130
C 0.0655 0.2338 -0.9833	C 0.0863 0.2041 -0.9022	C 0.1132 0.2074 -0.9200	C -0.0699 0.3018 0.9297
C -1.1912 0.5948 -1.4521	C -1.1040 0.5278 -1.4253	C -1.1085 0.5723 -1.4709	C -1.2714 0.6716 1.5095
C 1.6303 -0.9145 0.4575	C 1.5902 -0.8888 0.5660	C 1.6071 -0.9287 0.5866	C 1.6268 0.2258 -0.7443
C 2.4313 -0.1385 -0.6931	C 2.3967 -0.2433 -0.5976	C 2.4169 -0.2897 -0.5905	C 2.1642 -0.3163 0.5881
C 1.4064 0.7709 -1.4253	C 1.4307 0.7214 -1.3362	C 1.4731 0.7060 -1.3215	C 0.9405 -0.6752 1.4436
C 1.5801 2.2564 -1.0923	C 1.6254 2.1947 -0.9580	C 1.7098 2.1601 -0.9296	C 0.4967 -2.1128 1.2066
C 3.0607 -1.1407 -1.6181	C 2.9198 -1.2659 -1.6006	C 2.9236 -1.2743 -1.6266	C 2.9921 0.7374 1.2817
N 1.9073 -2.3403 0.6065	N 1.8667 -2.3279 0.7179	N 1.8914 -2.3328 0.8908	N 2.4986 1.2281 -1.2948
O 4.3187 -1.4565 -1.2430	O 4.1356 -0.9246 -2.0723	O 3.8531 -0.7177 -2.4352	O 3.5579 0.4721 2.4648
O 2.5254 -1.6621 -2.5696	O 2.3644 -2.2720 -1.9705	O 2.6077 -2.4411 -1.7839	O 3.1440 1.8206 0.7689
H 4.6527 -2.1385 -1.8498	H 4.4124 -1.5966 -2.7166	H 4.0925 -1.4402 -3.0543	C -3.3826 2.0128 1.4325
C -3.7119 0.3886 -1.3808	C -3.5680 0.3447 -1.5178	C -3.6108 0.3815 -1.5506	C 1.5088 -0.8218 -1.8435
H -3.0707 -1.3806 0.5593	H -3.0362 -1.3477 0.5219	H -3.1092 -1.3896 0.4827	H -2.3435 2.8099 -0.9421
H -0.8244 -2.0182 1.3928	H -0.8653 -1.9238 1.4652	H -0.9174 -2.0135 1.4470	H -0.1785 2.0985 -1.9831
H -1.2894 1.3230 -2.2507	H -1.1554 1.2419 -2.2622	H -1.1525 1.2946 -2.2818	H -1.5790 0.2597 2.4820
C 1.9027 -0.2487 1.8112	C 1.8490 -0.1901 1.9103	C 1.8383 -0.1539 1.8972	H 2.8143 -1.2050 0.4260
H 3.2319 0.4451 -0.2420	H 3.2811 0.3011 -0.1913	H 3.3163 0.2036 -0.1967	H 1.1167 -0.6157 2.5409
H 2.5700 2.6064 -1.3936	H 2.6441 2.5437 -1.2433	H 2.7331 2.4621 -1.1761	H 1.2700 -2.8099 1.6002
H 0.8377 2.8539 -1.6236	H 0.8895 2.8447 -1.4843	H 1.0264 2.8232 -1.4719	H -0.4675 -2.2942 1.7320
H 1.4542 2.4461 -0.0244	H 1.4997 2.3641 0.1339	H 1.5506 2.3383 0.1372	H 0.3636 -2.2846 0.1150
H 1.6971 -2.8521 -0.2428	H 1.4669 -2.8447 -0.0846	H 1.7525 -2.8826 0.0340	H 2.6310 1.9751 -0.6130
H 2.8789 -2.5050 0.8454	H 2.8822 -2.5225 0.6910	H 2.8760 -2.4505 1.1209	H 3.4301 0.8337 -1.4257
H 1.2529 -0.6847 2.5699	H 1.2144 -0.6236 2.7166	H 1.2306 -0.5592 2.7164	H 4.0355 1.2786 2.7221
H 2.9426 -0.4032 2.1110	H 2.9117 -0.3115 2.2209	H 2.8891 -0.2053 2.2068	H -3.8877 2.7339 0.7516
H 1.7099 0.8236 1.7640	H 1.6288 0.8989 1.8653	H 1.5716 0.9035 1.8089	H -4.0355 1.1238 1.5810
H 1.5191 0.6526 -2.5049	H 1.5294 0.6196 -2.4438	H 1.5700 0.6210 -2.4110	H -3.1907 2.5014 2.4139
H -4.4805 -0.1653 -0.8423	H -4.4124 -0.1979 -1.0373	H -4.4253 -0.2112 -1.1220	H 2.1161 -0.5084 -2.7221
H -3.9033 1.4555 -1.2360	H -3.7576 1.4356 -1.3965	H -3.8154 1.4372 -1.3465	H 1.8821 -1.7996 -1.4650
H -3.8267 0.1738 -2.4467	H -3.5905 0.1014 -2.6046	H -3.6130 0.2163 -2.6329	H 0.4427 -0.9228 -2.1467

SSR-eg

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -2.5961 -0.2124 -0.8864	C -2.5464 -0.2766 -0.8809	C -2.6140 -0.2774 -0.9162	C -2.3324 1.0591 0.4445
C -2.6425 -1.0613 0.2237	C -2.5833 -1.0391 0.2286	C -2.6341 -1.0957 0.2297	C -1.9413 1.5725 -0.8024
C -1.4541 -1.3464 0.9031	C -1.4370 -1.2724 0.8939	C -1.4470 -1.3271 0.9365	C -0.7332 1.1838 -1.3852
C -0.2566 -0.7918 0.4689	C -0.2881 -0.7501 0.4431	C -0.2669 -0.7628 0.4707	C 0.0704 0.2915 -0.6963
C -0.2313 0.0546 -0.6341	C -0.2619 0.0006 -0.6645	C -0.2474 0.0327 -0.6840	C -0.3194 -0.2202 0.5464
C -1.4012 0.3528 -1.3169	C -1.3961 0.2485 -1.3333	C -1.4227 0.2909 -1.3755	C -1.5209 0.1496 1.1262
C 1.1120 -0.9298 1.0997	C 1.0743 -0.9103 1.0648	C 1.1030 -0.8791 1.0861	C 1.3773 -0.2962 -1.1276
C 2.0271 -0.4140 -0.0508	C 1.9805 -0.4418 -0.1026	C 1.9823 -0.4498 -0.1160	C 1.9147 -0.8383 0.2048
C 1.1603 0.5759 -0.9297	C 1.1231 0.5035 -0.9923	C 1.1466 0.5556 -0.9714	C 0.6910 -1.1972 1.0603

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C 1.4913 0.4810 -2.4179	C 1.4799 0.3751 -2.4809	C 1.5184 0.4509 -2.4572	C 0.9663 -0.9865 2.5433
C 3.4043 0.0721 0.3149	C 3.3798 -0.0374 0.3232	C 3.3835 -0.0731 0.2913	C 2.7584 -2.0660 -0.0361
C 1.1534 -0.1945 2.4469	C 1.1991 -0.0788 2.3478	C 1.2201 -0.0339 2.3522	C 1.1781 -1.4128 -2.1441
O 4.0453 0.6213 -0.7423	O 3.7113 1.2166 -0.0319	O 3.8519 1.0742 -0.2275	O 3.3267 -2.7003 0.9958
O 3.9468 -0.0310 1.3891	O 4.1562 -0.7372 0.9288	O 4.1076 -0.7426 1.0134	O 2.9206 -2.4762 -1.1607
N 1.2299 1.9912 -0.5395	N 1.2151 1.9174 -0.5902	N 1.2109 1.9681 -0.5864	N 0.2849 -2.5732 0.9620
H -3.5150 0.0131 -1.4145	H -3.4810 -0.0738 -1.4303	H -3.5377 -0.0728 -1.4539	C -2.8281 2.5461 -1.5141
C -3.9519 -1.6574 0.6652	C -3.8969 -1.6101 0.7084	C -3.9349 -1.6662 0.7145	H -3.2874 1.3763 0.8888
H -1.4724 -1.9972 1.7712	H -1.4279 -1.8881 1.8068	H -1.4454 -1.9371 1.8355	H -0.4280 1.5765 -2.3663
H -1.3918 1.0168 -2.1719	H -1.3880 0.8774 -2.2371	H -1.4249 0.9353 -2.2493	H -1.8285 -0.2623 2.0987
H 1.3476 -1.9786 1.2936	H 1.2644 -1.9848 1.3044	H 1.3132 -1.9277 1.3311	H 2.0570 0.4243 -1.6352
H 2.2213 -1.2802 -0.6920	H 2.1607 -1.3569 -0.7265	H 2.1259 -1.3490 -0.7396	H 2.5535 -0.0837 0.7159
H 0.7732 1.0680 -2.9896	H 0.8706 1.0630 -3.1092	H 0.8945 1.1038 -3.0797	H 0.7764 -1.9344 3.0948
H 2.4882 0.8691 -2.6231	H 2.5494 0.6341 -2.6520	H 2.5629 0.7335 -2.6296	H 2.0278 -0.6825 2.6835
H 1.4396 -0.5544 -2.7569	H 1.3137 -0.6597 -2.8571	H 1.3828 -0.5699 -2.8340	H 0.2962 -0.1889 2.9352
H 0.4001 -0.6271 3.1083	H 0.4558 -0.4084 3.1092	H 0.4970 -0.3668 3.1056	H 0.7889 -0.9846 -3.0948
H 2.1265 -0.2878 2.9221	H 2.2066 -0.1887 2.8079	H 2.2150 -0.1243 2.7986	H 2.1516 -1.9170 -2.3360
H 0.9182 0.8669 2.3370	H 1.0270 1.0037 2.1594	H 1.0272 1.0268 2.1701	H 0.4486 -2.1523 -1.7444
H 4.9290 0.8921 -0.4419	H 4.6259 1.3809 0.2492	H 4.7730 1.1254 0.1091	H 3.8151 -3.4498 0.6157
H 0.9713 2.1198 0.4319	H 1.0569 2.0182 0.4258	H 0.8328 2.0946 0.3495	H 0.0795 -2.7983 -0.0114
H 2.1697 2.3524 -0.6634	H 2.1657 2.2879 -0.7594	H 2.1952 2.2436 -0.5418	H 1.0703 -3.1775 1.2036
H -4.1042 -1.5233 1.7389	H -3.7916 -2.1968 1.6482	H -3.7707 -2.4833 1.4248	H -2.3577 2.8405 -2.4789
H -4.7881 -1.1925 0.1433	H -4.6259 -0.7915 0.9060	H -4.5231 -0.8881 1.2107	H -3.8151 2.0716 -1.7128
H -3.9852 -2.7310 0.4569	H -4.3273 -2.2879 -0.0637	H -4.5107 -2.0723 -0.1236	H -2.9708 3.4498 -0.8802

