Lipid dependent interactions of biomimetic membrane models with the ENTH domain of epsin

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I, Nelli Teske, hereby certify that my doctoral thesis entitled "Lipid dependent interactions of biomimetic membrane models with the ENTH domain of epsin" has been written independently and with no other source and aids than quoted.

Göttingen, 2019

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Meiner Familie

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ABSTRACT. Membrane remodeling processes during endocytosis are highly regulated by the protein-protein and protein-membrane interactions. In clathrin-mediated endocytosis (CME) binding of the protein epsin to its receptor lipid phosphatidylinositol-(4,5)-bisphosphate (PIP₂) induces the deformation of the membrane. Upon binding of the epsin *N*-terminal homology domain (ENTH) to PIP₂, conformational changes in the protein lead to a newly formed helix, which inserts into the cytosolic leaflet.

By using artificial membrane models the binding affinity and the surface topology upon ENTH binding as a function of the lipid composition were analyzed. With increasing the PIP₂ concentration the experiments on solid supported lipid bi- and monolayers proved a direct relation between the protein occupancy and the lipid content. Subsequently, this indicates that an accumulation of PIP₂ on the cytosolic leaflet can facilitate the endocytosis rate.

Besides PIP₂ also negatively charged lipids with the head group phosphatidylserine (PS) can affect protein binding in endocytosis. In presence of PS, higher binding affinities and protein occupancies of ENTH to PIP₂ doped membranes were observed. Although ENTH is known to act as a monomeric protein, atomic force microscopy (AFM) measurements revealed the appearance of protein clusters induced by PS. Thus, also membrane crowding seems to have an impact on the curvature inducing step in CME. Mutation of the amino acid R114 showed its relevance in ENTH cluster formation as no oligomers were observed with the mutant R114A.

Moreover, monolayer penetration experiments were performed to analyze the surface activity of ENTH dependent on the lipid composition. Increasing the PIP₂ content increased the critical surface pressure. Addition of PS did not significantly increase the penetration of ENTH into monolayers in a PIP₂ dependent manner, although a higher protein occupancy on supported lipid bilayers was observed.

To investigate whether ENTH has an influence on mechanical properties of membranes in the presence of PS, giant unilamellar vesicles (GUVs) were adhered to NeutrAvidin coated surfaces. Incubation of these GUVs with ENTH resulted in the rupturing of those due to lipid packing defects in the membrane by helix insertion and PS induced clusters. These results allow to understand how PS alters the binding structure of ENTH to PIP₂ doped membranes. This in turn also show that the lipid composition contributes to the regulation of protein-dependent membrane deformation during CME.

ZUSAMMENFASSUNG. Während der Endozytose wird die Membran durch verschiedene Protein-Protein und Protein-Lipid Interaktionen reguliert. In der Clathrin-vermittelten Endozytose (CME) bewirkt die Bindung des Proteins Epsin an sein Rezeptorlipid Phosphatidyl-(4,5)-bisphosphat (PIP₂) eine Membrankrümmung. Durch diese Bindung wird in der Epsin *N*-terminalen Homologie Domäne (ENTH) eine neue Helix (α_0) ausgebildet, die die cytosolische Lipidmonoschicht der Plasmamembran penetriert.

Mit artifiziellen Membransystemen konnten die Bindungsaffinitäten und die Oberflächentopologien bei ENTH-Anbindung in Abhängigkeit der Lipidzusammensetzung analysiert werden. Mit zunehmender PIP₂-Konzentration zeigten die Experimente an fest unterstützten Lipiddoppelschichten und -monoschichten einen direkten Zusammenhang zwischen der Proteinbelegung und der Rezeptorlipidkonzentration. Dies wiederum deutet darauf hin, dass eine Anhäufung von PIP₂ in der Plasmamembran die Endozytose-Rate steigern kann.

Nicht nur bei der Erhöhung der PIP₂-Konzentration wurde eine höhere Proteinoberflächenbelegungen auf Membranen beobachtet, sondern auch in Gegenwart der Lipidkopfgruppe Phosphatidylserin (PS), welche die Bildung von ENTH Cluster induziert. Somit scheint neben der Helixinsertion auch das Zusammenlagern mehrerer ENTH-Moleküle (Crowding) einen Einfluss auf den krümmungsinduzierenden Schritt in der CME zu haben. Mit der Mutante ENTH R114A konnte zudem gezeigt werden, dass die Aminosäure R114 essentiell für die Clusterbildung ist.

Penetrationsexperimente an Lipidmonoschichten wurden durchgeführt, um die Oberflächenaktivität von ENTH in Abhängigkeit von der Lipidzusammensetzung zu analysieren. Dabei zeigte die PIP₂ Konzentration einen größeren Einfluss auf den kritischen Oberflächendruck als die Zugabe von PS.

Adhärierte Riesenvesikel (GUVs) ermöglichten zudem die mechanischen Eigenschaften von Membranen in Abhängigkeit von PS zu untersuchen. Inkubation dieser GUVs mit ENTH führte zum Rupturieren der Vesikel durch die Entstehung von Defekten während der Helixinsertion und der PS-induzierten Clusterbildung. Dies zeigt, dass PS einen deutlichen Einfluss auf die Bindungsstruktur von ENTH an PIP₂ Membranen hat. Dadurch konnte auch gezeigt werden, dass die Lipidzusammensetzung zur Regulation der proteinabhängigen Membrankrümmung während der CME beiträgt.

1 Introduction

Cells are complex systems, consisting of different cell organelles with various functions. The cell membrane surrounds the cytoplasm of the cell and serves as protection from the environment. This barrier function enables the selective control of cell communication and molecule exchange. Both processes rely on substance transport, which can either be passive by simple diffusion or by carrier and channel proteins. Contrary to that, active transport requires energy to regulate the uptake of particles. However, large molecules are not able to pass the membrane even though the cell supplies the energy. To overcome this barrier, the transport of large molecules is driven by exo- and endocytosis.

1.1 Clathrin-mediated endocytosis

One important and well known endocytotic pathway is the clathrin-mediated endocytosis (CME). Here macromolecules are internalized into vesicles derived from the plasma membrane. The receptor regulated process enables the controlled entry into cells, which is important for immune response or intercellular communication. The process is connected to highly coordinated and complex interactions, preventing the entry of unwanted viruses or toxins. Although under certain conditions, like dysfunctions, they can gain entry and cause serious diseases. The clathrin-mediated endocytosis (CME).

Since clathrin does not have a binding motif for the membrane, adaptor proteins connect clathrin with membrane lipids and other proteins.⁵ The uptake of cargo-molecules is driven by a complex system of proteins, proceeded in five steps.

By membrane binding of accessory proteins first a curvature of the bilayer occurs, followed by the cargo selection via receptor molecules. Then the recruitment of the clathrin triskelia molecules induces the polymerization of clathrin, forming a stable cage around the invagination (clathrin-coated pit).^{6–8} Scission trough proteins, like dynamin,^{9,10} leads to the formation of the clathrin-coated vesicle (CCV). In the last

step, the clathrin triskelia cage decays and consequently the cargo molecules are released. This five step mechanism of the clathrin-mediated endocytosis is illustrated in Figure 1.1.

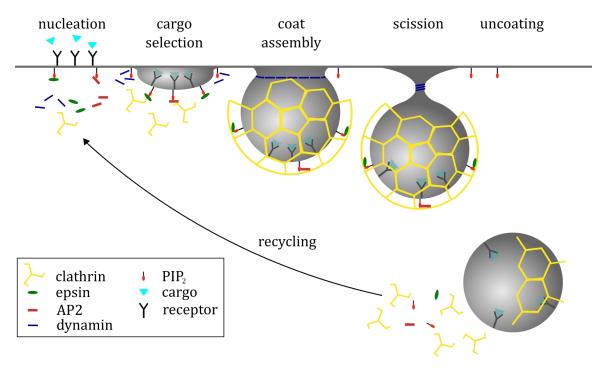


Figure 1.1: Schematic illustration of the clathrin-mediated endocytosis. Binding of proteins like epsin and AP2 to the receptor lipid PIP_2 induces an invagination, during which cargo molecules are recruited by receptor proteins. Afterwards polymerization of clathrin leads to a curvature stabilizing cage. Due to further curvature and coat assembly a vesicle is formed. Scission and uncoating result in a free vesicle, which is transported to the target cell compartment. Figure modified according to Messa $et\ al.^{11}$

The cargo uptake is highly regulated by receptor proteins. Adaptors like AP2, AP180/CALM and epsin are noteworthy as they exhibit binding sites for cargo molecules and bind clathrin as well as the membrane. Dephosphorylation of phosphatidylinositol-(4,5)-bisphosphate (PIP2) by synaptojanin induces a decreased binding to clathrin associated proteins, which destabilizes the clathrin coat and consequently results in uncoating. Uncoated vesicles typically are transported to early endosomes, where they fuse and release their cargo. Thereafter, the receptors are recycled and recruited to be available for the next CME circle. The role of various proteins in CME is still not completely understood. Nevertheless, all membrane-interacting adaptors in CME seem to have at least two functions. Cargo selecting receptors with

the ability to generate membrane curvature ensure the vesicle formation independent from the cargo molecule.¹ One molecule, which is crucial for membrane reorganization during CME is epsin.

1.2 Epsin *N*-terminal homology domain

The accessory protein epsin is among other clathrin-associated sorting proteins essential for the CME. It was first discovered in 1998 by Chen *et al.* as <u>Eps</u>15 <u>in</u>teracting protein (epsin). ¹⁸ Epsin 1 and the three isoforms are present in all vertebrates. ⁷ Other epsin-like homologues were also found in lower species. ENTH contains several binding motifs, important for different functions. For example, it is essential in signaling activation of the epidermal growth factor receptor (EGFR), contributed in cancer diseases, or the epithelial sodium channel as it regulates electrolyte balances. The *C*-Terminus of epsin has a binding site for proteins like Eps15 and intersectin. ¹⁸ Further motifs in the protein can interact with ubiquitin and AP2 (Figure 1.2). ^{19,20}

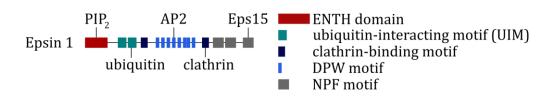


Figure 1.2: Schematic structure of epsin 1 with all important binding motifs, necessary for interactions with PIP₂, ubiquitin, AP₂, clathrin and Eps15. Derived from Horvath *et al.*⁷

In the CME epsin is recruited to the plasma membrane by the evolutionary well conserved epsin N-terminal homology (ENTH) domain.²⁰ It consists of about 150 amino acids structured to a superhelix of seven α -helices. At the plasma membrane the domain specifically recognizes and binds the receptor lipid PIP₂, resulting in a conformational change in the tertiary structure of its ENTH domain. This however leads to the formation of a PIP₂ binding pocket established by the amino acids R7, R8, K11, R25, R63, K69 and H73 of the ENTH domain, which accommodates the negative charge of the receptor lipid.^{8,21} An unstructured 14 amino acids large sequence of the domain becomes ordered, called the α_0 helix. The newly formed helix inserts into the

inner leaflet of the membrane, reduces the energy barrier of bending and consequently modifies membrane curvature. Ford *et al.* were able to co-crystallize the ENTH domain together with inositol-(1,4,5)-triphosphate (IP₃), the head group of PIP₂ (Figure 1.3)²², showing the structural change in the protein.

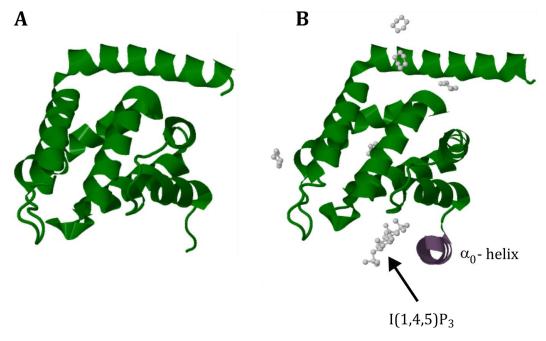


Figure 1.3: Crystal structures of unbound ENTH and the ENTH-IP $_3$ complex. Binding to IP $_3$ leads to the formation of a new helix (α_0). In the bound state the R114 loop of ENTH is in close contact to the membrane. ^{22,23}

Itoh *et al.* showed that the ENTH-PIP₂ binding is specific as the affinity to other inositol species and phospholipids was drastically decreased or missing.²⁴ Moreover the binding affinity can be influenced by the pH or the lipid composition of the membrane.^{8,25,26} The ENTH domain furthermore can induce tubulation in PIP₂-containing liposomes and also tubular structures in the plasma membrane of living cells.^{21,27,28} Besides the curvature inducing effect, epsin also binds to clathrin in the CME.²⁹ It connects clathrin with the membrane and causes its polymerization, which stabilizes the invagination and the subsequent vesicle formation (cf. chapter 1.1).

The complete contribution and interaction of ENTH domain in the network of clathrin associated proteins is still discussed and thus the investigation is important. Some studies indicate that not only the insertion of the amphiphilic helix leads to a curvature, but also the assembly of the protein is contributing the membrane curvature process.³⁰

1.3 Interplay of the ENTH assembly and amphiphilic helix insertion

Many biological processes depend on highly curved membrane structures, even though induction of curvature is associated with high energy requirements. To overcome this energy barrier an interplay between several proteins in collaboration with lipids is necessary. There are different ways to generate membrane deformation. The epsin ENTH domain is known to induce membrane asymmetry by the insertion of an amphiphilic helix. Ford *et al.* performed mutagenesis experiments, where they changed the hydrophobicity of the α_0 helix on position L6.²¹ Decreased hydrophobicity resulted in lower binding affinity and decreased tubulation, showing the impact of the amphiphilic helix of ENTH.

Nevertheless, computational studies challenge the efficiency of helix insertion to induce membrane deformation.^{31,32} According to these studies, about 10-25% of the membrane surface have to be occupied by helices. Considering the footprint of ENTH, a physiological improbable coverage of almost 100% would be necessary.²³ Amphiphilic helix insertion induced deformation of membranes was also found for other proteins, but usually this effect was supported by the shape of the protein. An example is the *N*-BAR domain of endophilin.³³ The banana-shaped protein domain is enriched in synapses and stabilizes the helix induced membrane curvature by its curved shape. In contrast to that, ENTH does not show an inherent curvature itself. This leads to the question if an additional effect supports the helix insertion and promotes membrane curvature.

Stachowiak *et al.* assumed that crowding of ENTH drives membrane bending due to lateral pressure resulting from collision of two bound proteins.³⁰ With increasing protein coverage the pressure increases until the energy barrier is reached to induce membrane deformation. They proved the relevance of the protein crowding mechanism with His-tagged ENTH lacking the α_0 helix. Fluorescence microscopy experiments with giant unilamellar vesicles showed that high coverage with His-tagged ENTH also resulted in tubular structures, indicating the contribution of protein-protein interaction on the membrane deformation process.

Contrary to that, Kozlov *et al.* exhibited with computational predictions the high relevance of the ENTH insertion mechanism.³⁴ They calculated the membrane tube radius for different models, which predict that the crowding effect is clearly weaker than the insertion effect. Protein crowding as a non-specific interaction with the membrane would lead to bending on both sides of the membrane, thus counteract the effect on each side. Further studies investigated the protein crowding induced deformation, assuming that it is co-regulated by membrane tension and protein coverage.³⁵ Recent studies showed that binding of ENTH to PIP₂-containing membranes resulted in the decrease of lateral membrane tension and consequently in a decreased bending modulus.³⁶ Hence, an interplay of both, the protein crowding and the protein helix insertion of ENTH, seems likely to induce membrane curvature. In Figure 1.4 the predicted mechanism involved in ENTH induced membrane bending is illustrated.

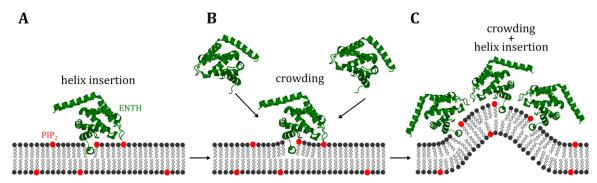


Figure 1.4: Schematic illustration of the predicted mechanism of ENTH induced membrane curvature. Recruitment of ENTH by PIP₂ results in the insertion of an amphipathic helix (A). The helix insertion of ENTH disturbs the structure of the lipid bilayer (B). Rearrangement leads to a deformation of the membrane. The assembly of several ENTH monomers facilitates further membrane curvature (C), which represents the first step in CME.

Besides the effect of proteins, also the lipid composition affects the membrane topology. Due to versatile lipid shapes, charges and intramolecular interactions, they can induce membrane asymmetries promoting membrane curvature.³⁷

1.4 Lipid dependent membrane deformation

The plasma membrane consists of various lipids and proteins interacting with each other. Amphiphilic lipids have a hydrophilic head group and long hydrocarbon chain residues. On one hand the lipids in the membrane built the barrier, but additionally they also have other functions for living cells. They are necessary for the energy storage or as signal molecules. 38,39 Furthermore, the lipid structure or geometry can influence the shape of the membrane due to different packing densities. 40,41 If lipids have a large area ratio of head groups to acyl chains, they can create an invagination on the other side of the membrane. The lipids in the membrane are usually distributed asymmetrically, which arises from dynamic flip-flop and lateral diffusion processes. 42,43 Among other effects, this asymmetry can lead to the deformation of the membrane. Certain proteins or molecules are localized in specific compartments of the cell. Asymmetric distribution or even crowding of the lipids in these regions would ensure the protein-lipid interactions, important for several processes. More often the relevance of lipids in membrane deformation is coupled to proteins, which induce this effect. One important lipid, which binds several proteins involved in membrane curvature is PIP₂. In mammalian cells it is the most abundant phosphoinositide species as it incorporates 1% of all phospholipids, occurring in the plasma membrane.44 It is the source of the two second messengers diacylglycerol as well as inositol-1,4,5-triphosphate and also interacts with proteins important for the membrane's attachment to the cytoskeleton. Moreover, several studies revealed the essential role of PIP₂ in exo- and endocytotic processes. 45-48 For example, it is an important regulator of the BAR domain membrane deformation activity, required in endocytosis. In the CME it is also the major regulator for various clathrin associated proteins. Due to recruitment of AP2 and epsin, the membrane is bent until a vesicle is formed (cf. chapter 1.1). Also the uncoating of CCVs is regulated by PIP₂, enabling the fusion of the vesicle with the membrane of the target cell compartment. 15,49 Even though PIP2 is probably the most essential lipid involved in membrane deformation processes, also other lipids affect them. For example, phosphatidylserine (PS) lipids can influence these processes. They are highly enriched in the plasma membrane (12%),^{37,50} especially in the cytosolic leaflet. The anionic lipid is involved in the coagulation cascade and also in the apoptosis process.^{51,52} Due to the negative charge it causes membrane deformation by the recruitment of proteins. In neuronal exocytosis PIP₂ and PS bind synaptotagmin-1, which triggers fusion.⁵³ In yeast, the absence of PS translocation leads to an insufficient membrane curvature as PS drives the assembly of involved proteins like clathrin.⁵⁴ Furthermore, in CME it recruits proteins to the membrane and significantly affects the efficiency of this process.⁵⁵

The exact influence of lipid compositions and their interplay with proteins in the CME is still poorly understood. To investigate such protein-membrane interactions often artificial membrane models were used to mimic biological membranes in a simplified way. This allows the analysis of the contribution of certain lipids in protein induced membrane deformation processes under controlled conditions.