SSR-ax

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C -2.1760 0.5025 -0.5766	C -2.0658 0.4423 -0.5899	C -2.1325 0.4448 -0.6279	C -2.0829 1.0759 0.4382
C -2.0534 -0.4164 0.4678	C -1.9507 -0.4050 0.4508	C -2.0314 -0.4255 0.4738	C -1.6918 1.5893 -0.8087
C -0.8071 -0.7773 0.9747	C -0.7550 -0.6832 0.9953	C -0.7900 -0.7311 1.0358	C -0.4837 1.2006 -1.3915
C 0.3284 -0.2110 0.4200	C 0.3381 -0.1046 0.4816	C 0.3501 -0.1603 0.4881	C 0.3199 0.3083 -0.7026
C 0.2287 0.6929 -0.6319	C 0.2400 0.7415 -0.5506	C 0.2581 0.7148 -0.6051	C -0.0699 -0.2034 0.5401
C -1.0140 1.0546 -1.1326	C -0.9558 1.0183 -1.0884	C -0.9749 1.0339 -1.1575	C -1.2714 0.1664 1.1199
C 1.7597 -0.3668 0.8921	C 1.7610 -0.2760 0.9456	C 1.7841 -0.3182 0.9422	C 1.6268 -0.2794 -1.1339
C 2.5779 0.2241 -0.3235	C 2.5335 0.2410 -0.3005	C 2.5486 0.1792 -0.3287	C 2.1642 -0.8215 0.1985
C 1.5916 1.2105 -1.0404	C 1.5872 1.2796 -0.9604	C 1.6197 1.2438 -0.9766	C 0.9405 -1.1804 1.0540
C 1.7927 1.3894 -2.5432	C 1.7729 1.4791 -2.4680	C 1.8270 1.5078 -2.4629	C 1.2158 -0.9696 2.5370
C 3.0656 -0.8333 -1.2773	C 2.8763 -0.8874 -1.2659	C 2.8886 -0.9354 -1.2998	C 2.9921 0.2323 0.8921
N 2.0478 -1.7371 1.2961	N 2.0579 -1.6860 1.2488	N 2.0873 -1.6791 1.3907	N 2.4986 0.7230 -1.6844
O 4.3087 -0.5667 -1.7311	O 4.0831 -0.7084 -1.8397	O 3.9860 -0.6497 -2.0346	O 3.5579 -0.0330 2.0752
O 2.4529 -1.8129 -1.6374	O 2.2007 -1.8550 -1.5214	O 2.3168 -2.0000 -1.4601	O 3.1440 1.3154 0.3793
H 4.5547 -1.2670 -2.3595	H 4.2438 -1.4485 -2.4476	H 4.0900 -1.4323 -2.6170	C -3.3826 1.5076 1.0429
C -3.5386 0.8976 -1.0827	C -3.4118 0.7561 -1.2001	C -3.4787 0.7930 -1.1931	C 1.5088 -1.3269 -2.2331
H -2.9505 -0.8501 0.8940	H -2.8498 -0.8827 0.8737	H -2.9268 -0.8701 0.9025	H -2.3435 2.3048 -1.3317
H -0.7252 -1.4847 1.7899	H -0.6714 -1.3733 1.8497	H -0.7182 -1.4023 1.8867	H -0.1785 1.5933 -2.3727
H -1.0944 1.7645 -1.9487	H -1.0322 1.7180 -1.9347	H -1.0476 1.7311 -1.9868	H -1.5790 -0.2455 2.0924
C 1.9845 0.5414 2.1089	C 2.0326 0.5870 2.1873	C 2.0515 0.6146 2.1378	H 2.8143 -1.7102 0.0364
H 3.4527 0.7658 0.0319	H 3.4946 0.7159 0.0092	H 3.5061 0.6287 -0.0314	H 0.6262 -2.2454 0.9779
H 2.8081 1.7228 -2.7658	H 2.8223 1.7733 -2.6978	H 2.8629 1.7991 -2.6643	H 1.9695 -1.7114 2.8840
H 1.1044 2.1415 -2.9301	H 1.1165 2.2916 -2.8539	H 1.1857 2.3285 -2.8031	H 0.2721 -1.1038 3.1117
H 1.6051 0.4590 -3.0822	H 1.5337 0.5569 -3.0436	H 1.5823 0.6347 -3.0772	H 1.6067 0.0598 2.6989
H 1.9652 -2.3659 0.5061	H 1.6277 -2.2916 0.5277	H 1.8597 -2.3271 0.6286	H 2.6310 1.4700 -1.0026
H 2.9890 -1.8193 1.6643	H 3.0736 -1.8763 1.2006	H 3.0857 -1.7830 1.5552	H 3.4301 0.3286 -1.8153
H 3.0152 0.4567 2.4634	H 3.0960 0.4973 2.5060	H 3.1051 0.5782 2.4392	H 4.0355 0.7735 2.3325
H 1.7888 1.5846 1.8577	H 1.8228 1.6638 1.9982	H 1.8119 1.6595 1.9084	H -3.8877 2.2288 0.3620

Appendix

H 1.3117 0.2481 2.9149	H 1.3969 0.2653 3.0436	H 1.4449 0.3353 3.0085	H -4.0355 0.6186 1.1914
H 1.7443 2.1871 -0.5698	H 1.7150 2.2662 -0.4491	H 1.7559 2.1947 -0.4436	H -3.1907 1.9963 2.0243
H -4.3230 0.3422 -0.5691	H -4.2438 0.1959 -0.7181	H -4.2022 -0.0078 -1.0080	H 2.1161 -1.0136 -3.1117
H -3.7225 1.9636 -0.9189	H -3.6364 1.8425 -1.0994	H -3.8493 1.7152 -0.7353	H 1.8821 -2.3048 -1.8546
H -3.6360 0.7067 -2.1546	H -3.4164 0.4936 -2.2826	H -3.4193 0.9322 -2.2777	H 0.4427 -1.4280 -2.5363