2 Scope of Thesis

Endocytosis processes are regulated by interactions of several proteins and lipids. It enables the transport of important nutrients and signal molecules into cells. Thereby, the clathrin-mediated endocytosis (CME) represents the main pathway of endocytosis, thus the investigation is of great interest. During the CME adaptor proteins, like epsin, bind to the receptor lipid phosphatidylinositol-(4,5)-bisphosphate (PIP₂) and induce a membrane curvature. Upon specific binding of the epsin *N*-terminal homology (ENTH) domain to PIP₂, a former unstructured region rearranges to a new ordered helix. Insertion of this helix induces asymmetry within the both bilayer leaflets and facilitates the membrane deformation. The influence of lipids in this mechanism is still not understood completely and will be characterized more precisely within the scope of the thesis.

Therefore, different methods will be used to analyze specific interactions of the ENTH with PIP2-doped artificial membranes as a function of the lipid composition. In particular the influence of the lipid receptor and the negatively charged lipid phosphatidylserine (PS) on binding affinities and the helix insertion of ENTH will be evaluated. Solid supported lipid membranes will be established to determine the surface occupancy and the dissociation coefficient of the ENTH domain. Thereby, the differences in the binding behavior to PIP₂ containing bi- and monolayers will be analyzed as well as the influence of PS. To investigate if besides the helix insertion also a membrane crowding effect induces a membrane deformation, atomic force microscopy will give information about the surface topologies upon ENTH binding. Moreover, mutagenesis experiments moreover shall show the relevance of specific regions within the protein. Monolayer penetration experiments on a Langmuir-Blodgett trough will clarify how the lipid packing influences the helix insertion and consequently the membrane deformation. Finally, the impact of ENTH on mechanical properties of free-standing membranes in the presence of PS will be analyzed using adhered GUVs to understand the contribution of lipid compositions and membrane tension as regulators for membrane deformation processes.

3 MATERIALS AND METHODS

3.1 Used Materials

3.1.1 Buffers

For the experiments in this work different buffer systems were used. In Table 3.1 the compositions of these are listed. All buffers were diluted in ultrapure water, degased and filtered (pore size \emptyset 0.2 nm).

Table 3.1: Used buffers and their compositions.

buffer	composition
citrate buffer	Na-citrate (20 mM)
	KCl (50 mM)
	EDTA (0.1 mM)
	NaN ₃ (0.1 mM)
phosphate buffered saline (PBS), pH 7.4	KCl (0.3 mM)
	NaCl (140 mM)
	KH ₂ PO ₄ (0.1 mM)
	Na ₂ HPO ₄ (10 mM)
protein buffer, pH 7.4	NaCl (200 mM)
	HEPES (10 mM)
HEPES buffer, pH 7.4	HEPES (10 mM)
	KCl (160 mM)
sucrose buffer	HEPES (2 mM)
	sucrose (298 mOsmol/kg)

3.1.2 Lipids

Artificial membrane systems were generated using different lipids. These exhibit various head groups, like phosphocholine (PC) and phosphoethanolamine (PE), which are found in plasma membranes in a high amount.⁵⁶ In order to mimic biological conditions experiments with these lipids were performed. Furthermore lipids with the negative lipid head group phosphatidylserine (PS) were used as the plasma membrane is also enriched with PS (12%), especially the cytosolic leaflet (cf. chapter 1.4).^{37,50}

Cholesterol was added in giant unilamellar vesicles (GUVs) as it is known to increase the mechanical stability of membranes. Moreover it can regulate the membrane fluidity.^{37,57}

1-palmitoyl-2-oleoyl-sn-glycero-3-phosphocholine (POPC), M = 760.1 g/mol

1-palmitoyl-2-oleoyl-sn-glycero-3-phosphoethanolamine (POPE), M = 718.0 g/mol

1-palmitoyl-2-oleoyl-sn-glycero-3-phospho-L-serine (POPS), M = 784.0 g/mol

1,2-dioleoyl-sn-glycero-3-phosphocholine (DOPC), M = 786.1 g/mol

1,2-dioleoyl-sn-glycero-3-phosphoethanolamine (DOPE), M = 744.0 g/mol

1,2-dioleoyl-sn-glycero-3-phospho-L-serine (DOPS), M = 810.0 g/mol

Cholesterol, M = 386.65 g/mol

Protein binding lipids

In this thesis two different lipids were used to bind proteins. L- α -phosphatidylinositol-4,5-bisphosphate (PIP₂) is the specific receptor lipid for several proteins and exhibits different functions in biological cells (cf. chapter 1.4). For the experiments in this thesis brain extracted PIP₂ was used to bind ENTH. It mainly consists of the fatty acids stearoylic and arachidonoylic acid and has a net charge of -3 to -5 dependent on the pH (-4 at pH 3.7).

Moreover 1,2-dioleoyl-*sn*-glycero-3-phosphoethanolamine-*N*-(cap biotinyl) (cap-biotin-DOPE) was used to bind the protein NeutrAvidin in order to immobilize vesicles via biotin-NeutrAvidin interactions.

L- α -phosphatidylinositol-4,5-bisphosphate (PIP₂), M_{Av} = 1096.4 g/mol

1,2-dioleoyl-sn-glycero-3-phosphoethanolamine-N-(cap biotinyl) (cap-biotin-DOPE), M = 992.3 g/mol

$$C_8H_{17}$$
 C_8H_{17}
 C_8H

Fluorophores

To image the artificial membrane models with fluorescence microscopy, the fluorophores Texas Red® 1,2-dihexadecanoyl-*sn*-glycero-3-phosphoetanolamine (TxR) and Atto488 1,2-dipalmitoyl-*sn*-glycero-3-phosphoetanolamine were added to the membranes.

Texas Red® 1,2-dihexadecanoyl-sn-glycero-3-phosphoetanolamine (TxR), M = 1381.9 g/mol

Atto488 1,2-dipalmitoyl-sn-glycero-3-phosphoetanolamine (Atto488-DPPE), M = 1263.6 g/mol

3.1.3 Proteins

3.1.3.1 Epsin N-terminal homology domain

The highly conserved epsin N-terminal homology (ENTH) domain is an adaptor protein involved in clathrin-mediated endocytosis. It consists of approximately 150 amino acids (M = 16.75 kDa) and can be found in many endocytotic proteins. The ENTH domain and its mutant R114A used in this work were isolated and purified by Dr. Benjamin Kroppen and Indrani Mukherjee from the group of Prof. Dr. Michael Meinecke (Department of Cellular Biochemistry, University Medical Center Göttingen). The protein was stored in protein buffer at -80 °C. Before use the protein was thawed on ice.

3.1.3.2 NeutrAvidin

NeutrAvidin is a 60 kDa protein and was used for the immobilization of biotinylated giant unilamellar vesicles via a NeutrAvidin-biotin interaction. The protein is the deglycosylated form of the tetrameric glycoprotein avidin. NeutrAvidin as well as avidin have high binding affinities to biotin, which makes them well-suited for immunological assays and biosensing applications. Deglycosylation does not reduce the biotin binding ability, but due to the more neutral isoelectrical point (pI = 6.3) than avidin (pI = 10.5), less non-specific interactions to surfaces occur.

3.2 Preparative methods

3.2.1 Preparation of unilamellar vesicles

Vesicles are well suited to investigate protein-membrane interactions. They differ in size and application. The radii vary from 20 nm (small unilamellar vesicles, SUVs), 100-200 nm (large unilamellar vesicles, LUVs) up to $100~\mu m$ (giant unilamellar vesicles, GUVs). 61-63

Small unilamellar vesicles

For the preparation of small unilamellar vesicles (SUVs) lipid films with 0.4 mg of lipid material were used. These films were obtained by merging different lipid stock solutions (dissolved in chloroform; c = 2-10 mg/mL) in a test tube. This enabled to prepare lipid films with defined lipid compositions. After removing the chloroform in a nitrogen flush, the films were dried under vacuum at 30°C. The films were then stored at 4 °C until use.

The films were rehydrated with citrate buffer (30 min) and subsequently the test tubes were vortexed three times for 30 s in a five-minute interval, resulting in the formation of multilamellar vesicles (MLVs). The MLVs were then treated in an ultrasonic bath for 30 min at RT to obtain SUVs.

Giant unilamellar vesicles

Giant unilamellar vesicles (GUVs) with a diameter of $\leq 1~\mu m$ were prepared by an electro-formation process, first described by Angelova *et al.*⁶⁴ A mixture of lipids dissolved in chloroform (c = 0.5~mg/mL) were added on two indium tin oxide (ITO) covered glass slides. The chloroform was removed under reduced pressure for at least 30 min, resulting in a lipid film on top of the ITO. These ITO slides were then assembled to a chamber and sealed with a silicon ring and two Teflon spacers. Afterwards a sucrose solution (298 mOsmol/kg) was filled into the chamber. The connection to the generator was achieved with copper stripes, so each ITO slide was linked to one pole (Figure 3.1). A sinusoidal alternating current voltage of 1.6 V (peak-to-peak) and 10 Hz was applied for 3 h, resulting in GUV formation. After collection of the GUVs they were stored at room temperature for a maximum 3 days.

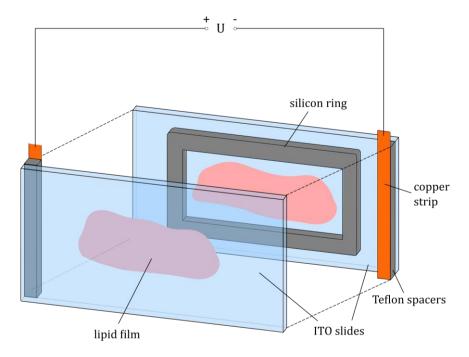


Figure 3.1: Schematic drawing of a GUV electro-formation chamber. On each ITO slide a self-adhesive copper stripe is placed via a Teflon spacer. The Teflon spacers connect the chamber with a voltage source, which results in the formation of GUVs.

3.2.2 Substrate surface preparation

Various experimental methods require different substrates, which vary in material or size. Due to this, surface functionalization strategies can be used to obtain stable membrane models. In this work two different functionalized substrate types were employed to generate supported lipid membranes and adhered giant unilamellar vesicles.

Lipid bilayer and monolayer on silicon dioxide wafers

Silicon wafers coated with silicon dioxide (SiO_2) from *Silicon Materials, Inc.* (PA, USA) were cut into 1.9 cm x 0.8 cm rectangles. For reflectometric interference spectroscopy (RIfS) experiments wafers with 5000 nm SiO_2 layer thickness and for atomic force microscopy (AFM) measurements wafers with 100 nm SiO_2 were used.

For both techniques the substrates were hydrophilized with an aqueous ammonia hydrogen peroxide solution (H_2O/NH_3 (25%)/ H_2O_2 (30%) 5:1:1 (v/v/v)) for 30 min at 70 °C. Hydrophilized substrates were rinsed with ultrapure water and stored in ultrapure water. Before use the substrates were dried in a nitrogen flush and treated with oxygen plasma (30 s, 0.2 mbar, 60 % power). Afterwards, SUVs were spread on the hydrophilic surfaces, resulting in a lipid bilayer.

Lipid monolayers were prepared on hydrophobic silicon wafers. Therefore, the wafers were first cleaned with Hellmanex (15 min) and ultrapure water (2 x 15 min) in an ultrasonic bath. Then they were treated with oxygen plasma (30 s, 0.2 mbar, 60% power) and subsequently incubated with 1,1,1,3,3,3-Hexamethyldisilazane (HMDS). Incubation was performed overnight in a sealed chamber at 120°C and under reduced pressure, yielding hydrophobic surfaces (Figure 3.2).

Figure 3.2: Hydrophobic functionalization of silicon dioxide wafers with HMDS.

Silanization and PEGylation to generate adhered GUVs

First, glass slides had to be functionalized with NeutrAvidin to obtain immobilized biotinylated giant unilamellar vesicles (GUVs). The glass slides (24 x 50 mm, ThermoFisher Scientific Gerhard Menzel, Brunswick, Germany) were cut into 1.0 cm x 1.0 cm substrates, cleaned with ethanol p.a. and ultrapure water. After treatment with oxygen plasma (30 s, 0.2 mbar, 60 % power) pure 3-glycidyloxi-propyltrimethoxysilane (GOPTS, stored in an argon atmosphere) was added between two substrates and incubated for one hour at 80 °C in a glass weighing bottle with an argon atmosphere. In that time, a mixture of methoxy- and biotin-functionalized PEGs (1:1, 2:1 and 3:1) was heated in a thermomixer at 85-95 °C (dependent on the PEG-length) and 1400 rpm. The substrates were separated from each other and then rinsed with acetone. Drying in a nitrogen flush and placing the substrates on a preheated aluminum block (80°C) enabled the addition of 200 μL of the molten PEG-mixture between two substrates without solidification. Afterwards, the substrates were placed in the weighing bottle and incubated at 85-95 °C for 4 h. The substrates were separated, intensely rinsed with ultrapure water and dried in a nitrogen flush. If the molten PEG mixture was still present on the substrates (impure surface), the rinsing step was repeated. Prior to use the substrates were stored under argon at 4 °C. The silanization and PEGylation is illustrated in Figure 3.3.

<u>silanization:</u> <u>PEGylation:</u>

Figure 3.3: Silanization and PEGylation of glass substrates, resulting in methoxy-PEG and biotin-PEG surfaces.

To obtain adhered GUVs, the functionalized substrates were first incubated with NeutrAvidin (50 μ g/mL in PBS) for 30 min. Then the NeutrAvidin-functionalized substrates were rinsed six times with PBS and further six times with sucrose buffer. 10-50 μ L of a biotin-containing GUV suspension (in sucrose) was added and incubated for 15 min in a humidity chamber to avoid osmolar changes. Via biotin-NeutrAvidin interactions the GUVs adhered to the surface.

To verify a complete coverage of the substrates with NeutrAvidin, the surface was incubated with DyLight® 594 labeled Neutravidin/unlabeled NeutrAvidin in a ratio of 99:1.

3.3 Biophysical methods

3.3.1 UV-Vis spectroscopy

The determination of ENTH solutions' concentration was done using UV-Vis spectroscopy. This method is based on the absorption of ultraviolet (UV) and visbible (Vis) light.⁶⁵ Proteins with aromatic amino acid residues (Phenylalanine, Trypthophan, Tyrosine) exhibit an absorption maximum at $\lambda = 280$ nm due to $\pi \rightarrow \pi^*$ transition within

the amino acids. Measurements of the absorption A_{280} at this wavelength allowed the calculation of the protein concentration via the Lambert-Beer law:

$$A_{280} = \varepsilon_{ENTH} \cdot c \cdot d \tag{3.1}$$

The extinction coefficient of the ENTH-domain is $25440 \,\mathrm{M}^{-1}\mathrm{cm}^{-1}$ and the cuvette thickness is $d=1 \,\mathrm{cm}$. The molecular masses and extinction coefficients of ENTH were calculated with the ExPASy ProtParam tool.⁶⁶ The measurements were performed using the NanoDropTM 2000c (Thermo Fisher Scientific, Waltham, MA, USA).

3.3.2 Reflectometric interference spectroscopy

Reflectometric interference spectroscopy (RIfS) was used to determine the binding of the ENTH domain to lipid bilayers. It is a label-free optical biosensing method, which detects time-resolved changes in optical thicknesses at interfaces between two media with different refractive indices (n).^{67,68} The method is based on the reflection and transmission of white light at these interfaces ⁶⁹ A light beam is directed to the interface, where it is partially reflected and refracted. Superimposition of the partial beams leads to a specific interference pattern. Changes of the surface thickness (physical thickness d) due to e.g. protein adsorption lead to an altered interference spectrum.

The law of Snellius describes the refraction of light on interfaces with the incident angle α , the refracting angle φ and the different refractive indices n (equation(3.2)).

$$n_1 \cdot \sin(\alpha) = n_2 \cdot \sin(\varphi) \tag{3.2}$$

Depending on the refractive indices of the media, the reflection and transmission of the light changes. In Figure 3.4 three possible light pathways are illustrated.

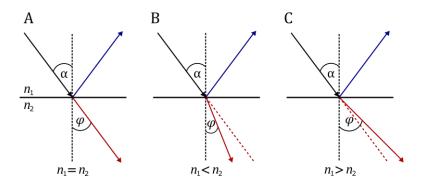


Figure 3.4: Schematic illustration of the SNELLIUS reflaction law. With an incident angle the light is directed to a second medium. Due to the different refractive indices of both media the light can have various light pathways. In case of $n_1 = n_2$ the light is reflected and reflected with the same angles ($\alpha = \varphi$). When the refractive index of the first medium is higher than the second ($n_1 > n_2$), the light is refracted to the vertical ($\alpha > \varphi$). If the relation of $n_1 < n_2$ occurs, the light is tilted out of the vertical ($\alpha < \varphi$).

If the refractive indices n_1 and n_2 have the same value (Figure 3.4 A), the angles α and φ are equal, then the light is reflected and transmitted in the same degree. With $n_1 < n_2$, the light is refracted when it is passing the second medium and the relation $\varphi < \alpha$ is valid (Figure 3.4 B). In the latter case, the transmitted light is tilted out of the vertical due to the higher refractive index of the second medium ($n_1 > n_2$, Figure 3.4 C).

In the experimental setup, utilized in this work, the light is directed to a three phase system: An opaque silicon layer followed by a silicon dioxide layer, which is surrounded by an aqueous phase (Figure 3.5).

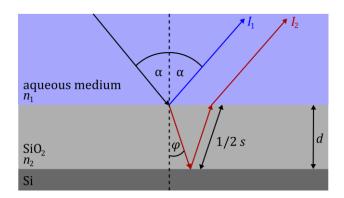


Figure 3.5: Schematic light pathway during a RIfS experiment. The light is directed to a system consisting of three phases with an incident angle α . Reaching the silicon dioxide layer with a thickness d after passing the aqueous medium, the light is reflected (I_1 , blue) and refracted (I_2 , red) to the vertical at an angle φ due to different refractive indices (I_2). As the silicon layer is intransparent the light is completely reflected, resulting in an additional traveling distance I_2 .

With the incident angle α the light passes the aqueous medium (n_1 = 1.33 at 20 °C).⁷⁰ When it is hitting the silicon dioxide layer (n_2 = 1.45)⁷¹ the light is refracted towards the normal at the angle φ . The opaque silicon layer reflects the light beam, which is then refracted once more at the interface of media with n_1 and n_2 . The result is a phase shift $\Delta \varphi$ due to the longer traveling distance (s) of the refracted light (I_2) compared to the reflected part (I_1). The distance can be calculated with equation (3.3).

$$s = \frac{2d}{\cos(\varphi)} \tag{3.3}$$

The optical distance is dependent on the refractive index of the second medium n_2 (equation (3.4)), which leads to equation (3.5).

$$s_{OT} = n_2 \cdot s \tag{3.4}$$

$$s_{OT} = \frac{2n_2d}{\cos(\varphi)} \tag{3.5}$$

The refraction of the light causes a phase shift, resulting in constructive and destructive interference. When constructive interference arises the $\Delta \varphi$ is a multiple of the wavelength, described by equation (3.6).

$$\Delta\varphi(\lambda) = \frac{2\pi \cdot OT}{\lambda} \cos(\varphi) \tag{3.6}$$

The optical thickness changes by adsorption of molecules on the surface as it is the product of n and d (equation (3.7)). Since the refractive indices of silicon dioxide and lipid or protein layers are similar,⁷² it hence enables to measure the molecule adsorption by detection of the optical thickness changes.

$$OT = n \cdot d \tag{3.7}$$

In case of a perpendicular incident light causing $\alpha = \varphi = 0^{\circ}$, equation (3.7) simplifies to equation (3.8). This allows to monitor the optical thickness change via phase shift detection.

$$\Delta\varphi(\lambda) = \frac{2\pi \cdot OT}{\lambda} \tag{3.8}$$

Evaluation of the data

A reflectivity spectrum was determined by the program Spectra Suite (Ocean Optics Inc., Dunedin, FL, USA) with the reference (I_{Γ}) and the dark ($I_{\rm d}$) spectrum according to equation (3.9). The reference spectrum is recorded at nearly complete reflectivity against a silver-coated surface. Measurement of the light intensity without illumination yielded the dark spectrum.

$$R(\lambda) = \frac{I_{\rm m} - I_{\rm d}}{I_{\rm r} - I_{\rm d}} \tag{3.9}$$

 $I_{\rm m}$ is the measured intensity spectrum. The reflectivity is described with the Fresnel reflectivity coefficients r_1 and r_2 in equation (3.10).

$$R(\varphi) = \frac{r_1^2 + 2r_1r_2\cos(2\varphi) + r_2^2}{1 + 2r_1r_2\cos(2\varphi) + r_1^2r_2^2}$$
(3.10)

Thereby, the Fresnel reflectivity coefficients are defined as:

$$r_1 = \frac{n_1 - n_2}{n_1 + n_2}$$
 and $r_2 = \frac{n_2 - n_3}{n_2 + n_3}$ (3.11)

Combination of equation (3.10) and (3.11) leads to the reflectivity, which is dependent on the wavelength and the optical thickness OT. This in turn enables the timeresolved tracking of OT (equation (3.12)).

$$R(\lambda, OT) = \frac{r_1^2 + 2r_1r_2\cos\left(\frac{4\pi}{\lambda}OT\right) + r_2^2}{1 + 2r_1r_2\cos\left(\frac{4\pi}{\lambda}OT\right) + r_1^2r_2^2}$$
(3.12)

Experimental procedure

The hydrophilic or hydrophobic functionalized silicon wafers (cf. chapter 3.2.2) were implemented into a RIfS chamber, consisting of an aluminum bottom, an acrylic glass cap and a rubber ring to seal the chamber (Figure 3.6).

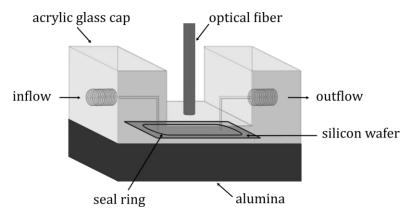


Figure 3.6: Schematic illustration of a RIfS chamber. A silicon wafer is mounted between an acrylic glass cap and an alumina plate. In and outflow of liquids is realized by inlets in the acrylic glass cap. Then the optical fiber is placed in the indentation of the acrylic glass cap.

The flow in and out of the chamber was conducted via inlets inside the acrylic glass cap, which enabled a controlled addition of liquids. A fanned optical fiber was orthogonally mounted to the wafers. This fanned fiber is made of one light-collecting fiber, which is surrounded by six emitting fibers connected to a tungsten halogen lamp ($\lambda = 350 - 950$ nm). The reflected and refracted light of the 1 mm² lighted area is detected by a UV/Vis-silicon-CCD-detector (2048 Pixel).

Every two seconds a reflectivity spectrum was recorded with the software *Spectra Suite* and then tracked with a MATLAB script using equation (3.12) written by Dr. I. Mey (Georg-August-University Göttingen).

For the experiment, the system was first rinsed with citrate buffer, which was defined as the baseline for the determination of the optical thickness changes. Then SUVs were added, leading to an increase of the optical thickness. This increase indicates the spreading and formation of a lipid bilayer (hydrophilic) or monolayer (hydrophobic). After a plateau was reached, PBS was rinsed through the chambers to remove nonspread SUVs. Then, BSA (1 mg/mL) was added for surface passivation, preventing possible unspecific interactions of ENTH with defects in the membranes. Before protein addition, the system was again rinsed with PBS, the same buffer in which the ENTH domain is solved. For the investigation of the influence of PIP2 on the ENTHmembrane interaction, measurements with a single protein addition with a concentration of 1 µM were implemented. Isotherms with gradual protein additions (c (ENTH) = 0.1-5 μ M)) were performed to detect the influence of the charged lipid Phosphatidylserine (PS) on the ENTH binding affinity. Due to the interaction of the membrane/surface and the added protein, an increase of the optical thickness was monitored. When a plateau was reached, PBS was rinsed through the system to detect the protein dissociation process or more protein was added to obtain a Langmuir adsorption isotherm.

3.3.3 Fluorescence microscopy

Fluorescence microscopic techniques were used to image solid-supported membranes and adhered GUVs. This non-invasive method bases on the physical phenomenon of fluorescence. After the absorption of light, fluorophores emit the light with a longer wavelength (stokes shift). Separation with filters and mirrors allows the collection of the emitted light, excluding illumination light. There are different types of fluorescence microscopy setups, which enable to image labeled biological structures in high resolution. In the next chapters two types of microscopes are described in more detail.

3.3.3.1 Confocal laser scanning microscope

The confocal laser scanning microscope (CLSM) uses a single laser beam, which is focused onto a specimen. As the emitted light is filtered through a pinhole rejecting out-of-focus light, the sample is scanned point-wise. Therefore, higher axial resolutions can be reached compared to an epifluorescence microscope.⁷³ Due to the illumination of a small focal volume of the sample moreover the bleaching of fluorophores is reduced. In Figure 3.7 an illustration of the setup and the light path of a CLSM is shown.

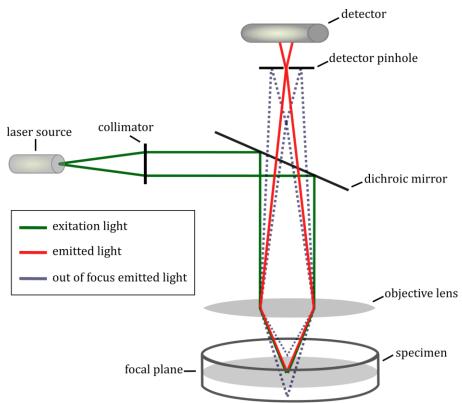


Figure 3.7: Schematic illustration of a confocal laser scanning microscope setup. The fluorophore of the specimen is excited by a laser. The emitted light (red) passes the dichroic mirror and is focused to the detector via a confocal pinhole, which enables to reject out-of-focus light (blue dotted lines).

The resolution of a CLSM using very small pinholes can be determined by the full width at half maximum (FWHM) of the point spread function according to Abbes law, enabling lateral and axial resolutions in the range of 200-250 and 570-950 nm⁷⁴ when the specimen was exited with visible light (equations (3.13) and (3.14)).

$$d_{\text{FWHM,lateral}} = \frac{0.37 \ \overline{\lambda}}{NA} \tag{3.13}$$

$$d_{\text{FWHM,axial}} = \frac{0.67\,\overline{\lambda}}{n - \sqrt{n^2 - NA^2}} \tag{3.14}$$

 $\bar{\lambda}$ is described by the wavelength of the emitted ($\lambda_{\rm em}$) and the exited light ($\lambda_{\rm ex}$) with $\sqrt{\lambda_{\rm em} \cdot \lambda_{\rm ex}}$, n is the refractive index of the imaged medium and NA the numerical aperture of the objective.

The CLSM setup was used to image solid-supported membranes as well as to determine diffusion coefficients of these membranes via fluorescence recovery after photobleaching experiments.

Fluorescence recovery after photobleaching

Fluorescence recovery after photobleaching (FRAP) is a method to determine the lateral mobility of labeled molecules. First, the fluorescence intensity of a defined circular region of interest (ROI, Figure 3.8, red circle) of the sample is measured. Subsequently, fluorophores in this region are irreversibly bleached by a laser beam. As far as the lipids are mobile, the fluorescence intensity in the ROI increases after a time due to diffusion of unbleached molecule into the ROI and diffusion of bleached molecules out of the ROI. In Figure 3.8 a typical FRAP experiment is illustrated.

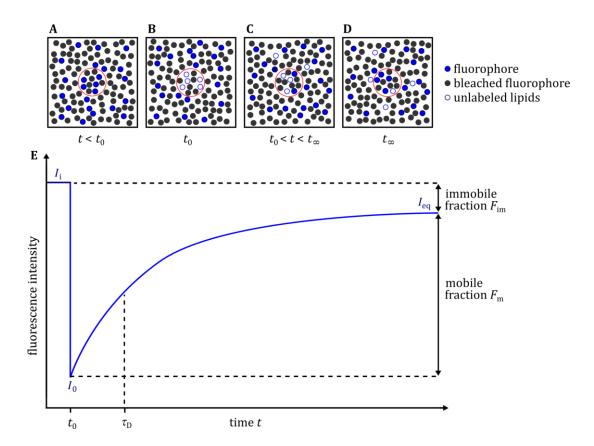


Figure 3.8: Schematic illustration of a FRAP experiment. (A) At the beginning the fluorescence intensity I_1 in a circular region of interest (ROI, red circle) is detected. The lipids are distributed homogenously (B) Bleaching of labeled lipids in the ROI leads to the decrease of the fluorescence intensity to I_0 in corresponding time trace (E). (C) Due to lateral mobility labeled lipids diffuse out of the ROI and non-labeled lipids into the ROI, which increases the fluorescence intensity in the ROI. (D) At the end the a new equilibrium of the fluorescence intensity I_{eq} is reached.⁷⁵

The graph in Figure 3.8 E shows an exponential increase of the intensity after bleaching of labeled lipids in the ROI. The time dependent fluorescence intensity I(t) can be described with the fluorescence intensity directly after the bleaching I_0 , the intensity at the new equilibrium I_{eq} , the time t and the time constant of the fluorescence recovery t_F (equation (3.15)).^{76,77}

$$I(t) = I_{\text{eq}} - \left(I_{\text{eq}} - I_{0}\right) \cdot \exp\left(\frac{-t}{t_{\text{F}}}\right)$$
(3.15)

The time constant is related to the characteristic diffusion time τ_D by equation (3.16).

$$\tau_{\rm D} = \ln (0.5) \cdot (-t_{\rm F})$$
 (3.16)

With the radius of the ROI (r) and τ_D it is possible to calculate the diffusion coefficient D using equation (3.17).

$$D = \frac{r^2}{4\tau_{\rm D}} \tag{3.17}$$

Dependent on the interactions of the substrate and the membrane, a certain amount of the lipids remains immobile. The mobile fraction F_m can be determined with equation (3.18).

$$F_{\rm m} = \frac{I_{\rm eq} - I_0}{I_{\rm i} - I_0} \tag{3.18}$$

describes the initial fluorescence intensity before bleaching. The diffusion coefficients and the mobile fraction were determined with a MatLab Script written by Jönsson et al. using a Hankel transformation.⁷⁸

Experimental procedure

The experiments were performed with an upright microscope (LSM 880 Examiner, Carl Zeiss, Jena, GER) equipped with a water immersion objective (WPlan-Apochromat 63x, NA 1.0, Carl Zeiss, Jena, GER). SUVs were prepared as described in chapter 3.2.1. To enable the visualization of the membranes, a fluorophore was incorporated. In this thesis the lipid bound dye TxR was used. TxR was excited at 561 nm and detected at 570-680 nm. After functionalization of the substrates (cf. chapter 3.2.2) they were implemented into Teflon chambers. Depending on the functionalization, spreading of SUVs for 1 h resulted in mono- (hydrophobic) or bilayers (hydrophilic). Nonspread vesicles were removed by rinsing with citrate buffer (3x) and PBS (3x). Subsequently, imaging and FRAP experiments were performed.

3.3.3.2 Spinning Disc

Compared to a conventional CLSM a spinning disc confocal laser microscopy (SDCLM) has the advantage that imaging works much faster due to multiple beam illumination. This is realized with two synchronically rotating discs. Excitation light first passes the micro lens array disc and is then focused onto the sample by a second Nipkow disc with multiple (about 1000-20000) spirally arranged pinholes (50 μ m). The SDCLM setup was used to image z-stacks of biotinylated vesicles adhered to a NeutrAvidin functionalized surface. In Figure 3.9 the setup of a SDCLM is illustrated.

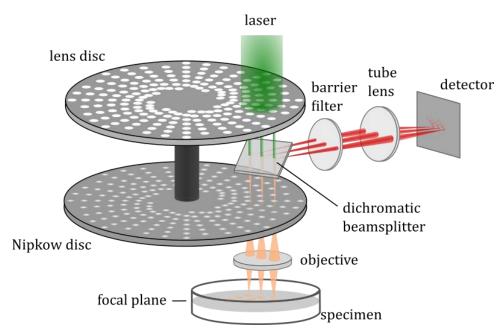


Figure 3.9: Setup of a spinning disc. Laser light is directed to a specimen by the micro lens disc passing through pinholes of a second disc (Nipkow disc). The emitted light is focused to the beamsplitter by the objective lens through the Nipkow disc pinholes, where the light is directed to the detector, passing a barrier filter and a tube lens. The figure was modified according to Gleisner.⁷⁹

Experimental procedure

The measurements were performed using an upright SDCLM (spinning disc: Yokogawa CSU-X, Rota Yokogaa GmbH & Co. KG, Wehr, stand: custom made Olympus Deutschland GmbH, Hamburg) equipped with an iXON 897Ultra camera (Andor Technology Ltd., Belfast, UK) and a water immersion objective (CFI Plan 100XW 100x, *NA* 1.1, Nikon, Düsseldorf). The dye Atto488-DPPE, incorporated in the GUVs, was excited at 488 nm. A LP496 ET filter was used to detect the emission of the fluorophore.

Adhered GUVs were prepared as described in chapter 3.2.2. The use of a humidity chamber avoided the evaporation of the buffer and consequently ensured iso-osmolar conditions during the experimental time. A series of z-stacks of the adhered GUVs were recorded with 5-13% laser power and with an exposure time of 20 ms. The distance between the slices was adjusted to 150-250 nm. The radii of the GUVs and the contact areas of the GUVs with the surface were determined with ImageJ.

3.3.4 Atomic force microscopy

Binning *et al.* first introduced the atomic force microscopy (AFM) technique in 1985.⁸⁰ It enabled scanning of non-conductive probes under physiological conditions with a high lateral (1.0 nm) and axial (0.1 nm) resolution.⁸¹ Hence it is well suitable for the analysis of biological systems, like e.g. receptor-ligand interactions.⁸² In this work AFM was used to analyze surface topographies of solid-supported membranes and ENTH clusters.

Principle of an AFM

A cantilever with a sharp tip is moved over a sample by x-y piezo actuators. Furthermore, a z-piezo actuator can approach the cantilever towards the sample until attractive or repulsive interactions of the sample and the cantilever lead to the deflection of it. The deflection is detected by a laser beam, which is reflected by the cantilever surface as well as by a mirror and then directed to a four quadrant diode. Lateral or vertical deflection results in the shift of the laser beam on the diode. Hence the deflection of the cantilever against the sample position allows to get a topographic image of the sample surface. The setup of an AFM is illustrated in Figure 3.10.

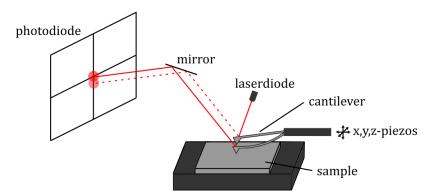


Figure 3.10: Schematic illustration of an AFM setup. The piezos move the cantilever on the surface. A light beam, which is directed to the cantilever, is reflected and directed to the four quadrant photodiode, detecting the deflection of the cantilever.

In AFM different imaging modes can be applied. In this work the contact mode was used, where the tip is in close contact to the sample. The contact mode can be realized in constant height or constant force. In case of constant force a topographic image can be illustrated due to the piezo element regulation. This mode appeared to be best suited for imaging ENTH clusters on solid-supported membranes.

Besides topographical information, AFM is utilized for the determination of mechanical properties of samples by force-distance curves. This method is often used to measure the elastic characteristics of biological systems.⁸³ Thereby, the force can be obtained by the vertical cantilever defection Z_c and the spring constant of the cantilever (k) by Hooke's law (equation (3.19)).⁸⁴

$$F = k \cdot Z_c \tag{3.19}$$

For force-distance curves the position of the piezo Z_p and the cantilever deflection Z_c are converted into force F and distance D (Figure 3.11 B). Applying a fit to the compliance region yielded the slope, which is equal to the conversion factor of the cantilever defection and the detector signal. The tip-sample distance can be obtained by equation (3.20).85

$$D = Z_{\rm p} + Z_{\rm c} \tag{3.20}$$

Force-distance curves were detected at single sample positions and predefined forces of 6-8 nN to obtain breakthrough forces which provides the membrane thicknesses as they are correlated (Figure 3.11).

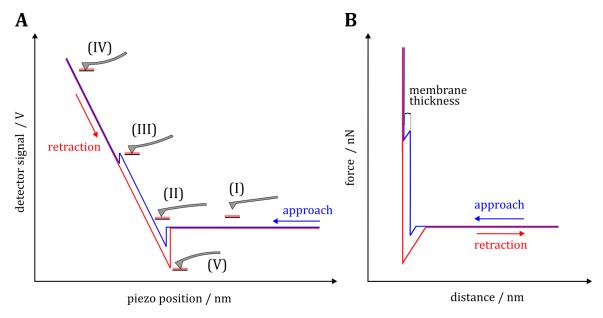


Figure 3.11: Schematic illustration of a signal versus piezo position (A) and a force-distance curve (B). During approach (blue) no interactions of cantilever and sample occur (I). Due to attractive forces the cantilever gets in contact with the sample and the measured force suddenly decreases (II). Further approach leads to increase of the force (IV). At high predefined forces the cantilever can break through the lipid bilayer resulting in a second drop of the force (III). Until the retraction energy exceeds the adhesion energy the cantilever stays in contact to the surface during retraction (red, V). At the end the baseline is reached again as no contact of cantilever and sample is present.