RSR-eg

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C 2.6738 0.2084 1.0688	C 2.6017 0.2413 1.0742	C 2.6813 0.1865 1.0599	C 2.0975 -1.6400 -0.9184
C 2.5914 1.3567 0.2753	C 2.5513 1.2715 0.2078	C 2.6072 1.3327 0.2456	C 1.7064 -2.1534 0.3285
C 1.3928 1.7467 -0.3154	C 1.4127 1.5971 -0.4259	C 1.4085 1.7014 -0.3699	C 0.4983 -1.7647 0.9113
C 0.2564 0.9829 -0.0989	C 0.3087 0.8782 -0.1837	C 0.2780 0.9255 -0.1533	C -0.3053 -0.8724 0.2224
C 0.3114 -0.1403 0.7208	C 0.3462 -0.1521 0.6701	C 0.3438 -0.2100 0.6643	C 0.0845 -0.3607 -1.0203
C 1.5112 -0.5337 1.2978	C 1.4841 -0.4740 1.3003	C 1.5317 -0.5850 1.2776	C 1.2860 -0.7305 -1.6001
C -1.1007 1.1344 -0.7633	C -1.0560 1.0719 -0.7959	C -1.1129 1.1226 -0.7221	C -1.6122 -0.2847 0.6537
C -1.9698 0.2454 0.1896	C -1.9133 0.2454 0.2024	C -1.9277 0.2301 0.2666	C -2.1496 0.2574 -0.6787
C -1.0312 -0.8272 0.8325	C -0.9885 -0.8450 0.8061	C -0.9885 -0.9092 0.7490	C -0.9259 0.6163 -1.5342
C -1.0438 -2.2171 0.1904	C -0.9541 -2.2095 0.1048	C -0.9776 -2.2070 -0.0552	C -0.4822 2.0538 -1.2972
C -3.2827 -0.2693 -0.3267	C -3.3219 -0.1024 -0.2340	C -3.2906 -0.1981 -0.2105	C -2.9933 1.4851 -0.4377
N -1.5387 2.5306 -0.8233	N -1.4653 2.4876 -0.7756	N -1.5165 2.5338 -0.6518	N -2.4841 -1.2870 1.2042
O -3.9633 -0.9117 0.6471	O -3.8470 -1.1183 0.4763	O -3.8311 -1.1782 0.5405	O -3.5617 2.1193 -1.4697
O -3.7416 -0.1282 -1.4355	O -3.9743 0.4762 -1.0688	O -3.9641 0.3313 -1.0786	O -3.1556 1.8953 0.6869
H -4.8076 -1.2148 0.2730	H -4.7688 -1.2410 0.1960	H -4.7334 -1.2934 0.1731	C 3.3971 -2.0717 -1.5230
C 3.9905 -0.2156 1.6629	C 3.8856 -0.1255 1.7809	C 3.9709 -0.1716 1.7394	C -1.4943 0.7628 1.7530
H 3.4887 1.9389 0.1016	H 3.4584 1.8653 0.0069	H 3.4917 1.9460 0.0853	H 2.3580 -2.8689 0.8515
H 1.3499 2.6251 -0.9469	H 1.3840 2.4393 -1.1352	H 1.3601 2.5834 -1.0026	H 0.1931 -2.1574 1.8925
H 1.5530 -1.4226 1.9188	H 1.4967 -1.3233 2.0014	H 1.5669 -1.4643 1.9150	H 1.5935 -0.3186 -2.5726
C -0.9769 0.6428 -2.2154	C -1.0573 0.5756 -2.2498	C -1.1321 0.7005 -2.1976	H -2.7885 -0.4972 -1.1898
H -2.2516 0.9073 1.0148	H -2.1161 0.9450 1.0589	H -2.1388 0.8415 1.1619	H -1.1022 0.5568 -2.6315
H -2.0204 -2.6897 0.3042	H -1.9122 -2.7629 0.2067	H -1.9430 -2.7182 -0.0133	H -1.2554 2.7510 -1.6908
H -0.3090 -2.8551 0.6839	H -0.1690 -2.8611 0.5522	H -0.2328 -2.9011 0.3522	H 0.4821 2.2353 -1.8226
H -0.7916 -2.1799 -0.8707	H -0.7284 -2.1261 -0.9789	H -0.7222 -2.0539 -1.1057	H -0.3491 2.2257 -0.2056
H -1.6935 2.9011 0.1086	H -1.4236 2.8611 0.1882	H -1.4729 2.8560 0.3136	H -2.6165 -2.0341 0.5224
H -2.4181 2.5966 -1.3251	H -2.4541 2.5778 -1.0684	H -2.4915 2.6149 -0.9441	H -3.4156 -0.8927 1.3351
H -1.9185 0.7514 -2.7488	H -2.0516 0.6941 -2.7325	H -2.1188 0.8305 -2.6532	H -4.0500 2.8689 -1.0895
H 4.8208 0.3071 1.1883	H 4.7688 0.4129 1.3703	H 4.8300 0.1376 1.1349	H 3.9023 -2.7929 -0.8422
H 4.1485 -1.2894 1.5365	H 4.0910 -1.2158 1.6831	H 4.0486 -1.2543 1.8838	H 4.0500 -1.1827 -1.6716
H 4.0304 0.0013 2.7348	H 3.8083 0.1229 2.8639	H 4.0284 0.3209 2.7149	H 3.2053 -2.5604 -2.5045
H -0.6704 -0.4017 -2.2608	H -0.7667 -0.4921 -2.3374	H -0.8376 -0.3417 -2.3403	H -2.1015 0.4495 2.6315
H -0.2149 1.2427 -2.7137	H -0.3348 1.1602 -2.8639	H -0.4303 1.3029 -2.7891	H -1.8676 1.7407 1.3744
H -1.3074 -0.9514 1.8817	H -1.2272 -1.0026 1.8864	H -1.2173 -1.1404 1.7968	H -0.4282 0.8639 2.0561

RSR-ax

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C 2.2949 0.3449 0.2319	C 2.1895 -0.2558 0.2771	C 2.2771 -0.2969 0.2836	C 2.3179 -0.9230 -0.2337
C 2.2424 0.5955 0.8049	C 2.1178 0.6257 -0.7386	C 2.1994 0.6506 -0.7552	C 1.9268 -1.4364 1.0132
C 0.9950 0.9619 1.3137	C 0.9166 0.8976 -1.2816	C 0.9545 0.9655 -1.3164	C 0.7187 -1.0477 1.5960
C 0.1626 0.3963 0.7909	C -0.1809 0.2939 -0.8048	C -0.1855 0.3574 -0.8090	C -0.0849 -0.1554 0.9071
C 0.0931 0.5338 0.2371	C -0.0957 -0.5770 0.2091	C -0.1017 -0.5784 0.2361	C 0.3049 0.3563 -0.3356
C 1.1378 0.9144 0.7539	C 1.0925 -0.8642 0.7564	C 1.1284 -0.9174 0.7804	C 1.5064 -0.0135 -0.9154
C 1.5894 0.6089 1.2437	C -1.6044 0.4674 -1.2689	C -1.6118 0.5098 -1.2706	C -1.3918 0.4323 1.3384
C 2.4293 0.0198 0.0935	C -2.3714 -0.0537 -0.0252	C -2.3870 -0.0397 -0.0442	C -1.9292 0.9744 0.0060
C 1.4608 1.0838 0.5956	C -1.4430 -1.1455 0.5773	C -1.4681 -1.1314 0.5850	C -0.7055 1.3333 -0.8495
C 1.6821 1.2483 2.0968	C -1.6537 -1.3989 2.0753	C -1.6866 -1.3723 2.0819	C -0.9808 1.1226 -2.3325
C 2.9414 0.9843 0.9055	C -2.6555 1.0646 0.9718	C -2.7396 1.0494 0.9515	C -2.7571 -0.0793 -0.6876

C 1.9327 2.0375 1.6628	C -1.9868 1.8735 -1.7405	C -2.0130 1.9021 -1.7405	C -2.3152 -0.6293 1.9215
O 4.2023 0.6891 1.2983	O -3.9677 1.1498 1.2680	O -4.0572 1.0567 1.2470	O -3.3229 0.1860 -1.8707
O 2.3409 1.9323 1.3563	O -1.8536 1.8155 1.4741	O -1.9985 1.8598 1.4784	O -2.9091 -1.1625 -0.1747
N 1.5887 2.4161 0.0120	N -1.6347 -2.4093 -0.1591	N -1.6386 -2.4302 -0.0933	N -0.2994 2.7093 -0.7511
H 3.2584 0.6373 0.6341	H 3.1702 -0.4881 0.7251	H 3.2444 -0.5604 0.7069	C 2.8136 -2.4100 1.7250
C 3.5102 1.1912 1.3613	C 3.3755 1.2883 -1.2502	C 3.4565 1.2768 -1.2851	H 3.2729 -1.2402 -0.6780
H 0.9381 1.6864 2.1173	H 0.8291 1.6110 -2.1149	H 0.8844 1.6752 -2.1356	H 0.4135 -1.4404 2.5772
H 1.2038 1.6481 1.5475	H 1.1711 -1.5881 1.5821	H 1.2054 -1.6589 1.5694	H 1.8140 0.3984 -1.8879
H 1.7571 0.0272 2.1205	H -1.7603 -0.2463 -2.1165	H -1.7556 -0.1838 -2.1105	H -1.3056 1.1979 2.1418
H 3.2949 0.5488 0.4863	H -3.3493 -0.4940 -0.3348	H -3.3357 -0.4759 -0.3880	H -2.5794 1.8631 0.1681
H 0.9755 1.9809 2.4853	H -1.0275 -2.2468 2.4339	H -1.0643 -2.1971 2.4513	H -0.7910 2.0705 -2.8840
H 2.6910 1.6136 2.3027	H -2.7162 -1.6642 2.2788	H -2.7310 -1.6301 2.2930	H -2.0423 0.8186 -2.4727
H 1.5332 0.3103 2.6293	H -1.3912 -0.5144 2.6967	H -1.4219 -0.4970 2.6850	H -0.3107 0.3250 -2.7243
H 1.3413 2.3198 2.5352	H -1.4773 2.1312 -2.6967	H -1.5071 2.1506 -2.6802	H -1.8956 -0.9986 2.8840
H 2.9861 2.1241 1.9396	H -3.0838 1.9371 -1.9239	H -3.0915 1.9530 -1.9229	H -3.3208 -0.1874 2.1010
H 1.7243 2.7466 0.8641	H -1.7138 2.6581 -1.0009	H -1.7445 2.6814 -1.0203	H -2.4038 -1.4772 1.2060
H 4.4676 1.3424 1.9674	H -4.0880 1.8702 1.9080	H -4.1477 1.7912 1.8915	H -3.8006 -0.6205 -2.1280
H 1.4016 2.3800 1.0086	H -1.3224 -2.2924 -1.1386	H -1.3997 -2.3386 -1.0803	H -0.0940 2.9344 0.2223
H 2.5318 2.7727 0.1086	H -2.6373 -2.6581 -0.2147	H -2.6137 -2.7204 -0.0504	H -1.0848 3.3137 -0.9928
H 3.2954 1.9209 2.1418	H 3.1773 2.0027 -2.0802	H 3.2354 2.1547 -1.9011	H 2.3432 -2.7044 2.6898
H 4.1509 0.4180 1.7929	H 4.0880 0.5216 -1.6312	H 4.0056 0.5544 -1.8969	H 3.8006 -1.9355 1.9236
H 4.0867 1.6935 0.5808	H 3.8729 1.8572 -0.4318	H 4.0952 1.6082 -0.4599	H 2.9562 -3.3137 1.0911