By the approach of the cantilever the force suddenly decreases due to attractive electrostatic or van-der-Waals interactions. As the cantilever is in contact further approach results in the increase of F until a predefined value. At high predefined forces a small drop in the approach force-distance curve can appear (cantilever tip breaks through the bilayer), enabling the determination of the membrane thickness. In the force-distance plot the membrane thickness is the distance between the decrease of F and the point where F begins to rise again (cf. Figure 3.11 B). Then retraction of the cantilever leads to the decrease of F. An even negative force value can be reached when adhesion causes cantilever bending. With further retraction the cantilever loses the contact as the retraction force exceeds the adhesion force.

Experimental setup

For the supported lipid bilayers (SLBs) previously described FRAP and CLSM experiments (cf. chapter 3.3.3.1) were done to verify the mobility of the membranes. Measurements were performed using a JPK Nanowizard 4 (JPK Instruments, Berlin, GER). After fixing the cantilever to a glass holder and implementing both into the AFM head, the laser was directed to the tip of the cantilever to get a maximal signal. The exact spring constant of the cantilever was determined by measuring the thermal noise spectrum.⁸⁷

As all settings were adjusted, micrographs of SLBs were taken to analyze the surface topography before and after protein adsorption. SLBs were prepared as described in chapter 3.2.2. Measurements were done in contact mode using BL-AC40TS-C2 cantilevers (BioLever mini, f = 85.4-139.1 kHz, k =0.03-0.12 N/m, Olympus). First of all break-through experiments were performed to measure the membrane thickness as it can be derived from the break-through force. Besides FRAP experiments, this ensured that a bilayer with a thickness of about 4 nm was formed, which is typical for a bilayer.⁸⁸Next, an area of $10 \times 10 \ \mu\text{m}^2$ of the SLBs was imaged. After incubation with $1 \ \mu\text{M}$ ENTH or ENTH R114A mutant for 2h at RT, the surface was scanned again. Roughnesses (root-mean-square, rms) were determined with the integrated JPK Data Processing software.

The protein height and occupancy were analyzed using a MatLab script written by Dr. Ingo Mey (Georg-August Universität Göttingen). To detect the membrane and the protein a threshold was set. Protein adsorption was identified by a 2D peak detection function, which marked the local maxima. A histogram out of the maxima was created and fitting a normal distribution yielded the protein height.

3.3.5 Langmuir-Blodgett trough

Monolayers have the advantage of a homogenous lipid distribution without the feasible asymmetry between two leaflets of bilayers. One possibility to generate monolayers without the necessity of a solid support is the Langmuir-Blodgett trough. This technique thus enables to prevent interactions of lipids with substrate surfaces. Besides the determination of lipid concentrations and properties, also the compressibility of monolayers can be analyzed. Furthermore, single or multiple monolayers on solid supports (so called Langmuir films) can be generated with this technique. Studies with spread lipids on water already were done by Agnes Pockels in 1882.⁸⁹ However, the technique as it is known and used today was extended and improved by Irving Langmuir and Katherine Blodgett in 1917 and 1934.

The Langmuir-Blodgett trough or film balance consists of a temperature controlled Teflon trough, Teflon barriers and a microbalance system with a Wilhelmy plate.⁹⁰ These barriers can be moved by a control system to change the trough area and enable the compression or expansion of spread films (Figure 3.12).

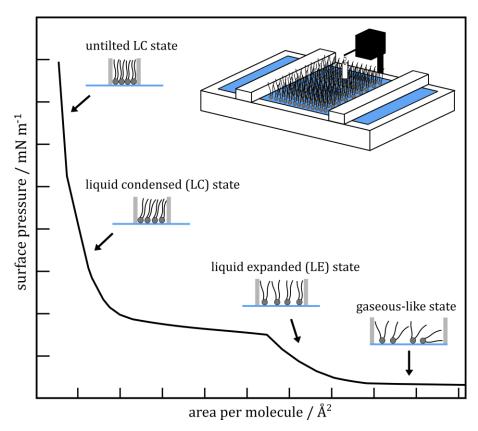


Figure 3.12: Schematic illustration of a Langmuir-Blodgett trough and a typical surface pressure-area per molecule isotherm. Due to compression the area per molecule gets smaller and the lipids change their orientation. The transition from a gas-like phase to a liquid-expanded state occurs. Further compression leads to a solid-like or untilted LC state via a liquid condensed (LC). The Figure was modified according to M Larsen.⁹¹

With a large trough area the lipid molecules were comparable to a 2D gas. 92,93 By reducing the area the lipids switch to a liquid-expanded (LE) state. At a certain surface pressure a co-existence region of the LE and the liquid condensed (LC) phase can occur, induced by a first-order transition, followed by the LC state, where the lipids are highly ordered. Further compression leads to the collapse of the film due to too high packing densities. Here the appearance of multilayers or the formation of micelles are possible.

Determination of the surface pressure

When amphiphilic molecules, like lipids, get in contact with a water-air interface a rearrangement of the lipids occurs. The hydrophilic head groups will get in contact

with the water (subphase) and the lipophilic tail protrudes into the air. This perpendicular alignment is associated with the decrease of the free energy of the system. An enrichment of the molecules causes the favored expansion of the interface with an interface or surface pressure π . The surface tension of the water $\gamma_{\rm H_2O}$ (72.5 mN/m) counteracts, resulting in a reduced surface tension. The surface pressure is hence the difference of the surface tension without $\gamma_{\rm H_2O}$ and with the lipid film $\gamma_{\rm Film}$ (equation (3.21)). 96

$$\gamma_{\text{Film}} = \gamma_{\text{H}_2\text{O}} - \pi \leftrightarrow \pi = \gamma_{\text{H}_2\text{O}} - \gamma_{\text{Film}}$$
 (3.21)

The resulting surface tension γ_{Film} was determined with the Wilhelmy plate method. Figure 3.13 shows the principle of this technique.

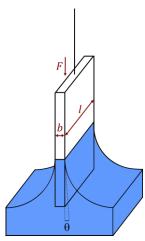


Figure 3.13: Wetting principle of a Wilhelmy plate. A force F acts on the Wilhelmy plate, which is proportional to the wetted perimeter L = (2l + 2b). l and b describe the plate width and thickness. The plate is wetted at an angle θ .

By wetting of a Wilhelmy plate a force *F* acts on this plate according to equation (3.22).

$$F = \gamma \cdot L \cdot \cos \Theta \leftrightarrow \gamma = \frac{F}{L \cdot \cos \Theta}$$
 (3.22)

This force correlates with the surface tension, the wetted perimeter L = (2l + 2b) (cf. Figure 3.13) and the contact angle (Θ). Wilhelmy plates are often made of filter papers, ensuring complete wetting ($\Theta = 0^{\circ}$), which simplifies the equation (3.22) to:

$$\gamma = \frac{F}{L} \longrightarrow \gamma_{\text{Film}} = \frac{F}{L} \tag{3.23}$$

The calculation of the surface tension γ , or more exactly the surface tension with the film γ_{Film} , then enables the determination of the surface pressure π (cf. equation (3.21)), which is displayed on the control system.

Experimental procedure

Before the experiment was started, the trough and the barriers were cleaned with Mucasol for 15 min. In this time the thermostat (20 °C) and the software were turned on and the filter paper was soaked in ultrapure water. After rinsing with ultrapure water (trough 4 times, barriers 10 times), the trough was filled with 120 mL subphase. Lipid concentrations were determined in ultrapure water and the investigation of the monolayer stability was performed in PBS. For measurements with a protein trough, first 50 mL PBS were added. The protein trough was then fixed with screws and further 70 mL PBS were filled into the middle of the trough. The barriers were fixed and the filter paper was dipped into the subphase to get a slight contact. During this contact the *Offset* was adjusted to 0 mN/m. Without contact to the water, the Calibration-screw was set to 72.75 mN/m ($\gamma_{\rm H_2O}$).95 This procedure was repeated three times for precise calibration.

The lipid mixture, dissolved in chloroform, was carefully added to the subphase with a Hamilton syringe, which was rinsed three times with chloroform before use. The system was then equilibrated for 15 min to remove the chloroform by vaporization. After adjusting the barrier velocity (4.7 mm²/s) and starting the measurement in the software, the barrier motors (setup: "auto") were turned on.

For determination of the lipid concentrations via an isotherm the trough area was compressed up to a surface pressure of π = 25 mN/m. For the isotherms of the lipid mixtures the surface pressure was set to π = 35 mN/m. To investigate the monolayer stability of the used lipid mixtures a maximum surface pressure of π = 30 mN/m was adjusted. When this surface pressure was reached, the barrier motors were turned

off to get a constant trough area. The surface pressure is then plotted against the time to see if the monolayer remained stable over one hour.

Protein insertion experiments were done with a special protein trough (Figure 3.14). This protein trough enables the insertion of ENTH into the subphase to investigate the surface activity in a smaller volume (V = 2 mL).

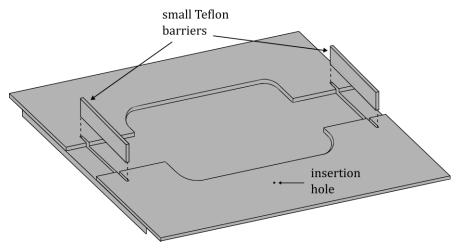


Figure 3.14: Schema of the used protein trough with a capacity of 2 mL, confined by small Teflon barriers. Through a small hole protein can be inserted into the subphase.

For this experiment different initial surface pressures were set. When the desired pressure was achieved, the barrier motors were turned off and small Teflon barriers were fixed into the pre-built notches. As the monolayer is disturbed during the inset of the small barriers, the system was equilibrated for 5-10 min until no more change of the surface pressure was observed. Then 1 μ M of the ENTH-domain was added to the subphase and subsequently distributed with a Hamilton syringe. The measurements were recorded until a maximum of the surface pressure change was reached. In this work following lipid mixtures were used to analyze the ENTH-PIP2 interaction with the film balance technique.

- POPC/POPE 80:20
- POPC/POPE/PIP₂ 79:20:1
- POPC/POPE/PIP₂ 75:20:5
- POPC/POPE/PIP₂ 70:20:10
- POPC/POPE/POPS 60:20:20
- POPC/POPE/POPS/PIP₂ 55:20:20:5

4 RESULTS

In this work, the impact of negatively charged lipids phosphatidylinositol-4,5-bisphosphate (PIP₂) and phosphatidylserine (PS) on the interaction of the epsin *N*-terminal homology (ENTH) domain and membranes was analyzed. First, the binding behavior on solid supported bi- and monolayers was investigated by means of reflectometric interference spectroscopy (RIfS). Furthermore, atomic force microscopy (AFM) experiments were performed to obtain the surface topology of membranes after ENTH binding. Finally, the penetration capability of ENTH and its influence on membrane tension were determined using the Langmuir-Blodgett technique and spinning disc confocal laser microscopy (SDCLM).

4.1 Influence of the PIP₂ concentration on the ENTH binding to solid supported lipid membranes

In the clathrin-mediated endocytosis the receptor lipid PIP₂ is necessary for the recruitment of ENTH to the membrane and consequently involved in one key step of this process (cf. chapter 1.1).^{8,21,25} It is predicted that ENTH binds PIP₂ in a 1:1 stoichiometry,⁷ which would assume a direct correlation of the PIP₂ content and the amount of bound protein. The impact of PIP₂ and especially its concentration on the ENTH-membrane interactions is still unexplored. To analyze how the concentration of the receptor lipid influences the protein occupancy on PIP₂ doped membranes, RIfS experiments were performed. Therefore, the amount of bound protein on solid supported bilayers and monolayers at different PIP₂ concentrations was determined.

4.1.1 Formation of solid-supported lipid bilayers on silicon dioxide

Since the formation of a homogenous and planar membrane is required to investigate the ENTH binding by means of RIfS, the spreading of small unilamellar vesicles (SUVs) on hydrophilic silicon dioxide (SiO₂) was considered first. When vesicles of POPC/PIP₂ with different receptor lipid concentrations (0-10 mol%) were added to a hydrophilic SiO₂ wafer, they spread on the surface, resulting in a lipid bilayer. During this spreading process the increase of the optical thickness (ΔOT) was observed, indicating the successful bilayer formation. In Figure 4.1 the formation of a POPC (100 mol%, A) and a POPC/PIP₂ (90:10, B) bilayer are shown.

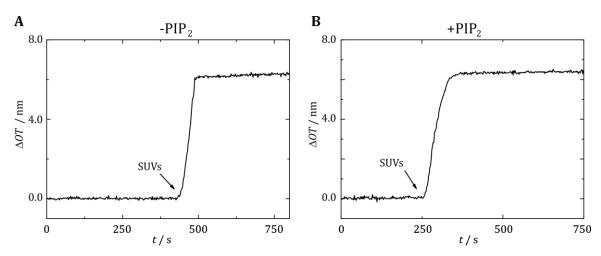


Figure 4.1: Exemplary plots of the change in optical thickness vs. the time during vesicle spreading of (A) pure POPC and (B) POPC/PIP₂ (90:10) SUVs onto hydrophilic silicon dioxide wafers. During spreading of SUVs the optical thickness increases until a plateau is reached, indicating the successful formation of a bilayer.

The baseline was generated by rinsing the system with citrate buffer at a pH of 4.8, which reduced the negative charge of the PIP₂ and the SiO₂ surface. In previous studies, it was shown that these conditions (ions, pH) improve the spreading of PIP₂-containing SUVs and hence the bilayer formation.⁹⁷

After rinsing with PBS to remove non-spread vesicles, the surface was passivated with BSA to avoid defects in the bilayer. The system was then rinsed with PBS again, yielding the ΔOT after vesicle spreading (Figure 4.2).

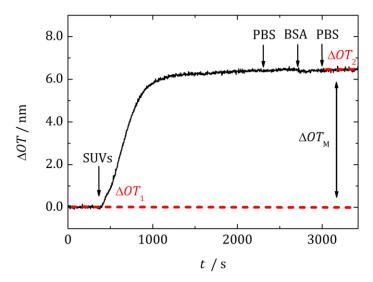


Figure 4.2: Time-trace of the spreading process of POPC/PIP₂-SUVs, followed by rinsing with PBS and BSA. PBS was added to remove adhered, non-spread vesicles and BSA for passivation. The optical thickness of the membrane ($\Delta OT_{\rm M}$) is determined by the difference of ΔOT_2 - ΔOT_1 . Here an example of the bilayer formation on hydrophilic SiO₂ is shown. The $\Delta OT_{\rm M}$ of a monolayer is calculated analogously.

To obtain the membrane thickness ($\Delta OT_{\rm M}$) the difference of the ΔOT before (baseline, ΔOT_1) and after (ΔOT_2) SUV spreading were calculated with equation (4.1).

$$\Delta OT_{\mathsf{M}} = \Delta OT_2 - \Delta OT_1 \tag{4.1}$$

From the optical thickness of the membrane ($\Delta OT_{\rm M}$) the physical thickness $d_{\rm M}$ could be determined with equation (3.7).

The lipid bilayer formation was investigated for different receptor lipid concentrations in the vesicles ($\chi(PIP_2 = 0-10 \text{ mol}\%)$). In Table 4.1 the determined values for ΔOT_M and d_M were collected.

Table 4.1: Averaged optical and physical thicknesses of the POPC/PIP₂ lipid bilayers with the corresponding standard deviations. The receptor lipid concentration was changed from 0-10 mol%. *n* number of performed experiments.

Lipid composition	$\Delta OT_{\rm M}$ / nm	<i>d</i> м / nm	n
POPC	6.4 ± 0.2	4.4 ± 0.2	5
POPC/PIP ₂ (99:1)	6.4 ± 0.3	4.3 ± 0.2	4
POPC/PIP ₂ (98:2)	6.4 ± 0.4	4.3 ± 0.3	8
POPC/PIP ₂ (95:5)	6.4 ± 0.5	4.3 ± 0.4	9
POPC/PIP ₂ (93:7)	6.3 ± 0.2	4.3 ± 0.2	4
POPC/PIP ₂ (92:8)	6.3 ± 0.4	4.3 ± 0.3	5
POPC/PIP ₂ (90:10)	6.3 ± 0.5	4.3 ± 0.3	10

The ΔOT of the membranes ranged from 6.3 to 6.4 nm and the d_M from 4.3 to 4.5 nm, not showing a significant change with increased PIP₂ concentration as it was presumed due to the larger head group of PIP₂.^{44,98} These results show that for all lipid compositions the bilayers were formed successfully.

To verify that the formed lipid bilayers are homogenous and mobile, additionally fluorophore labeled bilayers were investigated by means of confocal laser scanning microscopy (CLSM). Solid supported lipid bilayers composed of POPC, PIP₂ (1-10 mol%) and the lipid coupled fluorophore Texas Red® 1,2-dihexadecanoyl-*sn*-glycero-3-phosphoetanolamine (TxR) were prepared as described in chapter 3.2.2. Then FRAP-experiments were performed to obtain the diffusion coefficients (*D*) of TxR in the solid supported bilayers (cf. chapter 3.3.3.1). In Figure 4.3 an exemplary FRAP experiment on POPC/PIP₂/TxR (89.5:10:5) is shown.

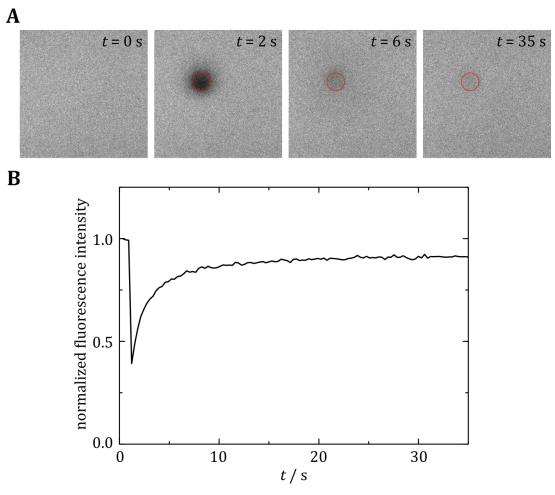


Figure 4.3: (A) Fluorescence images of a FRAP experiment and (B) the corresponding time-trace of the normalized fluorescence intensity on a POPC/PIP $_2$ /TxR (89.5:10:0.5) bilayer. After bleaching of a region of interest (ROI, red circle, t = 2 s) the fluorescence intensity increases, reaching a plateau after 20 s. Due to the immobile fraction of TxR in the bilayer, the fluorescence intensity does not reach the fluorescence intensity value at t = 0.

In all samples (1-10 mol% PIP₂) the fluorescence of TxR was distributed equally, indicating the formation of a homogenous bilayer without obvious defects. A region of interest (ROI) of the bilayer was bleached and the fluorescence recovery was recorded. All membranes were mobile as the fluorescence intensity increased over time. The diffusion coefficients were determined using a MatLab Scipt from Jönssen $et\ al.^{78}$ In Table 4.2 the determined diffusion coefficients and the mobile fractions are listed.

Table 4.2: Mean values of the diffusion coefficients D and the mobile fractions F_m of TxR in lipid bilayers with corresponding standard deviations yielded by FRAP experiments. n is the amount of performed FRAP experiments.

Lipid composition	Diffusion coefficient	Mobile fraction	n
POPC/PIP ₂ /TxR	D / μm²/s	F _m / %	
98.5:1:0.5	2.2 ± 0.3	99 ± 1	8
97.5:2:0.5	2.3 ± 0.2	98 ± 1	7
94.5:5:0.5	2.3 ± 0.3	93 ± 4	7
92.5:7:0.5	1.8 ± 0.3	89 ± 4	8
91.5:8:0.5	2.5 ± 0.4	90 ± 2	7
89.5:10:0.5	2.1 ± 0.6	90 ± 3	7

The diffusion coefficient of the bilayers ranged from $1.8-2.5~\mu m^2/s$. Within the standard deviation the values were similar and did not show a trend with increasing receptor lipid concentration. Considering the mobile fractions of the lipid bilayers the PIP₂ concentration seemed to have a minor influence. The mobile fraction generally slightly decreased from (99 ± 1) % for 1 mol% PIP₂ to (90 ± 3) % for 10 mol% PIP₂. The values were significantly different (Mann-Whitney U test, p<0.01). Together, the results indicate successful formation of homogenous and mobile lipid bilayers.

4.1.2 Formation of solid-supported lipid monolayers on silicon dioxide

Supported lipid bilayers are a versatile tool to investigate membrane-protein interaction as they exhibit high mechanical stability. However, the substrate surface can influence the lipid distribution within the both leaflets^{99,100}. This can have an effect on the accessibility of specific receptor lipids on the upper leaflet of the SLBs, necessary for protein binding. To investigate if the PIP₂ is asymmetrically distributed on supported bilayers and the ENTH binding is influenced, RIfS experiments with monolayers were done. In monolayers the receptor lipid should be distributed homogenously and completely accessible for ENTH binding due to the missing second leaflet.

For precise binding studies of ENTH the formation of a stable monolayer is required. SUVs composed of POPC and PIP₂ (0-10 mol%) were spread on hydrophobic functionalized silicon dioxide substrates (cf. chapter 3.2.2), leading to an increase of the ΔOT . Figure 4.4 shows exemplary RIfS time-traces of the spreading process of POPC (100 mol%) and POPC/PIP₂ (90:10) vesicles on these substrates.

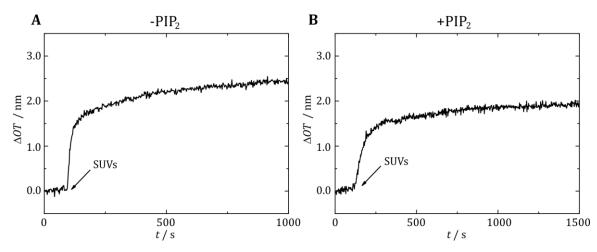


Figure 4.4: Exemplary RIfS time-traces of the vesicle spreading on hydrophobic functionalized silicon dioxide. Pure POPC (A) and POPC/PIP₂ (90:10, B) SUVs were added after a baseline with citrate buffer was observed, resulting in the formation of a lipid monolayer indicated by the ΔOT increase of about 2-3 nm.

After the formation of a baseline with the citrate buffer (cf. chapter 4.1.1) in either case the addition of SUVs led to the ΔOT increase up to 2.1-2.6 nm indicating that the bilayer was formed successfully. First, ΔOT increased fast, then slighter reaching a plateau. The lipid monolayer formation was investigated for different PIP₂ concentrations c (PIP₂ = 0-10 mol%). In Table 4.3 the determined values for ΔOT and $d_{\rm M}$ were listed.

Table 4.3: The optical and physical thicknesses of the POPC/PIP $_2$ lipid monolayers with receptor lipid concentration ranging from 0-10 mol%. The errors were given as standard deviations. n is the number of measurements.

Lipid composition	Δ <i>OT</i> / nm	dм / nm	n
POPC	2.7	1.6 ± 0.2	3
POPC/PIP ₂ (99:1)	2.6 ± 0.3	1.8 ± 0.2	4
POPC/PIP ₂ (98:2)	2.6 ± 0.6	1.7 ± 0.4	5
POPC/PIP ₂ (95:5)	2.1 ± 0.5	1.4 ± 0.3	4
POPC/PIP ₂ (93:7)	2.4 ± 0.7	1.6 ± 0.4	3
POPC/PIP ₂ (92:8)	2.4 ± 0.3	1.7 ± 0.2	3
POPC/PIP ₂ (90:10)	2.3 ± 0.2	1.6 ± 0.2	6

In all cases the monolayers were formed successfully with physical thicknesses of 1.4 to 1.8 nm. The experiments showed no significant influence of the PIP $_2$ concentration on the membrane thickness within the standard deviations.

Since interactions of membrane layers with the substrate surface can occur, which influence the membrane mobility,⁹⁹ diffusion coefficients of TxR in the lipid monolayers were determined by FRAP. Fluorescence images of POPC/PIP₂/TxR monolayers with different PIP₂ concentrations (1 – 10 mol%) revealed an even distribution of the fluorophore, indicating that homogenous monolayers were formed. Rarely occurring small bright spots were assigned to adhered vesicles. In Figure 4.5 exemplary fluorescence images of a FRAP experiment on POPC/PIP₂/TxR (89.5/10/0.5) and the corresponding time-trace of the normalized fluorescence intensity are illustrated.

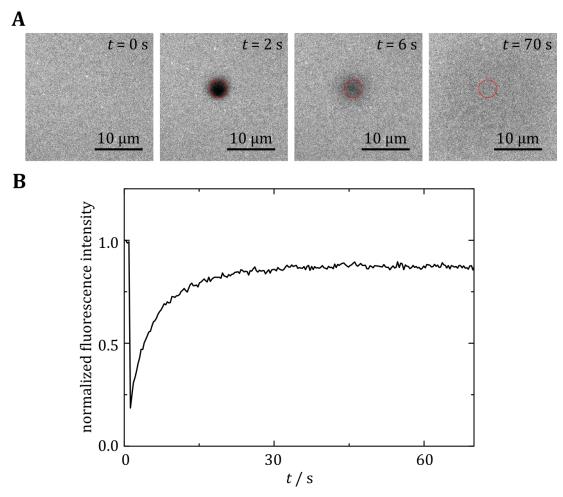


Figure 4.5: (A) Fluorescence images of a FRAP experiment and (B) the corresponding time-trace of the normalized fluorescence intensity on a POPC/PIP $_2$ /TxR (89.5:10:0.5) monolayer. After bleaching of a ROI (red circle, t = 2 s) the fluorescence intensity increases, reaching a plateau after 40 s. Due to the immobile fraction of TxR in the bilayer, the fluorescence intensity does not reach the fluorescence intensity value at t = 0.

Bleaching of a ROI led to the decrease of the fluorescence intensity (Figure 4.5 B). By evaluating the fluorescence recovery with the MatLab Scrip of Jönsson *et al.*⁷⁸ the diffusion coefficients and the mobile fractions ($F_{\rm m}$) of the fluorophore were determined (Table 4.4).

Table 4.4: Mean values of the diffusion coefficients D and the mobile fractions F_m of the fluorophores in lipid monolayers with corresponding standard deviations yielded by FRAP experiments. n is the amount of performed FRAP experiments.

Lipid composition	Diffusion coefficient	Mobile fraction	n
POPC/PIP ₂ /TxR	D / μm2/s	F _m / %	
98.5:1:0.5	0.36 ± 0.02	98 ± 1	5
97.5:2:0.5	0.37 ± 0.03	98 ± 1	6
94.5:5:0.5	0.36 ± 0.03	95 ± 2	5
92.5:7:0.5	0.41 ± 0.02	98.0 ± 0.1	6
91.5:8:0.5	0.42 ± 0.02	98 ± 3	6
89.5:10:0.5	0.41 ± 0.02	98 ± 8	6

The diffusion coefficient of TxR in all monolayers was calculated to be $D=0.4+0.1~\mu m^2/s$ (n=34). Moreover, mobile fractions of $F_m=98\pm2\%$ (n=34) were determined. For both values no influence of the receptor lipid content was observed. However these results indicate that the hydrophobic substrate surface affects the lateral diffusion of artificial membranes as the bilayer showed higher diffusion coefficient values (cf. section 4.1.1).

4.1.3 Analysis of the ENTH binding on lipid mono- and bilayers

It is known that the ENTH domain binds specifically to PIP_2 , 8,23 but the effect of the receptor lipid concentration on this binding is poorly investigated yet. Furthermore also the substrate seems to have a significant influence on the lipid distribution 99 and mobility and consequently on the protein adsorption to membranes. Therefore, the binding of ENTH to $POPC/PIP_2$ bi- and monolayers with varying PIP_2 content was analyzed via RIfS experiments. After the successful bi- and monolayer formation (cf. chapter 4.1.1 and 4.1.2), 1 μ M ENTH was added. In Figure 4.6 a schematic illustration of a whole RIfS experiment is shown.

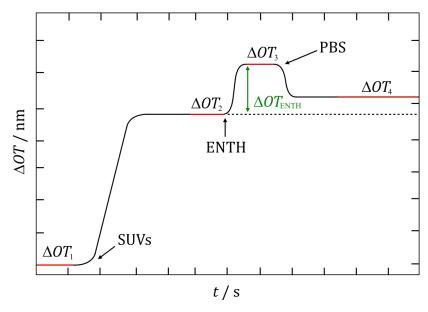


Figure 4.6: Schematic illustration of a RIfS experiment. After formation of a membrane (ΔOT_2) by SUV spreading, protein was added, resulting in a second increase of ΔOT . By rinsing with PBS ENTH desorbed from the surface, reaching a new plateau (ΔOT_4). The optical thickness of the protein (ΔOT_{ENTH}) was determined with the difference of ΔOT_3 and ΔOT_2 .

Addition of ENTH resulted in the increase of optical thickness. The difference of the ΔOT after protein adsorption (ΔOT_3) and ΔOT_2 yielded the ΔOT of the protein (ΔOT_{ENTH} , equation (4.3)), which can be translated into the physical thickness of the protein at complete surface coverage (d_{ENTH} , equation (3.7), with a refractive index $n_{\text{proteins}} = 1.47^{72}$).

$$\Delta OT_{\rm ENTH} = \Delta OT_3 - \Delta OT_2 \tag{4.2}$$

Dissociation of the ENTH domain from the membrane, bi- or monolayer, was reached by rinsing the system with PBS. The reversibility of the ENTH binding to PIP₂ was calculated using equation (4.3).

Reversibility / % =
$$\frac{\Delta OT_3 - \Delta OT_4}{\Delta OT_3 - \Delta OT_2}$$
 (4.3)

In order to show the specificity of the ENTH-PIP₂ binding, the experiments were first performed with pure POPC bi- and monolayers. In Figure 4.7 the time-resolved change in optical thickness during the ENTH addition is illustrated. To highlight the

changes upon ENTH incubation the time and ΔOT after bi- and monolayer formation (ΔOT_2) in the figures were set to zero henceforth.

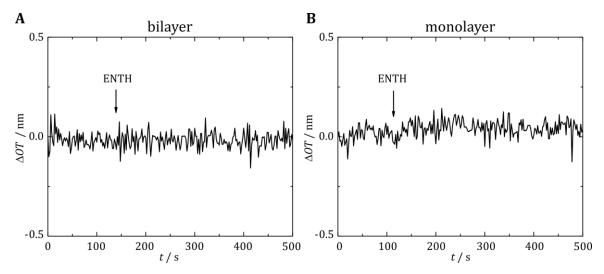


Figure 4.7: RIfS time-trace of the ENTH incubation (1 μ M) to pure POPC (A) bi- and (B) monolayers. After the protein addition no increase was observed, indicating that no ENTH was bound.

In both cases even after some time no increase of the optical thickness was observed (Figure 4.7), showing the specific binding of the ENTH domain to PIP_2 as no protein was bound in the absence of the receptor lipid. Afterwards the ENTH binding to $POPC/PIP_2$ mono- and bilayers dependent on the PIP_2 concentration was investigated. The content of the receptor lipid therefore was varied from one to ten mol%. The exemplary changes in ΔOT upon ENTH binding to $POPC/PIP_2$ (92:8) mono- and bilayers are illustrated in Figure 4.8.

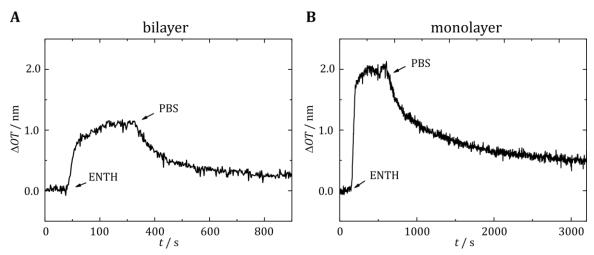


Figure 4.8: Representative time-resolved change in optical thickness observed by means of RIfS. showing the specific binding of 1 μ M ENTH to a (A) bi- and (B) monolayer composed of POPC/PIP₂ (92:8).

In either case ΔOT first increased fast and then reached a plateau when ENTH was added. For the bilayer the ΔOT in the example increased up to 1.1 nm. Compared to that protein adsorption to the monolayer resulted in a change in ΔOT up to 2.0 nm, being about twice as high under the same conditions. This on the one hand shows, that PIP₂ is necessary for ENTH binding to lipid membranes (Figure 4.8) and on the other hand proves the influence of the used membrane system (bilayer vs. monolayer) on binding. Afterwards rinsing with PBS led to the decrease of ΔOT in both cases, indicating reversible binding (Figure 4.8, Table 10.2).

Subsequently further binding studies were done for the other lipid compositions (0-10 mol%). Based on these experiments $\Delta OT_{\rm ENTH}$ was calculated using equation (4.3). In Figure 4.9, the time-resolved change in ΔOT upon protein addition and the obtained averaged $\Delta OT_{\rm ENTH}$ values on bi- and monolayers are plotted for the different PIP₂ concentrations.

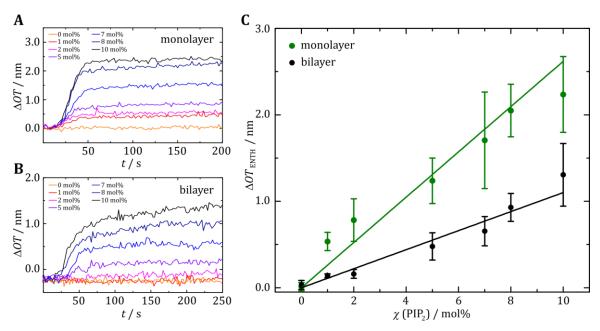


Figure 4.9: Time-resolved increase of ΔOT upon binding of 1 μ M ENTH to POPC (A) monolayers and (B) bilayers doped with PIP₂ (0-10 mol%). (C) Determined maximum changes in optical thickness after ENTH addition obtained from (A) and (B) plotted against the PIP₂ content. The error bars in (C) are given as standard deviations. A linear fit weighted by the standard deviation was plotted trough the data. The mean values with standard deviation and the number of measurements for each lipid composition are listed in Table 10.1 (cf. Appendix).

In both cases, the mono- and bilayer, the $\Delta OT_{\rm ENTH}$ increased linearly with increased receptor lipid concentration, showing the significant influence of the PIP₂ content. Under same conditions the amount of bound protein to POPC/PIP₂ monolayers was about twice as high as for the bilayers. $\Delta OT_{\rm ENTH}$ for the monolayers with 10 mol% PIP₂ reached a value of 2.2 ± 0.4 nm, indicating the beginning of a saturation as the values for 8 mol% (2.0 ± 0.3 nm) and 10 mol% PIP₂ are quite similar within the error. For the bilayers $\Delta OT_{\rm ENTH}$ at 10 mol% was determined to be 1.2 ± 0.4 nm. A saturation cannot be excluded within the errors.

Using equation (4.3) the reversibility of the ENTH binding to both membrane systems were calculated for the different PIP₂ concentrations (Figure 4.10).

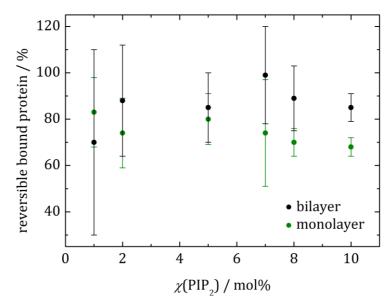


Figure 4.10: Reversibility of the ENTH binding to POPC/PIP₂ bi- and monolayers as a function of the PIP₂ concentration, indicating that a high amount of the protein binds reversible to PIP₂ doped membranes. The mean values with standard deviation and the number of measurements for each lipid composition are listed in Table 10.2 (cf. Appendix)

70-99% of the ENTH bound reversibly to the bilayer. No trend dependent on the PIP₂ concentration was observed. Having the noise to signal ratio in mind, qualified statements of the ΔOT at low PIP₂ concentrations are difficult to make as the ΔOT ENTH was only 0.1 ±0.1 nm at a receptor lipid concentration of 1 mol%. This systematical error also explains the high deviations when working with low PIP₂ concentrations.

Contrary to that the reversibility of the ENTH binding to monolayers showed a slight tendency. With increasing PIP₂ content the reversibility slightly decreased from $83 \pm 15 \%$ (1 mol%) to $68 \pm 4 \%$ (10 mol%), but within the error margins the differences were negligible (Figure 4.10). Using a Mann-Whitney U test revealed that the values are not significantly different (p=0.3).

In general the higher amount of bound protein to the monolayers showed that more PIP₂ was accessible in the monolayers compared to the bilayers under the same conditions.

4.2 The influence of PS on the ENTH binding to PIP₂ containing solid supported lipid bilayers

Not only the PIP₂ content in the membrane has an influence on the amount of bound ENTH, also other lipids showed to cause an increased protein binding. Hom *et al.* performed liposome binding assays, which evidenced the higher affinity of ENTH to SUVs composed of PC/PE/PS/C₁₆-PIP₂ compared to SUVs lacking PS.⁸ Moreover, B. Kroppen presented a co-sedimentation assay to show the appearance of protein homo-oligomers in liposomes in the presence of PS.²⁶

In this chapter the impact of PS on the binding affinity of ENTH to PIP₂ doped lipid bilayers was analyzed via RIfS. Furthermore AFM experiments were performed to investigate the surface topology of bilayers after ENTH incubation.

4.2.1 Influence of PS on the bilayer formation

For exact binding studies of ENTH to PS-containing lipid bilayers by means of RIfS the formation of a homogenous and mobile membrane is requisite. Since the lipid composition was changed by adding the negatively charged lipid PS, the quality of the formed bilayer and the thickness can be influenced. Therefore, SUVs composed of POPC/POPS (80:20) and POPC/POPS/PIP₂ (75:20:5) were spread on hydrophilic silicon dioxide (see Figure 4.11), resulting in the ΔOT increase. In all cases SUV spreading was performed in citrate buffer.