RSS-eg

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C 2.5758 0.4407 0.9448	C 2.5053 -0.4300 0.9184	C 2.5665 -0.4420 0.9541	C 2.0975 -1.6400 -0.4382
C 2.5465 0.4087 0.1627	C 2.4832 0.4040 -0.1391	C 2.5568 0.4203 -0.1585	C 1.7064 -2.1534 0.8087
C 1.3534 0.7040 0.8187	C 1.3403 0.6754 -0.7900	C 1.3665 0.7169 -0.8260	C 0.4983 -1.7647 1.3915
C 0.1803 0.1368 0.3537	C 0.2056 0.1047 -0.3659	C 0.1853 0.1460 -0.3723	C -0.3053 -0.8724 0.7026
C 0.1879 0.7211 0.7440	C 0.2108 -0.7284 0.6821	C 0.1865 -0.7184 0.7310	C 0.0845 -0.3607 -0.5401
C 1.3776 1.0098 1.3960	C 1.3550 -0.9987 1.3255	C 1.3694 -1.0303 1.3880	C 1.2860 -0.7305 -1.1199
C 1.2076 0.2939 0.9371	C -1.1700 0.2605 -0.9610	C -1.2040 0.2886 -0.9557	C -1.6122 -0.2847 1.1339
C 2.0786 0.2007 0.2921	C -2.0390 -0.2298 0.2254	C -2.0575 -0.1788 0.2610	C -2.1496 0.2574 -0.1985
C 1.2085 1.2337 1.0366	C -1.1691 -1.2605 0.9865	C -1.2011 -1.2411 0.9995	C -0.9259 0.6163 -1.0540
C 1.5451 1.3944 2.5153	C -1.4973 -1.3458 2.4812	C -1.5472 -1.3382 2.4826	C -1.2012 0.4055 -2.5370
C 3.4591 0.5950 0.1426	C -3.4473 -0.6140 -0.1822	C -3.4332 -0.6321 -0.1535	C -2.9933 1.4851 0.0425
N 1.4639 1.6629 1.3600	N -1.4756 1.6638 -1.2844	N -1.4887 1.6748 -1.3399	N -2.4841 -1.2870 1.6844
O 3.6931 1.9195 0.1231	O -3.7539 -1.8926 0.1024	O -3.7446 -1.8965 0.1752	O -3.5617 2.1193 -0.9895
O 4.2925 0.1990 0.5252	O -4.2361 0.1239 -0.7222	O -4.2522 0.0617 -0.7382	O -3.1556 1.8953 1.1670
C 1.3651 0.6265 2.1541	C -1.3094 -0.6132 -2.2164	C -1.3290 -0.6070 -2.1985	C -1.4943 0.7628 2.2331
C 3.8790 0.7441 1.6342	C 3.7924 -0.7370 1.6477	C 3.8615 -0.7844 1.6311	C 3.3971 -2.0717 -1.0429
H 3.4725 0.8495 0.5126	H 3.4170 0.8746 -0.4886	H 3.4848 0.8619 -0.5150	H 2.3580 -2.8689 1.3317
H 1.3410 1.3716 1.6706	H 1.3328 1.3516 -1.6595	H 1.3655 1.3777 -1.6882	H 0.1931 -2.1574 2.3727
H 1.3878 1.6676 2.2580	H 1.3554 -1.6873 2.1843	H 1.3717 -1.7225 2.2248	H 1.5935 -0.3186 -2.0924
H 2.2040 0.6749 0.9334	H -2.1736 0.6421 0.9160	H -2.2129 0.6649 0.9530	H -2.7885 -0.4972 -0.7096
H 1.3261 2.2072 0.5513	H -1.2440 -2.2692 0.5144	H -1.2946 -2.2299 0.5330	H -0.6116 1.6813 -0.9779
H 0.8810 2.1180 2.9900	H -0.8713 -2.1106 2.9937	H -0.9274 -2.0911 2.9811	H -1.9550 1.1473 -2.8840
H 2.5684 1.7538 2.6456	H -2.5610 -1.6344 2.6388	H -2.5935 -1.6255 2.6251	H -0.2576 0.5397 -3.1117
H 1.4387 0.4436 3.0424	H -1.3302 -0.3700 2.9911	H -1.3849 -0.3848 2.9982	H -1.5921 -0.6239 -2.6989
H 1.3531 2.3012 0.5791	H -1.2637 2.2692 -0.4728	H -1.3496 2.2845 -0.5347	H -2.6165 -2.0341 1.0026
H 2.4192 1.7555 1.6904	H -2.4890 1.7827 -1.4596	H -2.4780 1.7554 -1.5849	H -3.4156 -0.8927 1.8153
H 4.5910 2.0686 0.4649	H -4.6615 -2.0632 -0.1981	H -4.6590 -2.0169 -0.1599	H -4.0500 2.8689 -0.6094
H 2.3587 0.5241 2.5967	H -2.3210 -0.5165 -2.6704	H -2.3234 -0.5369 -2.6533	H -2.1015 0.4495 3.1117
H 1.2118 1.6704 1.8784	H -1.1386 -1.6896 -1.9888	H -1.1474 -1.6615 -1.9599	H -1.8676 1.7407 1.8546
H 0.6261 0.3566 2.9080	H -0.5714 -0.3102 -2.9937	H -0.6007 -0.3238 -2.9689	H -0.4282 0.8639 2.5363

Appendix

H 4.6978 0.1731 1.1976	H 4.6615 -0.1720 1.2428	H 4.5882 0.0275 1.5220	H 3.9023 -2.7929 -0.3620
H 4.1277 1.8057 1.5487	H 4.0304 -1.8222 1.5675	H 4.2822 -1.6955 1.1946	H 4.0500 -1.1827 -1.1914
H 3.8295 0.5018 2.6992	H 3.6985 -0.4751 2.7262	H 3.7091 -0.9413 2.7042	H 3.2053 -2.5604 -2.0243

RSS-ax

B3LYP/def2-TZVP	MM2	MMFF9	unoptimized
C 2.3860 0.1220 0.5755	C 2.2698 0.0020 0.6197	C 2.3648 -0.0465 0.6241	C 2.3179 -0.9230 -0.4445
C 2.3436 0.8199 0.4628	C 2.2088 0.8808 -0.3991	C 2.2966 0.8838 -0.4301	C 1.9268 -1.4364 0.8024
C 1.1011 1.2155 0.9562	C 1.0075 1.1818 -0.9258	C 1.0507 1.2272 -0.9703	C 0.7187 -1.0477 1.3852
C 0.0658 0.6780 0.4186	C -0.0975 0.6089 -0.4295	C -0.0952 0.6421 -0.4464	C -0.0849 -0.1554 0.6963
C 0.0055 0.2505 0.6074	C -0.0259 -0.2575 0.5876	C -0.0244 -0.2730 0.6117	C 0.3049 0.3563 -0.5464
C 1.2238 0.6611 1.1115	C 1.1632 -0.5738 1.1174	C 1.2068 -0.6313 1.1442	C 1.5064 -0.0135 -1.1262
C 1.4854 0.9495 0.8522	C -1.5111 0.8331 -0.8954	C -1.5125 0.8729 -0.8914	C -1.3918 0.4323 1.1276
C 2.3487 0.2300 0.2176	C -2.3124 0.2459 0.2967	C -2.3214 0.2311 0.2688	C -1.9292 0.9744 -0.2048
C 1.3712 0.7524 1.0250	C -1.3784 -0.7957 0.9781	C -1.3920 -0.7967 0.9854	C -0.7055 1.3333 -1.0603
C 1.5566 0.7108 2.5409	C -1.5852 -0.9342 2.4928	C -1.6037 -0.9154 2.4994	C -0.9808 1.1226 -2.5433
C 3.0463 1.2004 1.1316	C -2.7089 1.3661 1.2542	C -2.8063 1.3240 1.2006	C -2.7571 -0.0793 -0.8985
C 1.7605 0.4537 2.2768	C -1.7862 0.1237 -2.2275	C -1.8032 0.2364 -2.2476	C -1.1926 1.5489 2.1441
O 4.1784 0.6738 1.6483	O -3.9748 1.2236 1.6954	O -4.0254 1.0504 1.7126	O -3.3229 0.1860 -2.0816
O 2.6701 2.3138 1.4183	O -2.0230 2.2941 1.6111	O -2.2370 2.3558 1.5099	O -2.9091 -1.1625 -0.3856
N 1.5061 2.1523 0.6077	N -1.5520 -2.1178 0.3476	N -1.5489 -2.1509 0.4239	N -0.2994 2.7093 -0.9620
C 3.6205 1.3816 1.0323	C 3.4759 1.5067 -0.9324	C 3.5592 1.4954 -0.9621	C 2.8136 -2.4100 1.5141
H 3.3471 0.4371 0.9663	H 3.2500 -0.2561 1.0544	H 3.3300 -0.3235 1.0422	H 3.2729 -1.2402 -0.8888
H 1.0491 1.9409 1.7604	H 0.9171 1.8941 -1.7605	H 0.9760 1.9425 -1.7840	H 0.4135 -1.4404 2.3663
H 1.2790 1.3948 1.9060	H 1.2351 -1.2977 1.9438	H 1.2769 -1.3629 1.9430	H 1.8140 0.3984 -2.0987
H 1.6880 2.0215 0.8173	H -1.7121 1.9234 -1.0243	H -1.7038 1.9506 -0.9645	H -2.0715 -0.2882 1.6352
H 3.1210 0.3834 0.2421	H -3.2533 -0.2308 -0.0656	H -3.2254 -0.2500 -0.1301	H -2.5794 1.8631 -0.0428
H 0.8062 1.3454 3.0100	H -0.9429 -1.7373 2.9192	H -0.9536 -1.6821 2.9392	H -0.7910 2.0705 -3.0948
H 2.5415 1.0847 2.8263	H -2.6409 -1.2067 2.7201	H -2.6373 -1.1925 2.7367	H -2.0423 0.8186 -2.6835
H 1.4390 0.3007 2.9291	H -1.3389 0.0032 3.0395	H -1.3705 0.0202 3.0205	H -0.3107 0.3250 -2.9352
H 1.0948 0.9458 2.9875	H -1.1442 0.5355 -3.0395	H -1.1994 0.7043 -3.0331	H -0.8035 1.1207 3.0948
H 2.7897 0.6702 2.5719	H -2.8474 0.2623 -2.5366	H -2.8570 0.3669 -2.5146	H -2.1661 2.0531 2.3360
H 1.5937 0.6234 2.3580	H -1.5907 -0.9696 -2.1674	H -1.5797 -0.8346 -2.2631	H -0.4631 2.2884 1.7444
H 4.5609 1.3235 2.2616	H -4.1680 1.9548 2.3048	H -4.2149 1.8190 2.2926	H -3.8006 -0.6205 -2.3388
H 1.3571 2.2474 0.3910	H -1.0421 -2.1417 -0.5520	H -1.2799 -2.1520 -0.5579	H -0.0940 2.9344 0.0114
H 2.4400 2.4946 0.8111	H -2.5432 -2.2941 0.1103	H -2.5256 -2.4338 0.4658	H -1.0848 3.3137 -1.2036
H 3.4177 2.1074 1.8195	H 3.2854 2.2113 -1.7725	H 3.3550 2.1860 -1.7869	H 2.3432 -2.7044 2.4789
H 4.2412 0.5891 1.4581	H 4.1680 0.7185 -1.3073	H 4.2280 0.7139 -1.3370	H 3.8006 -1.9355 1.7128
H 4.2119 1.8789 0.2597	H 3.9930 2.0774 -0.1277	H 4.0678 2.0568 -0.1717	H 2.9562 -3.3137 0.8802