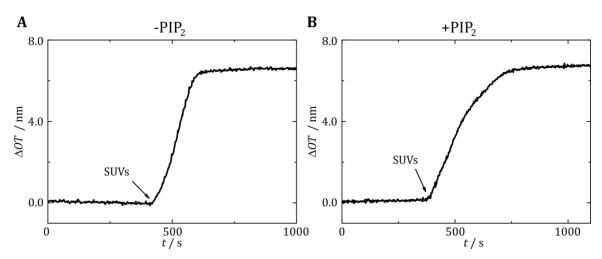


Figure 4.11: Exemplary time-trace of (A) POPC/POPS (80:20) and (B) POPC/POPS/PIP₂ (75:20:5) vesicles spreading on hydrophilic silicon dioxide substrates.

The graphs in Figure 4.11 show that the formation of bilayers with both lipid composition were successful. The influence of the lipid head groups on the membrane height was then investigated using equation (4.3). Table 4.5 shows the calculated values of $\Delta OT_{\rm M}$ and d of POPC/POPS (80:20) and POPC/POPS/PIP₂ bilayers.

Table 4.5: Mean values of the optical and physical thicknesses of the POPC/POPS (80:20) and POPC/POPS/PIP $_2$ (75:20:5) lipid bilayers. The errors were given as standard deviations. n describes the amount of performed experiments.

Lipid composition	$\Delta OT_{\rm M}$ / nm	<i>d</i> / nm	n
POPC/POPS (80:20)	6.4 ± 0.3	4.3 ± 0.2	6
POPC/POPS/PIP ₂ (75:20:5)	6.3 ± 0.2	4.2 ± 0.1	13

Like shown in Table 4.5, increasing of the POPS content did not change the membrane thickness (cf. Table 4.1). There was also no difference in the POPC/POPS bilayer thickness containing or lacking PIP₂. These results show that the negative charges of POPS and PIP₂ changed the curve shape during the bilayer formation but did not influence the membrane thickness. The lateral diffusion of TxR in the POPC/POPS and POPC/POPS/PIP₂ bilayers was measured to be $D = 1.9 \pm 0.3 \, \mu m^2/s$ (without PIP₂) and $D = 1.4 \pm 0.4 \, \mu m^2/s$ (with PIP₂). Although the diffusion coefficient of TxR in the POPC/POPS/PIP₂ bilayers slightly decreased compared to POPC/POPS and POPC/PIP₂ bilayers (1.8-2.5 $\, \mu m^2/s$), the difference got negligible within the deviations.

Furthermore mobile fractions of $F_m = 86 \pm 8\%$ (POPC/PIP₂) and $F_m = 76 \pm 4\%$ (POPC/POPS/PIP₂) were determined. These results show that the formation of mobile bilayers composed of POPC/POPS and POPC/POPS/PIP₂ was successful. Moreover; no obvious defects were detected by fluorescence microscopy (Figure 4.12).

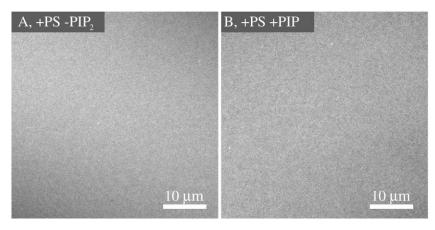


Figure 4.12: Fluorescence images of (A) POPC/POPS/TxR (79.5:20:0.5) and (B) POPC/POPS/PIP₂/TxR (74.9:20:5:0.5) bilayers on silicon dioxide, showing a homogenous distribution of the fluorophore.

PIP₂ is the receptor lipid for ENTH, has a negative charge of -3 to -5 under physiological conditions⁴⁴ and is binding in an ENTH binding pocket surrounded by positively charged amino acids. Moreover, previous studies on the ENTH-PIP₂ interaction on various membrane models obtained higher binding affinity of ENTH to PIP₂ when lipids with the head group PS (negatively charged) are involved.^{8,26} Therefore the next chapters focus on the influence of PS on the ENTH-PIP₂ interaction.

4.2.2 Influence of PS on the binding affinity of ENTH to PIP2

The influence of PS on the binding affinity of ENTH was examined with RIfS measurements. By addition of small unilamellar vesicles (SUVs) to hydrophilic silicon dioxide wafers, the SUVs spread on the surface. Thereby the optical thickness increased, indicating the successful formation of a bilayer (cf. Section 3.3.2). To investigate if ENTH also binds to the negatively charged lipid PS, ENTH was added to a POPC/POPS (80:20) bilayer lacking PIP₂ (Figure 4.13).

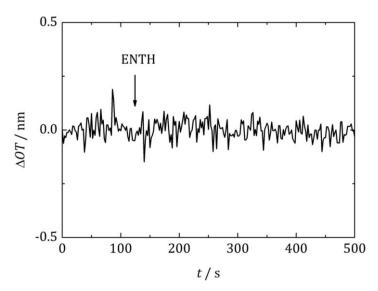


Figure 4.13: Exemplary time-trace of a RIfS blind experiment. Addition of ENTH (1 μ M) to a preformed POPC/POPS (80:20) bilayer did not show changes in the optical thickness, which proves that the ENTH binds specifically to PIP₂ and does not interact with the negative charge of PS in the absence of PIP₂.

No significant increase of ΔOT was observed when the ENTH domain was added to the POPC/POPS bilayer, showing no binding of ENTH to PS in the absence of PIP₂. This proves the specific protein binding to its receptor lipid. Based on this, binding studies of ENTH to POPC/PIP₂ (95:5) and POPC/POPS/PIP₂ (75:20:5) bilayers were performed with RIfS experiments to investigate the influence of PS on the binding affinity of ENTH when PIP₂ is present. In Figure 4.14 representative time traces upon addition of different ENTH concentrations to these bilayers are illustrated.

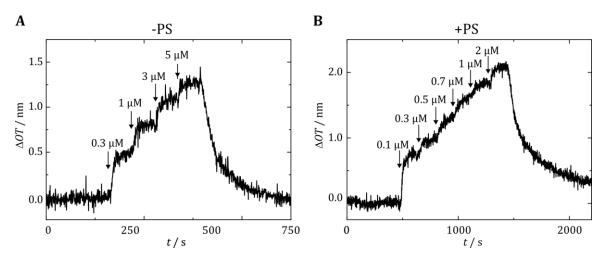


Figure 4.14: Representative RIfS time-traces of the ENTH adsorption upon different protein concentrations to POPC/PIP₂ bilayers in (A) absence and (B) presence of PS. The protein additions are marked by arrows. After rinsing with PBS the protein desorbed from the surface showing the reversible ENTH-PIP₂ binding. Lipid compositions: POPC/PIP₂ (A, 95:5), POPC/POPS/PIP₂ (B, 75:20:5).

For both lipid compositions the optical thickness increased with the addition of the protein domain. Higher ENTH concentrations led to increased ΔOT values. The difference between the experiments were the maximum ΔOT values, which reached values of $\Delta OT = 1.3$ nm at 5.0 μ M ENTH in absence and $\Delta OT = 2.1$ nm at 2.0 μ M ENTH in presence of PS, indicating a higher binding affinity of the protein to bilayers containing PIP2 and PS. With POPC/PIP2 bilayers five and with POPC/POPS/PIP2 bilayers eight individual experiments were performed. The isotherms revealed a high reversible binding of ENTH to POPC/PIP2 (95:5). $96 \pm 10\%$ of the ENTH desorbed from the membrane surface during rinsing with PBS. Compared to that $89 \pm 3\%$ of the ENTH bound reversible to POPC/POPS/PIP2 (75:20:5) bilayers. A Mann-Whitney U test revealed that the differences are not significant (p=0.2).

Based on this data the Langmuir adsorption model was used to calculate the dissociation constant (K_D) of the ENTH binding to PIP₂-containing membranes as well as the maximum of the optical thickness after protein addition ($\Delta OT_{ENTH, max}$). This model is in a good accordance with the used system, when the surface is homogenous, all binding sites are equal as well as a monolayer is formed and when no interactions of the adjacent molecules occur. ^{101,102} Furthermore a dynamic equilibrium of the bound and free molecules is assumed. As the vesicle spreading was successful (cf. 4.1.1 and 4.2.1), indicating a homogenous membrane surface and the ENTH binding is almost

completely reversible the used system seems to match these assumptions. In Figure 4.15 the adsorption isotherms of ENTH to POPC/PIP₂ (95:5) and POPC/POPS/PIP₂ (75:20:5) bilayers are illustrated, showing the influence of the lipid composition on the binding affinity.

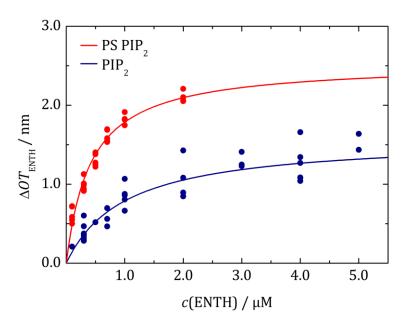


Figure 4.15: Adsorption isotherms of ENTH to POPC/PIP₂ (95:5) (blue circles) and POPC/POPS/PIP₂ (75:20:5) (red circles) bilayers. The values for $\Delta OT_{\text{ENTH, max}}$ and K_D were obtained by fitting a Langmuir adsorption isotherm (solid lines) to the data. Non-linear regression weighted by the corresponding number of measurements that went into each concentration was carried out using a Levenberg-Marquardt algorithm. 103,104

Figure 4.15 shows ΔOT_{ENTH} values after the addition of different ENTH concentrations to POPC/PIP₂ (blue circles) and POPC/POPS/PIP₂ (red circles) bilayers.

The data was fitted with a Langmuir adsorption isotherm (solid lines). Due to different numbers of data points, a non-linear regression with a Levenberg-Marquardt algorithm weighted by the number of realized measurements for each concentration was carried out. 103,104 This enabled the determination of the maximal change in optical thickness after protein addition $\Delta OT_{\rm ENTH,\, max}$ and the dissociation constant $K_{\rm D}$ by equation (4.4).

$$\Delta OT = \Delta OT_{\text{ENTH, max}} \frac{c}{K_{\text{D}} + c}$$
 (4.4)

When no PS was present, mean values of $\Delta OT_{\rm ENTH,\,max} = 1.6 \pm 0.1$ nm and $K_{\rm D} = 1.0 \pm 0.2$ µM were obtained. For the POPC/POPS/PIP₂ system the mean values of $\Delta OT_{\rm ENTH,\,max} = 2.5 \pm 0.1$ nm and $K_{\rm D} = 0.42 \pm 0.05$ µM were calculated, indicating the higher binding affinity of ENTH to bilayers containing POPS in the presence of PIP₂.

4.2.3 Topographical analysis of ENTH cluster formation

In the previous chapters the specific binding of ENTH to PIP₂ containing bilayers was shown. The topography of the substrate surface before and after protein binding was investigated by means of atomic force microscopy (AFM). This technique enables to obtain information about the protein organization on supported membrane models. 105,106 The RIfS measurements showed an increased binding affinity and a higher occupancy of ENTH when lipids with the head group PS were present in the bilayer (cf. chapter 4.2.2). Therefore the surface topographies of PIP2 doped supported bilayers with and without PS were analyzed. For better comparison with previous results,^{26,79} AFM experiments were performed with bilayers composed of DOPC/DOPE/PIP₂/TxR (64.9:30:5:0.1) and DOPC/DOPE/DOPS/PIP2/TxR (44.9:30:20:5:0.1). Fluorescently labeled supported lipid bilayers were generated as described in chapter 3.2.2. Afterwards force-distance curves were recorded via AFM, which enabled to determine the breakthrough force and the membrane thicknesses (cf. chapter 3.3.4). The breakthrough force is the maximum vertical force at which the bilayer is able to stand before the cantilever tip breaks through the membrane. 107 In Figure 4.16 an exemplary force distance curve with a breakthrough event is illustrated.

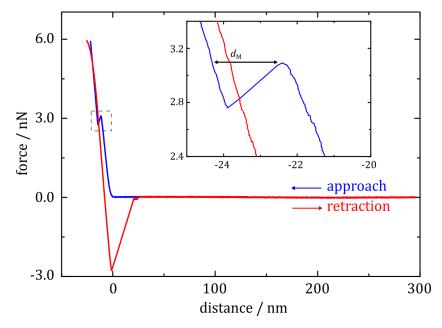


Figure 4.16: A representative force-distance curve measured on DOPC/DOPE/DOPS/PIP₂/TxR (44.9:30:20:5:0.1). The inlet shows the breakthrough event (dashed square) which enabled to determine the breakthrough force and the membrane thickness $d_{\rm M}$

Breakthrough forces were measured to be between 2-6 nN. Furthermore the AFM experiments revealed lipid bilayers with physical thicknesses of $d=3.9\pm0.8$ nm (without PS) and $d=3.7\pm07$ nm (with PS). Also the FRAP experiments proved the formation of mobile bilayers. After bleaching a region of interest (ROI) the fluorescence recovered in the ROI. Diffusion coefficients of $2.5\pm0.6~\mu m^2/s$ were measured for TxR in the DOPC/DOPE/PIP2/TxR (64.9:30:5:0.1) and $1.7\pm0.4~\mu m^2/s$ in the DOPC/DOPE/DOPS/PIP2/TxR (44.9:30:20:5:0.1) bilayers on hydrophilic silicon dioxide using a MATLAB Script by Jönsson *et al.*⁷⁸ The mobile fraction of TxR in the bilayers without PS were determined to be $83\pm3\%$ and the value for bilayers with PS resulted in a mobile fraction of $71\pm5\%$. Hence the AFM and FRAP experiments proved the formation and mobility of the bilayers.

When the bilayer formation was successful the surface was imaged by AFM using contact mode (cf. chapter 3.3.4). This enabled to compare the surface topography before and after protein addition. Then ENTH (1 μ M) was added and the solution was mixed by stirring for 2 h. Subsequently again images of the surface were taken. The atomic force micrographs before and after ENTH incubation are shown in Figure 4.17.

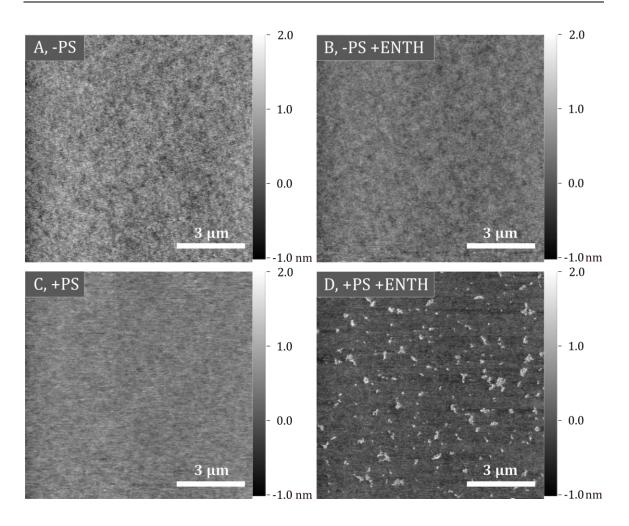


Figure 4.17: Atomic force micrographs of (A) DOPC/DOPE/PIP $_2$ /TxR (64.9:30:5:0.1) and (C) DOPC/DOPE/DOPS/PIP $_2$ /TxR (44.9:30:20:0.1) bilayers on hydrophilic silicon dioxide wafers prior ENTH addition. (B) and (D) show the corresponding micrographs after 2 h of ENTH incubation (1 μ M). Only in the presence of PS, protein clusters were observed on the membrane surface.

The micrographs show that before ENTH addition the DOPC/DOPE/PIP₂/TxR (Figure 4.17 A) and DOPC/DOPE/DOPS/PIP₂/TxR (Figure 4.17 C) bilayers had a homogenous surface with roughnesses of rms = 174 ± 17 pm (n = 7) without PS and rms = 158 ± 14 pm (n = 5) with PS.

Addition of ENTH to the DOPC/DOPE/PIP2/TxR bilayers (Figure 4.17 B) led to a roughness of rms = 196 ± 22 pm (n = 21), but no structures were observed (Figure 4.17 C). When PS was present in the bilayers incubation with ENTH for 2 h caused the appearance of small structures, which were about 0.5 μ m larger or smaller. To investigate if these structures could be assigned to protein adsorption the heights of these clusters were analyzed using a MatLab Script written by Dr. Ingo Mey (Georg-August-

University Göttingen). A threshold of the membrane height enabled to distinguish between membrane and adsorbed molecules, so that the protein occupancy and height could be determined (cf. chapter 3.3.4). From the AFM images the protein occupancy of $6 \pm 1\%$ and the protein height of 1.2 ± 0.5 nm (mean value \pm SD, n = 32, with n the number of evaluated images from three independent experiments) were calculated (Figure 4.18).

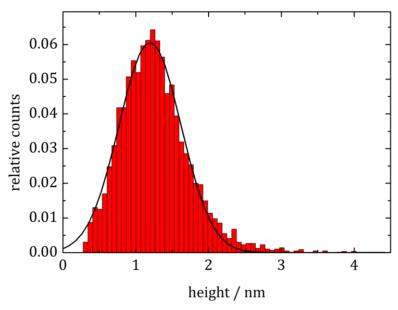


Figure 4.18: Histogram plot of the protein heights averaged over 32 micrographs. A protein height of 1.2 ± 0.5 nm was determined by fitting a normal distribution. The bin size (0.066 nm) was calculated according to Freedman *et al.*¹⁰⁸

With these results it seemed likely that the small structures were caused by protein adsorption and were attributed to ENTH clusters. Due to the ENTH cluster formation the surface roughness increased up to rms = 263 ± 33 pm (n = 32).

To ensure that the protein really bound to the membrane and not only to defects, occurring in the membrane, fluorescence images of these bilayers before and after ENTH incubation were taken. As control also the surface of membranes without ENTH incubation were imaged after 2h, showing a homogenous fluorescence over the whole substrates without visible defects in all cases (Figure 4.19). The few small bright spots were attributed to SUVs adhered to the surface.

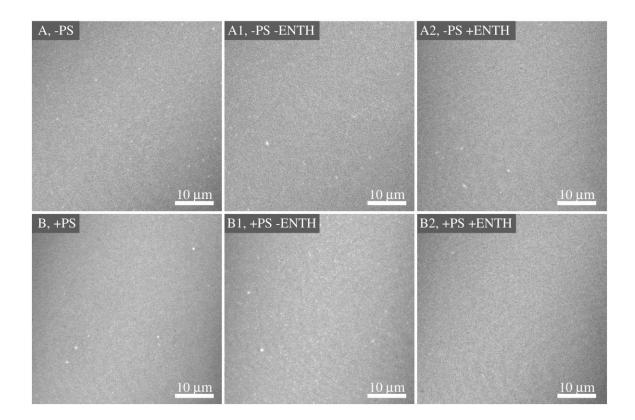


Figure 4.19: Fluorescence images of (A) DOPC/DOPE/PIP $_2$ /TxR (64.9:30:5:0.1) and (B) DOPC/DOPE/DOPS/PIP $_2$ /TxR (44.9:30:20:5:0.1) bilayers. After 2h of incubation without (A1, B1) and with 1 μ M ENTH (A2, B2) the bilayers were imaged again. In all cases a homogenous distribution of TxR was observed.

AFM allows imaging in a nanometer range resolution¹⁰⁹ so even single protein adsorption should be possible to observe. The discrepancy of the protein coverage in RIfS (cf. chapter 4.2.2) and AFM is not obvious. Hence FRAP measurements were performed to analyze the mobility of the bilayers after the ENTH incubation. The diffusion coefficients of DOPC/DOPE/PIP₂/TxR (64.9:30:5:0.1) and DOPC/DOPE/DOPS/PIP₂/TxR (44.9:30:20:5:0.1) bilayers after ENTH addition were determined. The bilayers were still mobile, but the mobility was decreased. In Table 4.6 the diffusion coefficients of TxR in the bilayers before and after ENTH addition are listed.

Table 4.6: Diffusion coefficients of TxR in DOPC/DOPE/PIP $_2$ /TxR (64.9:30:5:0.1) and DOPC/DOPE/DOPS/PIP $_2$ /TxR (44.9:30:20:5:0.1) bilayers before and after the incubation of 1 μ M ENTH.

lipid composition	before ENTH addition	after ENTH addition
DOPC/DOPE/PIP ₂ /TxR	$2.5 \pm 0.6 \mu\text{m}^2/\text{s}$	$0.6 \pm 0.2 \mu \text{m/s}^2$
(64.9:30:5:0.1)		
DOPC/DOPE/DOPS/PIP ₂ /TxR	$1.7 \pm 0.4 \mu m^2/s$	$0.3 \pm 0.2 \mu \text{m/s}^2$
(44.9:30:20:5:0.1)		

The results show that the diffusion coefficient of TxR in the PS containing membranes after ENTH incubation was about the half of the diffusion coefficient of TxR in the DOPC/DOPE/PIP₂ bilayers when ENTH was present. Although the protein adsorption seemed to modulate the membrane mobility, imaging of single protein molecules was prevented due to a dynamic protein distribution in contrast to more immobile ENTH clusters (when PS is present).

4.2.4 Investigation of the ENTH R114A mutant binding to PIP2

The ENTH domain consists of 150 amino acids, which are essential for various intraand intermolecular interactions. Mutation experiments were done by Ford *et al.* to
investigate the role of the amphipathic helix in the membrane deformation process. 21 With co-sedimentations assays they presented that the binding affinity and membrane deformation activity change when the sixth amino acid in the helix was mutated. The binding affinity decreased with hydrophilicity of the amino acid residue.
Furthermore, Yoon *et al.* used electron paramagnetic resonance (EPR) and fluorescence microscopy showing that clustering of the ENTH domain appears. 28 With the
mutants K23E and E42K only binding, but no clustering of ENTH was observed, showing the importance of these amino acids for protein-protein interactions. Lai *et al.* presented atomistic MD simulations, where they identified three important ENTHmembrane interaction regions: The binding pocket for PIP2, the α 0-helix and the
R114 loop (cf. Figure 1.3). 110

In order to specify the regions of the ENTH domain involved in certain interactions, experiments with the ENTH R114A mutant were performed. In this mutant the arginine in position 114 in the loop between the sixth and seventh helix is exchanged by an alanine. The mutant was expressed and purified by Dr. Ben Kroppen and Indrani Mukherjee from the group of Prof. Dr. Michael Meinecke (Department of Cellular Biochemistry, University Medical Center Göttingen). First of all, RIfS measurements were done to investigate the binding behavior of the mutant depending on the lipid ofcomposition. Therefore vesicles composed POPC/PIP₂ (95:5)POPC/POPS/PIP₂ (75:20:5) were added to hydrophilic silicon dioxide substrates (cf. chapter 3.3.1), resulting in vesicle spreading and the formation of a bilayer. After rinsing with PBS and passivation with BSA, the ENTH R114A mutant was added. The ENTH concentration was increased stepwise (Figure 4.20, indicated by arrows). A representative time-traces of the optical thickness (ΔOT) after ENTH R114A addition to POPC/PIP2 and POPC/POPS/PIP2 bilayers are plotted in Figure 4.20.

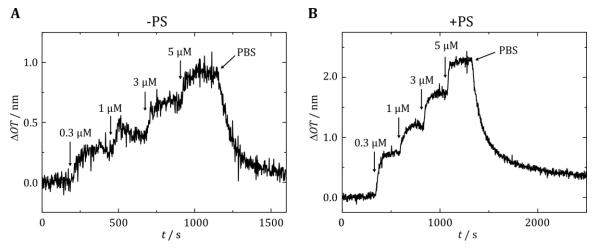


Figure 4.20: Time-resolved change in optical thickness upon addition of different ENTH R114A concentrations (marked by arrows) to (A) POPC/PIP₂ (95:5) and (B) POPC/POPS/PIP₂ (75:20:5) bilayers. The protein desorbs after rinsing with PBS showing the almost complete reversibility of the ENTH R114A binding.

In both cases the optical thickness increased by ENTH R114A addition until it reached a plateau. With every further and higher concentrated protein addition the optical thickness was stepwise increased due to higher protein occupancy on the surface. Rinsing with PBS led to desorption, which indicated the reversible binding. When the

bilayer was lacking POPS the ΔOT reached a value of 0.9 nm at an ENTH R114A concentration of 5.0 μ M (Figure 4.20 A). In case of bilayers with 20 mol% PS the ΔOT increased up to 2.3 nm when 5.0 μ M ENTH R114A was added (Figure 4.20 B). This indicated a higher amount of bound mutant when PS is present. To obtain the maximum change in optical thickness and the binding affinity of the mutant to POPC/PIP₂ and POPC/POPS/PIP₂ bilayers, $\Delta OT_{\rm ENTH~R114A}$ was plotted against the protein concentration (Figure 4.21).

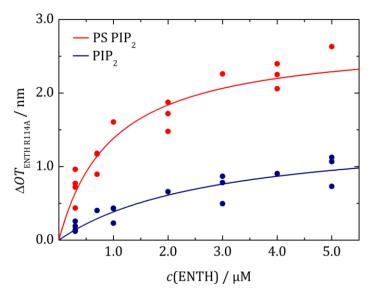


Figure 4.21: Adsorption isotherms of ENTH R114A to POPC/PIP₂ (95:5) (blue circles) and POPC/POPS/PIP₂ (75:20:5) (red circles) bilayers. The values for $\Delta\Delta OT_{max}$ and K_D (Tab. 1) were obtained by fitting a Langmuir adsorption isotherm (solid lines) to the data. Non-linear regression weighted by the corresponding number of measurements that went into each concentration was carried out using a Levenberg-Marquardt algorithm. 103,104

Figure 4.21 shows the changes in optical thicknesses upon addition of different ENTH R114A concentrations to POPC/PIP2 (blue circles) and POPC/POPS/PIP2 (red circles) bilayers. Four individual measurements without POPS and four with POPS at protein concentrations of 0.3 to 5.0 μ M were performed. By fitting the data with a Langmuir adsorption isotherm the maximum of the optical thickness change $\Delta OT_{\rm ENTH~R114A,max}$ and the dissociation constant K_D were obtained, showing the higher binding affinity of the ENTH R114 mutant to PS containing lipid bilayers. A non-linear regression was done using a Levenberg-Marquardt algorithm. Thereby the regression was weighted by the amount of measurements for each concentration. For the POPC/PIP2

(95:5) system a K_D of 2.9 \pm 0.8 μ M and a $\Delta OT_{ENTH~R114A,~max}$ = 1.5 \pm 0.2 nm were determined. Bilayers composed of POPC/POPS/PIP₂ (75:20:5) exhibited a three times smaller K_D value of 1.0 \pm 0.3 μ M and $\Delta OT_{ENTH~R114A,~max}$ of 2.7 \pm 0.2 nm, which is twice as high as for the system lacking POPS. The $\Delta OT_{ENTH~R114A,~max}$ values for the mutant are similar to those of the ENTH wildtype (wt), but the binding affinity of the mutant to both lipid systems decreased compared to that of ENTH wt. All values for the ENTH wt and the mutant are summarized in Table 4.7.

Table 4.7: Summary of the fit results of the Langmuir adsorption isotherms for ENTH wt and ENTH R114A. The K_D and the $\Delta OT_{\text{ENTH, max}}$ are given as parameter \pm SE.

lipid	<i>K</i> _D / μM		$\Delta OT_{\rm ENTH,max}$ / nm	
composition	ENTH wt	ENTH R114A	ENTH wt	ENTH R114A
POPC/ PIP ₂	1.0 ± 0.2	2.9 ± 0.8	1.6 ± 0.1	1.5 ± 0.2
(95:5)				
POPC/POPS/PIP2	0.42 ± 0.05	1.0 ± 0.3	2.5 ± 0.1	2.7 ± 0.2
(75:20:5)				

Table 4.7 shows that the mutant was still binding to the PIP₂-containing membranes, but its binding affinity to PIP₂ doped membranes changed. The mutant also bound reversible to those membranes. For the POPC/PIP₂ (95:5) bilayers a reversibility of $84 \pm 9\%$ was determined. Moreover $89 \pm 3\%$ of the ENTH R114 mutant bound reversibly to POPC/POPS/PIP₂ (75:20:5) bilayers.

Furthermore ENTH wt formed clusters on bilayers composed of DOPC/DOPE/DOPS/PIP2/TxR, which was shown by AFM experiments (cf. Figure 4.17). To compare if also ENTH R114A mutant clusters on PS containing supported bilayers appear, atomic force micrographs of the surface before and after protein addition were taken. In Figure 4.22 the micrographs of DOPC/DOPE/PIP2/TxR (64.9:30:5:0.1) and DOPC/DOPE/DOPS/PIP2/TxR (44.9:30:20:0.1) bilayers before (A, C) and after (B, D) ENTH addition are illustrated.

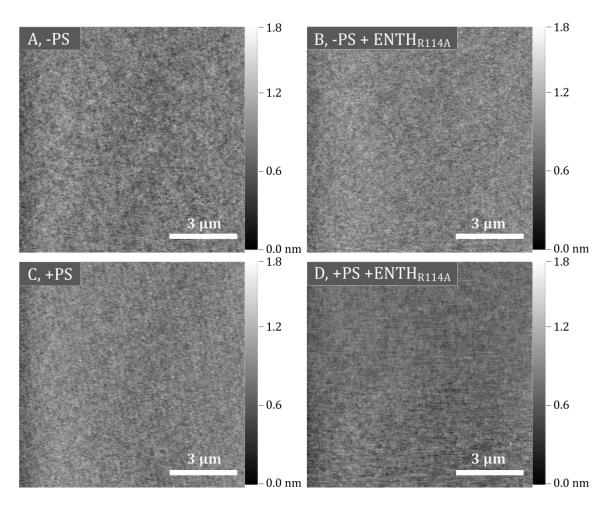


Figure 4.22: Atomic force micrographs of (A) DOPC/DOPE/PIP $_2$ /TxR (64.9/30/5/0.1) and (C) DOPC/DOPE/DOPS/PIP $_2$ /TxR (44.9/30/20/0.1) bilayers on hydrophilic silicon dioxide wafers prior protein addition. The corresponding micrographs (B, D) were obtained after 2 h of ENTH R114A incubation (1 μ M). Even in the presence of DOPS, no protein clusters were observed on the membrane surface indicating the necessity of the amino acid R114 for ENTH cluster formation.

The surfaces of the DOPC/DOPE/PIP2/TxR (Figure 4.22 A) and DOPC/DOPE/DOPS/PIP2/TxR (Figure 4.22 C) bilayers showed a homogenous distribution of the membrane. For the bilayers without PS a surface roughness of rms = 142 ± 13 pm (n = 3) was determined. When PS-containing bilayers were imaged surface roughness of rms = 142 ± 12 pm (n = 4) was obtained. The micrographs also showed that in both cases no specific structures were observed on the surface after the protein addition (Figure 4.22 B and D) in contrast to the ENTH wt (cf. Figure 4.17). The roughness of the surfaces after ENTH incubation were calculated to be slightly higher with rms = 152 ± 7 pm (n = 3) lacking PS and rms = 197 ± 41 pm (n = 11) in presence of PS. This showed that the surface nature changed due to ENTH

binding, although no ENTH R114A clusters were observed, indicating that the amino acid R114A is important for clustering of ENTH.

4.3 Investigation of the ENTH-PIP₂ interactions on lipid monolayers

The concentration of the receptor lipid PIP₂ and the presence of PS have an influence on the amount of bound ENTH (cf. sections 4.1.3 and 4.2.2). Moreover, it was shown that substrate interactions with the membrane lead to asymmetric artificial lipid bilayers on hydrophilic silicon dioxide surfaces. This was shown by binding studies on supported lipid monolayers (cf. chapter 4.1.3), indicating a higher protein occupancy compared to supported bilayers.

Another possibility to investigate the interactions of ENTH with monolayers is the Langmuir-Blodgett trough with the advantage that interactions with substrate surfaces are prevented. This method allowed to investigate the surface activity of ENTH as a function of the lipid composition. By measuring surface pressure changes upon ENTH addition, the penetration ability of the ENTH helix could be determined. The influence of the lipids PIP2 and PS on the ENTH penetration ability was investigated. First of all, surface pressure-area isotherms of all lipid compositions were recorded to analyze how the lipid composition influences the monolayer. Moreover, it was important to analyze the monolayer stability of the desired lipid mixtures. Only if the monolayers were stable enough over the experimental time, meaning lipids did not submerge into the subphase, it was guaranteed that enough PIP2 molecules were available for ENTH binding. Finally, penetration experiments with different lipid mixtures were performed.

4.3.1 PIP2 influence on the ENTH penetration ability

Binding of ENTH to bi- and monolayers is dependent on the concentration of the receptor lipid PIP₂ (cf. chapter 4.1.3). Due to this the intermolecular interactions of the

lipids within these membrane models are of great interest. Depending on the head groups, the number of carbon atoms in the fatty acid chains and their conformation, the isotherms can differ. These isotherms help to understand the lateral intermolecular interactions in the monolayer and the influence on protein-membrane interactions. The Langmuir-Blodgett trough technique was used to record pressurearea isotherms of POPC/POPE monolayers with different PIP2 content. To mimic the natural plasma membrane most closely POPC and POPE were used as these are present in membranes in a high amount. Furthermore, PIP2 was added as receptor lipid for the ENTH domain. Isotherms of POPC/POPE (80:20) and POPC/POPE/PIP2 (80-x:20:x, with x=1, 5, 10 mol%) were measured in ultrapure water. As the following penetration experiments with protein were performed in HEPES buffer, also isotherms were recorded in HEPES buffer. The corresponding averaged surface pressure (π) versus area per molecule $(A_{\rm M})$ diagrams are plotted in Figure 4.23.

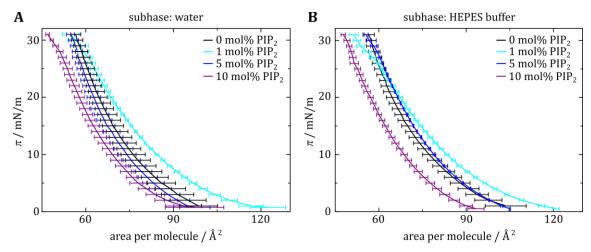


Figure 4.23: Averaged surface pressure-area per molecule isotherms of POPC/POPE (80:20, black), POPC/POPE/PIP2 (79:20:1, light blue), POPC/POPE/PIP2 (75:20:5, blue) and POPC/POPE/PIP2 (70:20:10, purple) at 20 °C. The subphases were ultrapure water (A) and HEPES buffer (B). The error bars represents the standard deviations of the area per molecule values at the same surface pressure with a surface pressure deviation of 0.05 mN/m. The isotherms were recorded from π = 0 to 32 mN/m showing an increase of the lipid packing with increased PIP2 concentration.

Figure 4.23 shows averaged isotherms of POPC/POPE (80:20) and POPC/POPE/PIP $_2$ (80-x:20:x, with x=1, 5, 10 mol%) measured in water (Figure 4.23 A) and HEPES buffer (Figure 4.23 B). In all cases, the surface pressure increased with reduced area

per molecule Am. None of the isotherms showed a coexistence region (cf. Section 3.3.5). The isotherms were shifted to lower $A_{\rm M}$ values when the PIP₂ concentration increased from 1 mol% PIP₂ to 10 mol% PIP₂. Thus, increasing the PIP₂ content led to a higher packing density or to the loss of lipids to the subphase. Within the errors, no significant differences of the isotherms in buffer or water were observed. The isotherms of the lipid composition POPC/POPE (80:20) were similar to those of POPC/POPE/PIP2 (75:20:5). Some monolayers with 1 mol% PIP2 on HEPES buffer collapsed at high surface pressures (about 27-30 mN/m), indicated by the drop of π . In the averaged isotherm of this lipid composition it is represented by a sudden shift to smaller A_M above 25 mN/m (cf. Figure 4.23, light blue, Figure 4.23 B). Thus, the maximum of the lipid packing was reached. To compare the isotherms, the A_{20} -values were listed in Table 4.8. The A_{20} value describes the surface pressure at an area per molecule of $A_{\rm M}$ = 20 Å² and allows to compare the lipid packing of different lipids at a certain surface pressure. At this certain surface pressure a collapse of all monolayer is excluded. The area per molecule values $(A_{\rm M})$ were averaged over all lipids in the ratio of their molecular fraction.

Table 4.8: A_{20} mean values with standard deviations obtained from the isotherms with the different lipid compositions. The subphases were ultrapure water and HEPES buffer (10 mM HEPES, 160 mM KCl, pH 4.7). n describes the amount of recorded isotherms.

	Subphase: ultrapure water		Subphase: HEPES buffer	
Lipid composition	A_{20} /Å ²	n	A_{20} /Å ²	n
POPC/POPE (80:20)	64.6 ± 2.0	4	64.5 ± 2.0	4
POPC/POPE/PIP ₂	68.6 ± 0.8	3	67.3 ± 0.8	3
(79:20:1)				
POPC/POPE/PIP ₂	63.1 ± 2.4	3	66.5 ± 0.1	3
(75:20:5)				
POPC/POPE/PIP ₂	57.8 ± 1.7	4	57.4 ± 0.8	3
(70:20:10)				

The A_{20} values show that the buffer did not influence the lipid packing, but the PIP₂ content did. With an increasing PIP₂ concentration, the area per molecule slightly decreased indicating a denser lipid packing of the monolayers. Within the errors the

POPC/POPE lipid mixture showed similar area per molecules values compared to the lipid compositions with 1 mol% and 5 mol%.