RDC Data β -aa-H*

B3LYP/def2-TZVP

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{RS-}ax)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}ax)/\text{Hz}$	
		(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)
1-CH ₃	-6.5	-0.1	-0.1	1.3	-0.5	-7.9	-6.2	-5.1	-5.8
C2-H2	16.1	15.6	15.6	15.5	15.6	16.0	16.2	15.9	15.9
C3-H3(I)	4.5	5.2	5.1	5.4	4.3	4.6	4.4	5.1	5.2
C3-H3(II)	10.2	10.9	10.9	10.0	10.9	10.2	10.1	10.8	10.7
C4-H4	12.7	12.1	12.0	13.3	13.1	12.6	12.3	12.3	12.4
5-CH ₃	6.6	6.5	6.6	0.0	0.6	6.1	5.8	0.8	0.2
C6-H6	-11.4	-11.9	-11.7	-11.2	-11.0	-11.4	-11.3	-11.1	-11.4
Q factor		0.25	0.22	0.37	0.31	0.06	0.04	0.24	0.22

MM2

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{RS-}ax)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}ax)/\text{Hz}$	
		(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)
1-CH ₃	-6.5	1.0	1.1	1.3	-0.4	-7.9	-6.7	-5.6	-5.3
C2-H2	16.1	14.9	14.8	15.3	15.4	15.9	16.1	15.7	15.8
C3-H3(I)	4.5	5.6	5.6	5.4	4.2	4.6	4.4	5.3	4.9
C3-H3(II)	10.2	11.4	11.3	9.9	11.0	10.2	10.2	10.6	11.0
C4-H4	12.7	11.7	11.6	13.4	13.1	12.4	12.7	12.6	12.7
5-CH ₃	6.6	7.6	7.6	-2.0	-1.1	7.1	7.0	-1.9	-1.2
C6-H6	-11.4	-12.0	-11.9	-10.8	-10.8	-11.4	-11.5	-10.7	-10.9
Q factor		0.29	0.29	0.43	0.36	0.06	0.03	0.32	0.27

MMFF9

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{RS-}ax)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}ax)/\text{Hz}$	
		(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)
1-CH ₃	-6.5	0.8	0.8	1.5	-0.2	-8.3	-7.4	-4.9	-4.7
C2-H2	16.1	15.2	15.2	15.5	15.6	15.8	16.0	15.8	15.9
C3-H3(I)	4.5	5.5	5.4	5.4	4.3	4.6	4.5	5.3	5.0
C3-H3(II)	10.2	11.1	11.2	9.9	11.0	10.3	10.3	10.7	11.0
C4-H4	12.7	11.9	11.8	13.4	13.2	12.4	12.6	12.7	12.8
5-CH ₃	6.6	6.9	6.6	-0.4	0.5	6.8	6.6	-0.8	0.0
C6-H6	-11.4	-11.9	-11.8	-11.0	-10.9	-11.4	-11.5	-10.8	-11.0
Q factor		0.28	0.27	0.39	0.32	0.07	0.04	0.28	0.25

Unoptimized

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{RS-}ax)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}eq)/\text{Hz}$		$D_{\text{Calc}}(\text{SS-}ax)/\text{Hz}$	
		(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)	(-i)	(-ii)
1-CH ₃	-6.5	0.8	0.2	1.9	0.1	-7.5	-6.1	-7.7	-6.8
C2-H2	16.1	15.4	15.5	15.3	15.4	16.0	16.2	15.6	15.8
C3-H3(I)	4.5	5.4	5.1	5.4	4.3	4.5	4.4	5.1	4.9
C3-H3(II)	10.2	10.9	11.1	9.9	11.0	10.2	10.1	10.6	10.9
C4-H4	12.7	11.9	11.8	13.3	13.1	12.5	12.8	12.7	12.8
5-CH ₃	6.6	6.5	6.2	-1.0	-0.5	6.7	6.5	-1.0	-0.5
C6-H6	-11.4	-12.0	-11.8	-11.1	-10.9	-11.3	-11.5	-10.7	-10.9
Q factor		0.27	0.25	0.42	0.36	0.05	0.03	0.28	0.26

RDC Data β -aa-Me*

B3LYP/def2-TZVP

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{SSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	1.4	-8.3	-0.0	-4.2
C2-H2	1.1	4.3	1.1	2.6	1.3
C3-H3	13.5	11.5	13.3	13.4	13.7
3-CH ₃	-8.8	1.4	-8.3	0.6	-7.2
C4-H4	-4.1	-1.5	-4.1	-2.4	-4.2
5-CH ₃	10.2	4.9	11.8	5.3	12.7
C6-H6	12.0	12.9	11.8	12.5	11.6
Q factor		0.66	0.08	0.57	0.23
Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	-11.1	0.8	-11.1	-2.8
C2-H2	1.1	1.3	0.5	2.4	0.8
C3-H3	13.5	13.2	14.0	13.0	12.8
3-CH ₃	-8.8	-7.0	-1.1	-6.5	-8.9
C4-H4	-4.1	-3.8	-3.1	-3.1	-4.0
5-CH ₃	10.2	9.4	8.0	3.7	14.1
C6-H6	12.0	12.1	12.0	12.8	11.5
Q factor		0.13	0.52	0.31	0.30

MM2

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{SSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	3.4	-9.1	0.1	-5.1
C2-H2	1.1	3.9	1.1	2.6	1.3
C3-H3	13.5	11.6	13.4	13.4	13.6
3-CH ₃	-8.8	0.7	-8.5	0.8	-7.1
C4-H4	-4.1	-1.7	-4.1	-2.4	-4.2
5-CH ₃	10.2	4.5	10.5	5.0	12.8
C6-H6	12.0	12.5	11.9	12.5	11.6
Q factor		0.69	0.04	0.58	0.20
Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	-7.3	0.9	-10.8	1.7
C2-H2	1.1	3.3	0.3	2.9	0.7
C3-H3	13.5	12.9	13.9	12.7	12.5
3-CH ₃	-8.8	-0.2	-0.4	-4.5	-8.7
C4-H4	-4.1	-2.8	-3.1	-2.9	-3.5
5-CH ₃	10.2	2.0	6.8	2.5	13.2
C6-H6	12.0	12.9	12.2	13.0	11.5
Q factor		0.50	0.55	0.38	0.45

MMFF9

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{SSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	-0.4	-8.7	-1.6	-5.9
C2-H2	1.1	3.9	1.1	3.4	1.3
C3-H3	13.5	11.9	13.4	12.5	13.5
3-CH ₃	-8.8	3.0	-8.6	2.6	-7.2
C4-H4	-4.1	-1.3	-4.1	-2.2	-4.1
5-CH ₃	10.2	5.4	11.0	4.7	12.9
C6-H6	12.0	12.6	11.9	12.8	11.6
Q factor		0.65	0.05	0.61	0.18

Appendix

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	-7.4	0.6	-11.5	0.6
C2-H2	1.1	3.4	0.4	2.0	0.7
C3-H3	13.5	12.9	13.9	13.2	12.5
3-CH ₃	-8.8	-1.2	-0.7	-6.3	-8.5
C4-H4	-4.1	-2.7	-3.1	-3.3	-3.7
5-CH ₃	10.2	1.4	7.6	5.4	13.8
C6-H6	12.0	13.0	12.2	12.6	11.5
Q factor		0.49	0.52	0.25	0.42

Unoptimized

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{SSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{SSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	5.0	-8.5	0.8	-6.7
C2-H2	1.1	2.6	1.1	2.4	1.3
C3-H3	13.5	11.8	13.3	13.6	13.6
3-CH ₃	-8.8	-7.3	-8.3	-0.2	-7.8
C4-H4	-4.1	-1.8	-4.1	-2.4	-4.2
5-CH ₃	10.2	9.1	11.6	5.9	12.3
C6-H6	12.0	12.2	11.8	12.4	11.6
Q factor		0.59	0.07	0.56	0.14

Coupling	D_{Exp}/Hz	$D_{\text{Calc}}(\text{RSR-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSR-}ax)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}eq)/\text{Hz}$	$D_{\text{Calc}}(\text{RSS-}ax)/\text{Hz}$
1-CH ₃	-8.9	-10.7	1.3	-10.8	1.8
C2-H2	1.1	1.4	0.3	1.4	0.5
C3-H3	13.5	13.3	14.0	13.4	12.2
3-CH ₃	-8.8	-7.9	-2.1	-7.2	-8.9
C4-H4	-4.1	-3.7	-3.5	-3.7	-3.3
5-CH ₃	10.2	8.8	8.1	8.1	13.3
C6-H6	12.0	12.2	12.0	12.3	11.6
Q factor		0.11	0.51	0.14	0.46

Exemplary MSpin Input and Output

Input file:

```
rdc_data {
  2 16 -11.4
  3 17 11.9
  6 18 12.7
  9 21 10.2
  9 20 4.5
  8 19 16.1
  14 25 -6.5
  14 26 -6.5
  14 27 -6.5
  15 28 6.6
  15 29 6.6
  15 30 6.6
}
permutations {
  20 21
}
```

Output file:

```
Solution1
Conformer 1
Distribution size=128
Distribution type=Gaussian
RDC general Std.Dev=0.5 Hz
Alignment tensor
<A`x> = -1.369e-04 Std.Dev = 6.718e-06
<A`y> = -3.696e-04 Std.Dev = 9.399e-06
<A`z> = 5.065e-04 Std.Dev = 8.449e-06
Quality factors statistic
<Q> = 0.203
StdDev(Q) = 0.012
Highest Q = 0.231
Lowest Q = 0.169
```

Structural Noise Script

```
filename = raw_input("Enter xyz file name ")
n_fixed_atoms = input("Enter number of coordinates to vary (one block at the end of the file)")
n_distances = input("Enter number of radii you want to apply")

distances = []
for _ in range(n_distances):
    random_distance = input("Enter radius of sphere in Angstrom (random displacement of atoms within this sphere)")
    distances.append(random_distance)
molecules_varied = input("Enter desired number of molecules for each distance")
for k in range(n_distances):
    atoms = []
    coordinates = []
    xyz_startmolecule = open(str(os.path.realpath("__file__")[:-8])+str(filename)+".xyz", "r")
    n_atoms = int(xyz_startmolecule.readline())
    title = xyz_startmolecule.readline()
    for line in xyz_startmolecule:
        atom,x_coordinate,y_coordinate,z_coordinate = line.split()
        atoms.append(atom)
        coordinates.append([float(x_coordinate), float(y_coordinate), float(z_coordinate)])
    xyz_varied = open(str(os.path.realpath("__file__")[:-8])+str(filename)+"_var_"+str(distances[k])+".xyz", "a")
    for _ in range(int(molecules_varied)):
        xyz_varied.write(str(n_atoms)+"\n\n")
        for i in range((n_atoms)):
```

```
xyz_varied.write(str(atoms[i])+" ")
if i < (n_atoms - n_fixed_atoms):
    for j in range(3):
        old_coods = "{0:.5f}".format(coordinates[i][j])
        xyz_varied.write(str(old_coods)+" ")
    xyz_varied.write("\n")
else:
    u = random.uniform(0, 1)
    r = distances[k]*sp.special.cbrt(u)
    costheta = random.uniform(-1, 1)
    theta = np.arccos(costheta)
    phi = random.uniform(0, 2*np.pi)
    new_x = "{0:.5f}".format(float(coordinates[i][0])+r*np.sin(theta)*np.cos(phi))
    new_y = "{0:.5f}".format(float(coordinates[i][1])+r*np.sin(theta)*np.sin(phi))
    new_z = "{0:.5f}".format(float(coordinates[i][2])+r*np.cos(theta))
    xyz_varied.write(str(new_x)+" "+str(new_y)+" "+str(new_z)+"\n")
xyz_varied.write("\n")
xyz_varied.close()
xyz_startmolecule.close()
```


Extended NMR Spectroscopy of Synthetic Indanes

J-coupling Analysis of β -aa-H and β -aa-Me

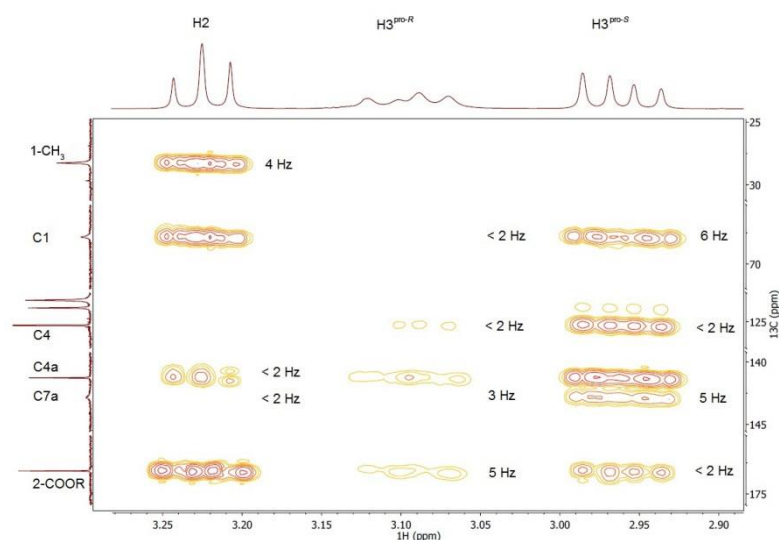


Figure B-5: Excerpt of HMBC spectrum of β -aa-H. For listed *J*-couplings see Table B-2.

Table B-2: Comparison of dihedral angles (in $^{\circ}$) in the two conformations (ester group on C2 equatorial or axial) of the *SS* and *RS* isomers of β -aa-H with experimental 3-bond coupling constants as estimated from high-resolution HMBC experiments. The structures were generated with MM2 geometry optimization.^[151] Angles that are at variance with the experiment are labeled red. Angles that are at variance with the experiment if the assignment of $H3^{pro-S}$ (2.96 ppm) and $H3^{pro-R}$ (3.09 ppm) is swapped are labeled blue.

Atoms	Dihedral Angle		Dihedral Angle		$^3J_{HH}$ or $^3J_{CH}$ / Hz
	<i>SS</i> - <i>eq</i>	<i>SS</i> - <i>ax</i>	<i>RS</i> - <i>eq</i>	<i>RS</i> - <i>ax</i>	
H2-C2-C3- $H3^{pro-S}$	31.8	-33.1	34.2	-34.0	8.5
H2-C2-C3- $H3^{pro-R}$	156.1	90.7	158.7	89.6	9.3
H2-C2-C3-C4a	-89.0	-148.7	-86.3	-149.3	< 2
H2-C2-C1-C7a	88.5	149.8	84.8	149.1	< 2
H2-C2-C1-(1-CH ₃)	-31.9	32.6	-157.3	-87.8	~ 5
$H3^{pro-S}$ -C3-C2-CO ₂ Me	-84.2	-151.2	-83.6	-152.0	< 2
$H3^{pro-R}$ -C3-C2-CO ₂ Me	40.2	-27.4	40.9	-29.4	~ 5
$H3^{pro-S}$ -C3-C2-C1	146.7	86.9	148.8	86.2	~ 6
$H3^{pro-R}$ -C3-C2-C1	-89.0	-149.4	-86.7	-150.3	< 2
$H3^{pro-S}$ -C3-C4a-C4	42.7	81.4	41.5	81.8	< 2
$H3^{pro-R}$ -C3-C4a-C4	-78.9	-40.2	-80.2	-39.9	< 2
$H3^{pro-S}$ -C3-C4a-C7a	-137.2	-99.3	-138.2	-99.2	~ 5
$H3^{pro-R}$ -C3-C4a-C7a	101.2	139.1	100.2	139.2	~ 3

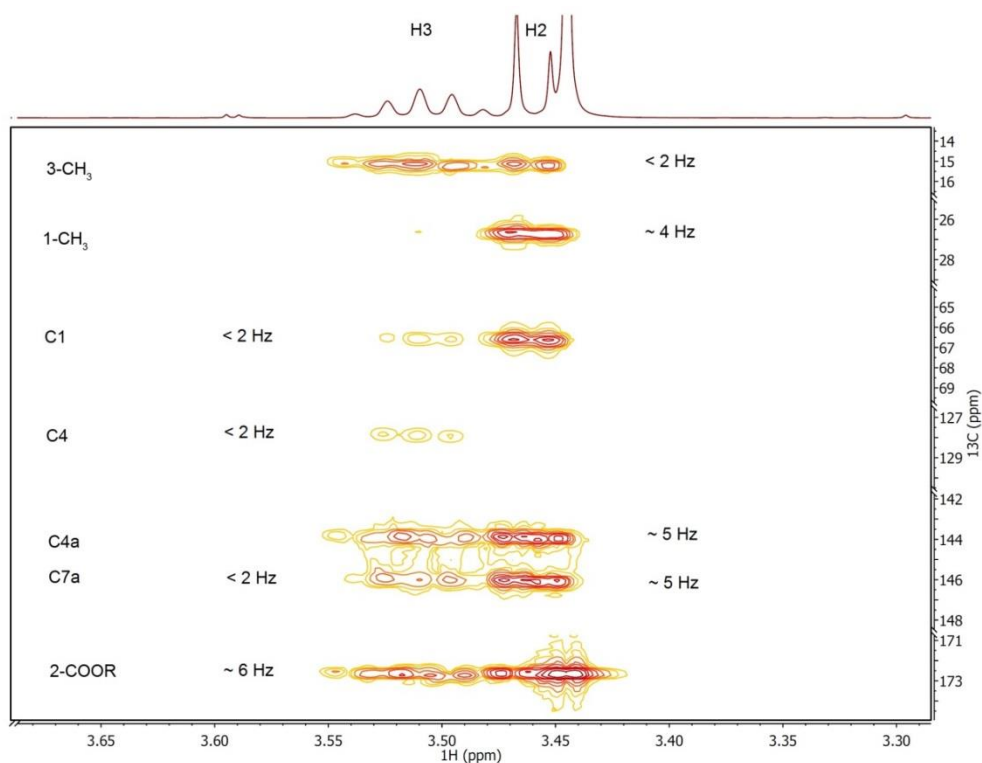


Figure B-6: Excerpt of HMBC spectrum of β -aa-Me. For listed J -couplings see Table B-3.

Table B-3: Comparison of dihedral angles (in $^\circ$) in the two conformations (ester group on C2 equatorial or axial) of *SSR*, *SSS*, *RSR* and *RSS* isomers of β -aa-Me with experimental 3-bond coupling constants as estimated from high-resolution HMBC experiments. The structures were generated with MM2 geometry optimization.^[151] Angles that are at variance with the experiment are labeled red.

Atoms	Dihedral Angle		Dihedral Angle		Dihedral Angle		Dihedral Angle		$^3J_{\text{HX}}/\text{Hz}$
	<i>SSR</i> -eq	<i>SSR</i> -ax	<i>SSS</i> -eq	<i>SSS</i> -ax	<i>RSR</i> -eq	<i>RSR</i> -ax	<i>RSS</i> -eq	<i>RSS</i> -ax	
H2-C2-C3-H3	29.3	-35.6	156.7	103.6	31.6	-38.5	162.6	96.6	7.4
H2-C2-C3-(3-CH ₃)	153.4	87.2	32.7	-19.4	156.9	83.9	37.9	-26.4	< 2
H2-C2-C3-C4a	-87.2	-148.1	-90.6	-139.8	-84.0	-150.7	-84.5	-146.0	~ 5
H2-C2-C1-C7a	86.9	148.6	90.4	140.9	83.4	150.6	83.6	145.9	~ 5
H2-C2-C1-(1-CH ₃)	-33.6	31.1	-30.0	22.6	-158.6	-85.9	-159.0	-91.3	~ 4
H3-C3-C2-CO ₂ Me	-83.0	-155.0	39.8	-14.0	-80.9	-158.4	162.6	-21.5	~ 6
H3-C3-C2-C1	141.3	83.6	-88.3	-136.3	142.6	80.4	-83.8	-142.6	< 2
H3-C3-C4a-C4	47.8	84.2	-79.4	-49.9	47.3	86.5	-82.2	-46.2	< 2
H3-C3-C4a-C7a	-132.1	-95.7	100.0	130.0	-137.7	-93.9	97.8	133.5	< 2

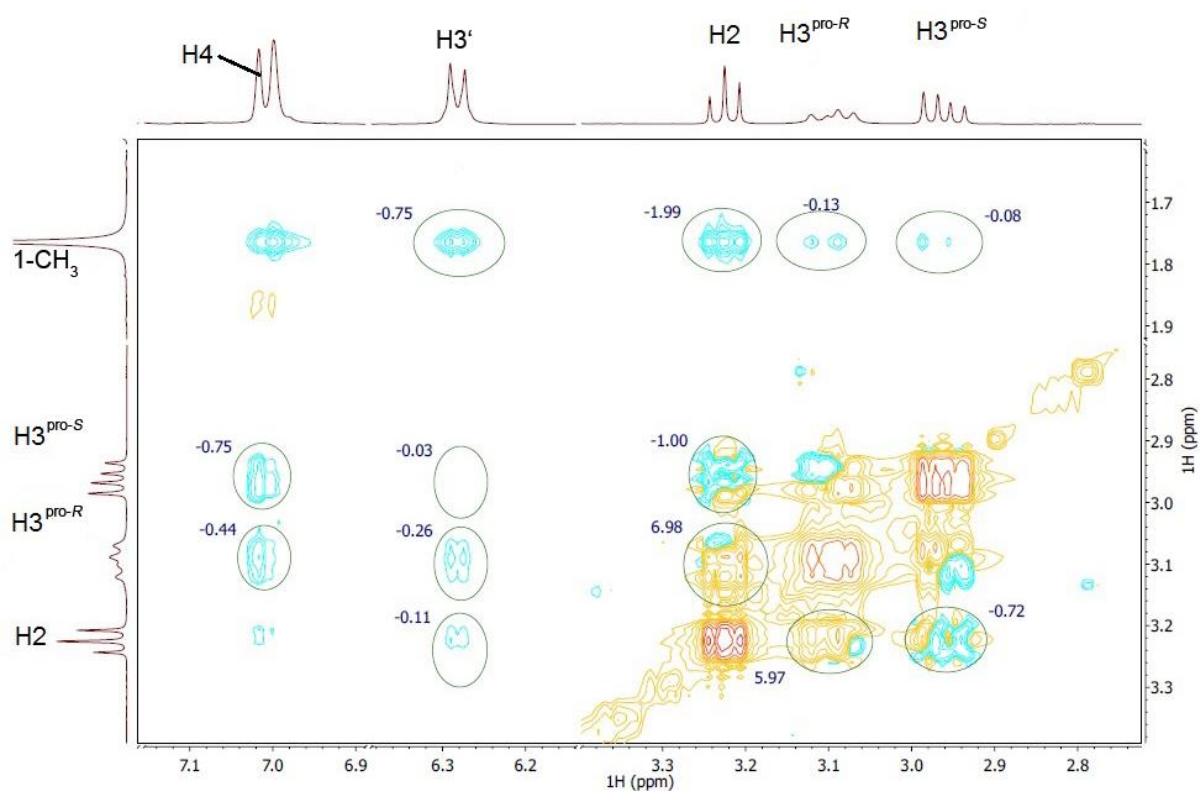
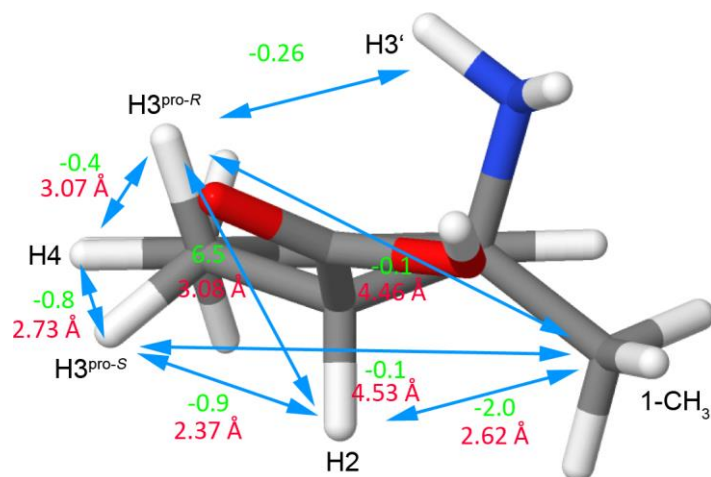
NOE Analysis of β -aa-H and β -aa-MeFigure B-7: Excerpt of NOESY spectrum of β -aa-H. For listed NOE contacts see Table B-4.Figure B-8: Correct structure *SS-eq* (MM2 optimized^[151]). Green: NOE Integral. Red: Distances.

Table B-4: Comparison of inter-proton distances in the two conformations (ester group on C2 equatorial or axial) of the *SS* and *RS* isomers of β -aa-H with experimental NOE contacts as estimated from NOESY integrals. The structures were generated with MM2 geometry optimization.^[151] For methyl groups, the position of the carbon atom was used.

	Distances / Å		Distances / Å		NOE Integral
	<i>SS</i> - <i>eq</i>	<i>SS</i> - <i>ax</i>	<i>RS</i> - <i>eq</i>	<i>RS</i> - <i>ax</i>	
(1-CH ₃)-H2	2.62	2.63	3.49	3.11	-2.0
(1-CH ₃)-H3 ^{pro-S}	4.53	3.20	4.36	4.43	-0.1
(1-CH ₃)-H3 ^{pro-R}	4.46	4.32	3.27	4.62	-0.1
H2-H3 ^{pro-S}	2.37	2.44	2.37	2.42	-0.9
H2-H3 ^{pro-R}	3.08	2.81	3.08	2.79	6.5
H3 ^{pro-S} -H4	2.73	3.13	2.74	3.13	-0.8
H3 ^{pro-R} -H4	3.07	2.73	3.10	2.73	-0.4

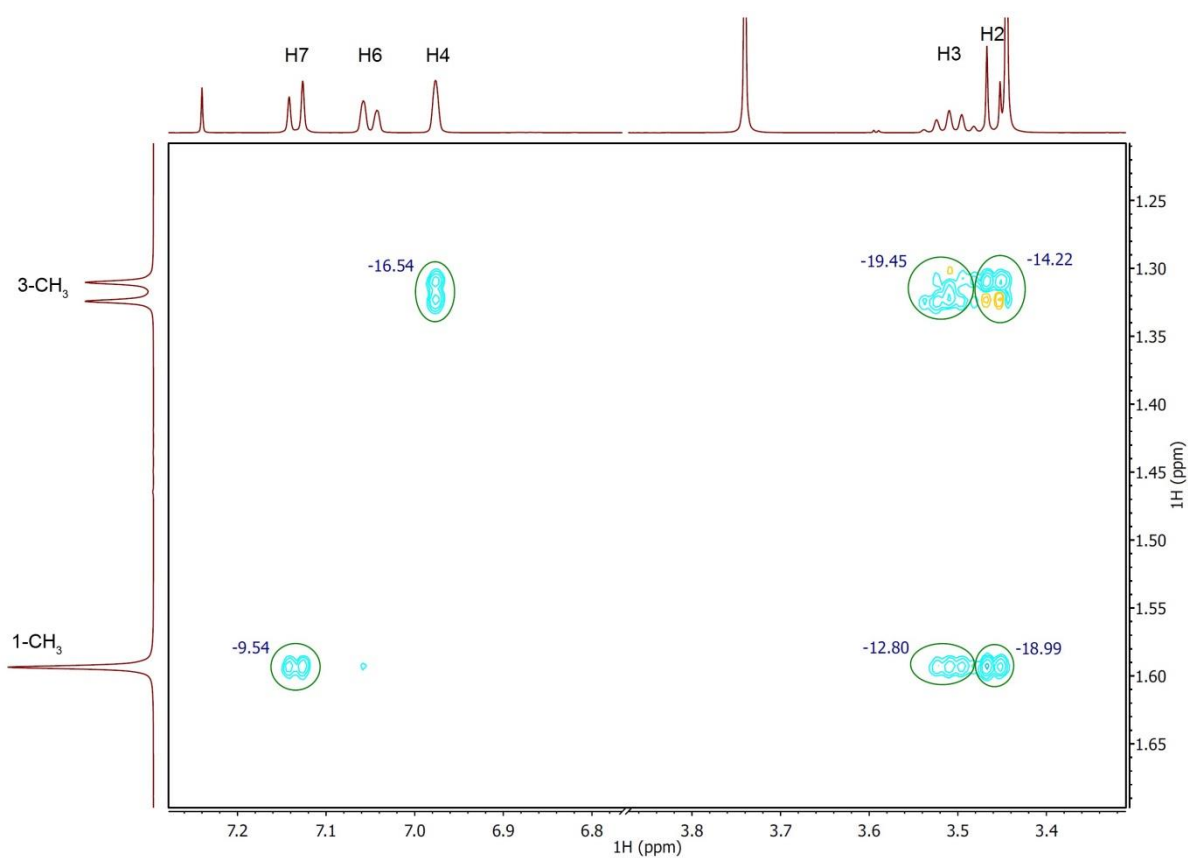


Figure B-9: Excerpt of NOESY spectrum of β -aa-Me. For listed NOE contacts see Table B-5.

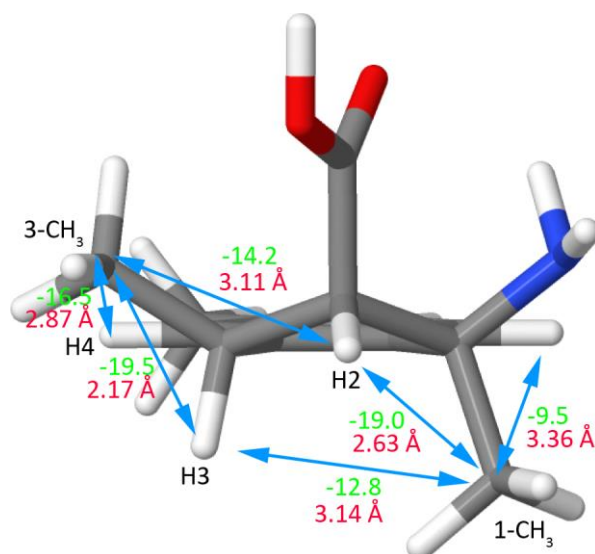


Figure B-10: Correct structure *SSR-ax* (MM2 optimized^[151]). Green: NOE Integral. Red: Distances.

Table B-5: Comparison of inter-proton distances in the two conformations (ester group on C2 equatorial or axial) of the *SSR*, *SSS*, *RSR* and *RSS* isomers of β -aa-Me with experimental NOE contacts as estimated from NOESY integrals. The structures were generated with MM2 geometry optimization. For methyl groups, the position of the carbon atom was used.

	Distances /Å		Distances /Å		Distances /Å		Distances /Å		NOE Integral
	<i>SSR-eq</i>	<i>SSR-ax</i>	<i>SSS-eq</i>	<i>SSS-ax</i>	<i>RSR-eq</i>	<i>RSR-ax</i>	<i>RSS-eq</i>	<i>RSS-ax</i>	
(1-CH ₃)-H2	2.56	2.63	2.60	2.59	3.49	3.08	3.48	3.13	-19.0
(1-CH ₃)-H3	4.47	3.14	4.44	4.40	4.43	4.35	3.19	4.53	-12.8
(1-CH ₃)-H7	2.92	3.36	2.90	3.25	3.27	2.87	3.34	2.90	-9.5
(3-CH ₃)-H2	3.47	3.11	2.60	2.63	3.49	3.07	2.62	2.64	-14.2
(3-CH ₃)-H3	2.17	2.17	2.19	2.17	2.17	2.17	2.19	2.17	-19.5
(3-CH ₃)-H4	3.24	2.87	2.89	3.22	3.22	2.85	2.89	3.27	-16.5

Quantitative Analysis of Residual Quadrupolar Couplings Data

Exemplary ORCA input files EFG tensor calculation

EFG gradient:

```
# ORCA JOB on ORCA optimized structure
#-----
%maxcore X
%pal nprocs Y
#choose X and Y according to hardware
! RKS NORI B3LYP 6-311++G(d,p) Decontract Grid5 NOFINALGRID VERYTIGHTSCF SlowConv Cosmo(solvens)
*xyzfile A B NAME.xyz
#coordinate file forcefield optimized structure; choose A and B according to charge and multiplicity
%eprnmr nuclei = all Q {fgrad}
#choose nuclei Q to calculate EFG for
end
```

Structure Data of [CpLi(OMe₂)], [CpLi(OMe₂)₂], [CpLi(OMe₂)₃] and [Cp₂Li]⁻

[CpLi(OMe ₂)]	[CpLi(OMe ₂) ₂]	[CpLi(OMe ₂) ₃]	[Cp ₂ Li] ⁻
C -1.7566 5.6908 4.7673	O 7.7273 6.8289 3.0828	Li 12.5445 1.8870 9.1681	Li -0.0001 -0.0002 0.0000
C -2.0669 6.9758 5.2737	O 7.9540 5.1835 0.5477	O 11.1578 0.3379 9.5584	C -0.8556 0.5971 2.0554
C -2.8114 6.8078 6.4655	C 5.0524 4.3924 2.9735	O 12.5136 2.3882 11.1430	C -0.6046 1.7702 1.3074
C -2.9625 5.4199 6.6951	C 4.7508 5.7128 3.3770	O 14.4172 0.9200 9.5003	C -1.4431 1.7504 0.1687
C -2.3067 4.7296 5.6473	C 4.4285 6.4591 2.2243	C 12.9404 3.9490 8.1842	C -2.2134 0.5642 0.2144
Li -3.9368 5.9480 4.8130	C 4.5124 5.5967 1.1045	C 13.4617 2.9700 7.3101	C -1.8510 -0.1462 1.3810
O -5.7251 6.4960 4.3190	C 4.8909 4.3166 1.5693	C 12.3836 2.1973 6.8231	H -0.3503 0.3024 2.9650
C -6.3574 6.3404 3.0509	C 8.2667 6.3052 4.2965	C 11.1928 2.7101 7.3845	H 0.1225 2.5336 1.5483
C -6.6395 6.4669 5.4170	C 8.6226 4.0295 1.0516	C 11.5350 3.7863 8.2300	H -1.4700 2.4959 -0.6141
H -3.1994 7.6005 7.0888	C 7.4706 4.9713 -0.7772	C 10.9786 -0.6998 8.6083	H -2.9304 0.2406 -0.5273
H -2.2255 3.6567 5.5416	C 8.6649 7.6641 2.4035	C 9.9320 0.7194 10.1687	H -2.2473 -1.1026 1.6929
H -1.1862 5.4772 3.8747	Li 6.5349 5.7522 1.9701	C 12.7672 1.3906 12.1251	C 0.8574 -0.5939 -2.0553
H -1.7780 7.9175 4.8290	H 5.3448 3.5807 3.6248	C 12.2891 3.6654 11.7229	C 0.6041 -1.7686 -1.3108
H -3.4819 4.9678 7.5280	H 4.7989 6.0970 4.3864	C 15.4567 1.7242 10.0421	C 1.4411 -1.7527 -0.1709
H -5.5815 6.4211 2.2922	H 4.1715 7.5091 2.2001	C 14.9203 -0.0061 8.5476	C 2.2127 -0.5671 -0.2125
H -7.0990 7.1287 2.9014	H 4.3186 5.8752 0.0783	H 13.5167 4.6740 8.7421	C 1.8527 0.1466 -1.3777
H -6.8452 5.3639 2.9792	H 5.0448 3.4378 0.9606	H 14.5066 2.8157 7.0799	H 0.3535 -0.2962 -2.9648
H -6.0585 6.6305 6.3227	H 8.5935 7.1189 4.9503	H 12.4595 1.3543 6.1507	H -0.1236 -2.5305 -1.5545
H -7.1420 5.4973 5.4719	H 7.4799 5.7327 4.7813	H 10.1971 2.3218 7.2220	H 1.4660 -2.5002 0.6101
H -7.3835 7.2601 5.3109	H 7.9306 3.1851 1.1270	H 10.8441 4.3682 8.8244	H 2.9288 -0.2461 0.5312
	H 9.4586 3.7554 0.4017	H 11.9538 -0.9426 8.1928	H 2.2505 1.1033 -1.6866
	H 6.7087 4.1875 -0.7913	H 10.3181 -0.3757 7.7982	
	H 7.0293 5.9066 -1.1139	H 10.5547 -1.5898 9.0863	
	H 9.0040 8.4674 3.0637	H 10.1529 1.5267 10.8612	
	H 9.5285 7.0838 2.0654	H 9.4952 -0.1253 10.7126	
	H 9.1227 5.6529 4.0930	H 9.2211 1.0728 9.4157	
	H 9.0035 4.2768 2.0408	H 12.9207 0.4525 11.6003	
	H 8.2951 4.6879 -1.4379	H 11.9182 1.2996 12.8103	
	H 8.1585 8.0915 1.5410	H 13.6620 1.6372 12.7063	
		H 12.0675 4.3544 10.9145	
		H 13.1784 4.0019 12.2667	
		H 11.4446 3.6280 12.4193	
		H 14.9982 2.4333 10.7248	
		H 15.9753 2.2721 9.2487	
		H 16.1798 1.1023 10.5812	
		H 14.0784 -0.5618 8.1422	
		H 15.6215 -0.7010 9.0221	
		H 15.4288 0.5165 7.7313	

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CONFERENCES AND WORKSHOPS

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- Summer School on the Theory of NMR financed by the Volkswagen Stiftung, 04.-10.09.2016 in Windischleuba (Workshop)
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SCIENTIFIC PUBLICATIONS

1. T. Niklas, D. Stalke, M. John, "Single-shot titrations and reaction monitoring by slice-selective NMR spectroscopy" *Chem. Commun.*, **2015**, 51, 1275-1277.
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