In order to investigate if the monolayers composed of POPC/POPE and POPC/POPE/PIP2 were stable over time, the change in surface pressure after reaching a value above 30 mN/m was recorded. In cell membranes, the lipid packing density has an area per molecule of 40-70 Ų, which is equivalent to a monolayer at this surface pressure. At the beginning, the monolayers were compressed to the desired surface pressure, then the barrier motor was stopped to keep the trough area constant and the surface pressure decay was observed. The monolayer with 1 mol% did not collapse during these experiments. Figure 4.24 shows an exemplary plot of the surface pressure over 1 h for a POPC/POPE/PIP2 (75:20:5) monolayer.

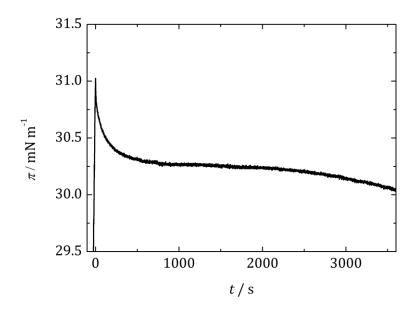


Figure 4.24: Exemplary plot of the surface pressure vs. time of a POPC/POPE/PIP $_2$ (75:20:5) monolayer. After reaching the desired surface pressure (above 30 mN/m) the barriers were stopped at t=0 s, keeping the trough area constant. The change in surface pressure was detected over one hour to determine the change in surface pressure and to obtain information about the monolayer stability. The measurement was performed at 20°C with HEPES buffer as subphase.

After the trough area was held constant, the surface pressure decreased fast about 0.7 mN/m, which was attributed to the barrier motor stop. Even though this behavior was observed often, the surface pressure stabilized after a time and the change in surface pressure got small. The total surface pressure loss was $2 \pm 3\%$ for

1 mol% PIP₂ (n=2), 4 ± 1% for 5 mol% PIP₂ (n=2) and 2 ± 1% for 10 mol% PIP₂ (n=3) over 1 h. The POPC/POPE monolayers without PIP₂ showed a surface pressure loss of 1 ± 3%. If the surface pressure decreased less than 1 mN/m per minute the monolayer was defined as stable.¹¹⁵ Although it could not be excluded that PIP₂ or other lipids submerged into the subphase during the compression, once 30 mN/m was reached, the monolayers with the different lipid compositions were sufficiently stable to perform insertion experiments.

As monolayer stability was guaranteed, penetration experiments of ENTH into POPC/POPE and POPC/POPE/PIP₂ monolayers were done at different initial surface pressures (π_0) with a protein concentration of 1 μ M. $\Delta \pi$ is dependent on the protein concentration and reaches a saturation at $c_{\text{ENTH}} = 0.2 \, \mu \text{M}.^{25}$ Therefore, at higher concentrations the maximal change in the surface pressure ($\Delta \pi$) should not be influenced by the protein concentration and consequently represents the maximum value for $\Delta \pi$. The isotherms were performed on a Teflon trough, comprising a volume of 120 mL. To reduce the amount of necessary protein to yield a concentration of 1 μM, the monolayer penetration experiments were performed with a custom-built protein trough (V = 2 mL, cf. Section 3.3.5). After the desired surface pressure was reached, the small Teflon barriers were inserted into the protein trough. If the surface pressure remained stable, ENTH was added to the subphase via a syringe. To distribute the protein the system was then mixed with the syringe. Addition of the protein resulted in an increase of the surface pressure. Figure 4.25 represents an exemplary time-resolved increase of π after injection of 1 μ M ENTH to a POPC/POPE/PIP₂ (70:20:10) monolayer at an initial surface pressure of 14.6 mN/m.

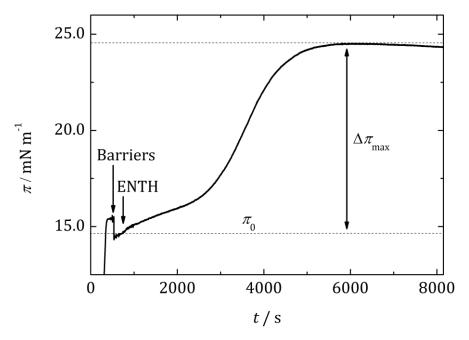


Figure 4.25: Exemplary plot of the surface pressure vs. time of a POPC/POPE/PIP $_2$ (70:20:10) monolayer after ENTH addition at 20 °C using a protein trough. When the desired surface pressure was reached (14.6 mN/m) the small Teflon barriers were implemented in the protein trough, resulting in a drop in surface pressure. When no change in surface pressure was observed, the ENTH domain was added via a small hole in the trough. Due to protein-monolayer interactions and penetration ability of the ENTH domain, the surface pressure increased reaching a maximum after 6000 s.

After the insert of the small protein trough barriers, a sudden decrease of π of about 1 mN/m was observed. Once π stabilized (usually within a few minutes), 1 μ M ENTH was injected into the subphase. Dependent on the initial surface pressure (π_0), π reached a maximum after 10-210 min. In general the time to reach the maximum increased with decreased initial surface pressure as the surface pressure change $\Delta\pi$ was higher at those conditions due to a looser lipid packing of the monolayer. The increase of π usually did not increased exponential, exhibiting local extrema. Directly after insertion of ENTH the π increased fast until it almost reached a plateau (Figure 4.25) or even decreased. Then π increased fast again. This behavior was observed several times and was attributed to the uneven distribution of the protein as the system was just mixed with a syringe directly after protein addition. Afterwards the ENTH was only distributed by diffusion.

To calculate the exclusion or critical surface pressure (π_c) at which no further protein insertion was possible into POPC/POPE (80:20) and POPC/POPE/PIP₂ (80-x:20:x,

with x = 1;5 and 10 mol%) monolayers, $\Delta \pi$ was plotted against π_0 (Figure 4.26). Interpolation yielded the intersection with the x-axis, the critical surface pressure.

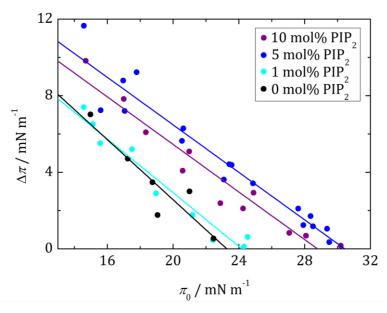


Figure 4.26: Monolayer penetration of the ENTH domain into POPC/POPE (80:20) and POPC/POPE/PIP₂ (80-x:20:x, with x=1,5,10) monolayers. Maximum increase in surface pressure ($\Delta\pi$) was monitored as a function of π_0 values. The subphase was 10 mM HEPES, 160 mM KCl, pH 7.4.

Figure 4.26 shows the $\Delta\pi$ at initial surface pressures of π_0 =14-30 mN/m for four different lipid mixtures. First of all, monolayers lacking PIP₂ were used to check the surface activity or the penetration ability of ENTH independently of the receptor binding. Towards this lipid mixture, the ENTH domain had a low intrinsic penetration capability with a π_c = 23.2 ±0.8 mN/m.

When the concentration of PIP₂ was increased to 1 mol%, the π_c just slightly increased to 24.1 ± 0.5 mN/m. At higher PIP₂ content, the penetration ability of ENTH into the monolayers increased (π_c = 30.3 ± 0.4 mN/m for 5 mol% and π_c = 28.7 ±0.6 mN/m for 10 mol%). The errors were given as weighted standard deviations of the fits. In Figure 4.27 the area per molecule at 20 mN/m (A) and the critical surface pressure after ENTH addition (B) were plotted against the PIP₂ concentration, showing the influence of the receptor lipid concentration as well as the lipid packing on the ENTH penetration ability. To ensure that the monolayer did not collapse at high pressures, the area per molecule at 20 mN/m was chosen to compare the lipid packing density. Considering the π_c value of the monolayers with 5 mol% (30.3 ± 0.4mN/m) to 10 mol% PIP₂ (28.7 ± 0.6 mN/m) no significant differences were observed indicating

that a maximum of the ENTH penetration was reached. Monolayers with 10 mol% PIP₂ showed a denser lipid packing compared to the monolayer with 5 mol% PIP₂.

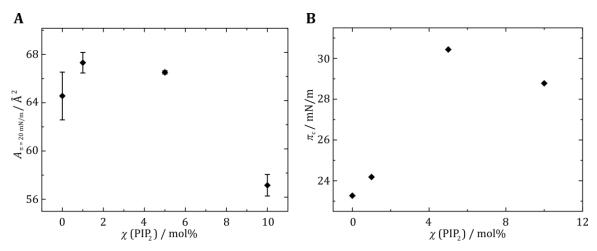


Figure 4.27: (A) Lipid packing against PIP₂ concentration and (B) the critical surface pressure against the PIP₂ content. The area per molecule at the beginning increases by addition of PIP₂ and then decreases with increased PIP₂ concentration. The critical surface pressure also first rises by the addition of PIP₂, reaches a maximum at 5 mol% and then decreases again indicating that at some point the lipid packing prevents further protein penetration.

In general, the $\Delta\pi$ values for lipid mixtures with 5 mol% PIP₂ were about 1 mN/m higher as the data measured on 10 mol% PIP₂ monolayers. Even though the scattering of the data points is quite high, this indicated a maximum of the penetration capability of ENTH into monolayers with 5 mol% PIP₂.

4.3.2 PS influence on the ENTH penetration ability

The influence of PS on the ENTH-PIP₂ interaction was already described earlier.^{8,26} Also RIfS measurements showed a higher binding affinity of ENTH to PIP₂ containing supported bilayers in the presence of PS (cf. Section 4.2.2). To investigate the influence of PS on the penetration capability of ENTH to PIP₂-containing monolayers, first of all the intermolecular interactions within POPC/POPE (80:20), POPC/POPE/POPS (60:20:20) and POPC/POPE/POPS/PIP₂ (55:20:20:5) monolayers were analyzed

with surface pressure-area isotherms on ultrapure water and HEPES as subphases. Averaged isotherms of these lipid compositions are shown in Figure 4.28.

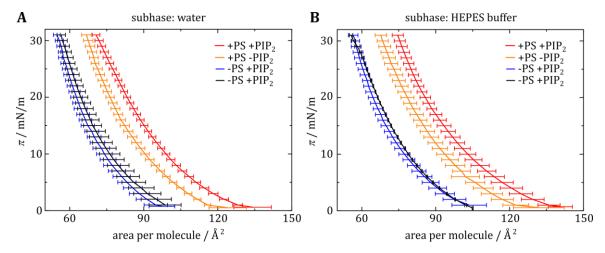


Figure 4.28: Averaged surface pressure-area per molecule isotherms of POPC/POPE (80:20, black, n=4), POPC/POPE/PIP2 (75:20:5, blue, n=3), POPC/POPE/POPS (60:20:20, orange, n=3) and POPC/POPE/POPS/PIP2 (55:20:20:5, red, n=3) at 20 °C. The subphases were ultrapure water (A) and HEPES buffer (B). The error bars represent the standard deviations of the area per molecule at the same surface pressure with a surface pressure deviation of 0.05 mN/m. The isotherms were recorded from $\pi=0$ to 31 mN/m, showing the influence of the PS on the lipid packing of the monolayers as the isotherms are shifted to higher molecule per area values. No significant difference between the isotherms on water or HEPES buffer were observed.

Independent of the lipid mixture the curve shape of all isotherms resemble one another. When decreasing the trough area, the surface pressure steadily increased and no co-existence region was observed (cf. Section 3.3.5). It was possible to compress the POPC/POPE and POPC/POPE/PIP2 monolayers further than the monolayers containing PS. The isotherm of the POPC/POPE/POPS/PIP2 monolayer showed the highest area per molecule values of the lipids for the same surface pressure compared to the other mixtures, indicating that POPS decreased the lipid packing of the monolayers. Especially in combination with PIP2 the lipid packing in the monolayers was decreased. Within the error margins the isotherms on HEPES buffer did not differ from the isotherms on water. Furthermore the lipid mixtures were stable up to the π values of 31 mN/m, allowing the measurement of the ENTH penetration capability. As al-

ready mentioned in chapter 4.3.1, the A_{20} values were determined to obtain information about the lipid packing (Table 4.9) and to compare them. The area per molecule values ($A_{\rm M}$) were averaged over all lipids in the ratio of their molecular fraction.

Table 4.9: A_{20} mean values with standard deviations obtained from the isotherms with the different lipid compositions. The subphases were ultrapure water and HEPES buffer. n is the number of recorded isotherms.

	Subphase: ultrapure water		Subphase: HEPES buffer	
Lipid composition	A_{20} /Å ²	n	$A_{20}/\text{Å}^2$	n
POPC/POPE (80:20)	64.6 ± 2.0	4	64.5 ± 2.0	4
POPC/POPE/POPS	76.3 ± 2.0	3	79.9 ± 4.0	5
(60:20:20)				
POPC/POPE/PIP2	63.1± 2.4	3	64.6 ± 0.1	3
(55:20:5)				
POPC/POPE/POPS/PIP2	83.3 ± 1.4	3	83.9 ± 4.2	6
(55:20:20:5)				

The A_{20} values show the influence of PS on the lipid packing within the monolayers. Comparing POPC/POPE (80:20) and POPC/POPE/PIP₂ (75:20:5) monolayers the lipid packing got looser when PS was added, indicated by the increase of the A_{20} value. Monolayers containing PS and PIP₂ even showed higher A_{20} values and consequently had a lower packing density.

Although, no collapse of the monolayers was observed until 31 mN/m, it was necessary to investigate the monolayer stability over time to ensure the accessibility of the receptor lipid during the experiment (cf. Section 3.3.5). Therefore, the surface pressure change ($\Delta\pi$) was recorded over time to obtain the percentage loss in surface pressure after 1 h for POPC/POPE/POPS (60:20:20) and POPC/POPE/POPS/PIP₂ (55:20:20:5) monolayers. Over 1 h the surface pressure increased about 1 ± 1% for the monolayers without PIP₂ and decreased about 6 ± 2% when PIP₂ was present, indicating sufficient monolayer stability.¹¹⁵

With these stable monolayers penetration experiments of ENTH were possible. The experimental conditions were similar to the experiments shown in chapter 4.3.1. After the adjusted surface pressure was reached, the small barriers were implemented. When almost no change in π was observed, indicating stable conditions, 1 μ M ENTH was added into the subphase and then mixed with the syringe. To investigate how the lipid PS influences the penetration ability of ENTH, the increases of $\Delta\pi$ for POPC/POPE/POPS (60:20:20) and POPC/POPE/POPS/PIP2 (55:20:20:5) monolayers dependent on the initial surface pressure (π_0) were analyzed. At low π_0 the change in surface pressure was greater compared to higher π_0 values due to the lower packing of the monolayer. When PS and PIP2 were present the maximum of π was reached after 60-270 min, whereby the time usually increased with decreasing π_0 . Some irregularities in the surface pressure increase were attributed to inhomogeneous ENTH distribution (cf. chapter 4.3.1)

To obtain the penetration capability of ENTH into POPC/POPE/POPS (60:20:20) and POPC/POPE/POPS/PIP₂ (55:20:20:5) monolayers, the data was plotted in a $\Delta\pi$ – π_0 diagram. Furthermore the data for the POPC/POPE/PIP₂ (cf. section 4.3.1) lipid mixture was added. Consequently , this allowed the comparision of all lipid mixtures and the influence of PS on the penetration ability. By interpolating π_c was determined (Figure 4.29).

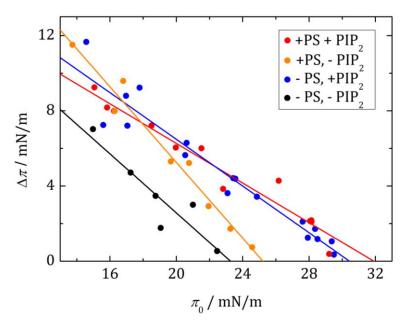


Figure 4.29: Monolayer penetration of the ENTH domain into POPC/POPE (80:20,), POPC/POPE/POPS (60:20:20), POPC/POPE/PIP₂ (75:20:5) and POPC/POPE/POPS/PIP₂ (55:20:20:5) monolayers. Maximum increase in surface pressure ($\Delta\pi$) was monitored as a function of π_0 values. The subphase was 10 mM HEPES, 160 mM KCl, pH 7.4.

The x-intercepts in Figure 4.29 of the fits reveal a mean $\Delta \pi_c$ of 30.3 ± 0.4 mN/m for POPC/POPE/PIP2 monolayers. Penetration experiments of **ENTH** into POPC/POPE/POPS/PIP₂ monolayers yielded a π_c of 32.0 ±0.9 mN/m. The fit with PS shows a slighter decrease of $\Delta\pi$ with increasing π_0 . Within the scatter of the measurements the values are similar. The experiments without PIP2 showed that the ENTH penetration ability decreased down to 23.2 ± 0.8 mN/m for POPC/POPE monolayers and 25.0 ±1.0 mN/m for the POPC/POPE/POPS monolayers. The errors were given as weighted standard deviations of the fits. This reveals that the penetration capability of the ENTH domain is clearly decreased when PIP₂ is lacking even though PS is present. Furthermore, the penetration ability of ENTH into monolayers composed of POPC/POPE/POPS/PIP2 only slightly increased compared to POPC/POPE/PIP2 monolayers, indicating that PS has almost no impact on the penetration capability of ENTH.

4.4 Adhered GUVs as a membrane model to investigate proteinmembrane interactions

The surface activity and especially the insertion of the amphipathic helix of ENTH resulted in an increase of the surface pressure in monolayer penetration experiments (cf. chapters 4.3.1 and 4.3.2). It was shown that the lipid composition has an influence on the exclusion pressure, although the effect of PS was only slight. Another possibility to investigate the activity of ENTH is the analysis of the activity as a function of membrane tension in bilayers.³⁶

In this work, GUVs were adhered to a biotin-PEG functionalized glass surface to immobilize them. By using fluorescence microscopy the vesicle geometry was determined, which could be evaluated to get the membrane tension. Afterwards, ENTH was added to see how the protein effects the membrane tension. Based on this the influence of the lipid head group Phosphatidylserine (PS) on the binding behavior of ENTH as a function of lateral membrane tension was investigated.

4.4.1 Theoretical model

Adhesion of vesicles or cells were already used to investigate membrane-protein interactions. ¹¹⁶ In this thesis, the membrane tension of adhered GUVs, generated as described in chapter 3.2.2, was analyzed to determine the impact of ENTH on the lateral membrane tension as a function of the lipid composition.

Attractive forces between the GUV and the surface result in the adhesion of the GUV. During this process the volume of the GUV remains constant under iso-osmolar conditions due to low water permeability of membranes. In contrast to this, the vesicle shape changes from a sphere with a radius $r_{\rm v}$ to a spherical cap with a radius $r_{\rm ad}$. Additionally, an adhesion area or contact area with a radius $r_{\rm i}$ is observed (Figure 4.30). From the geometry of the adhered vesicles it is then possible to calculate the membrane tension.

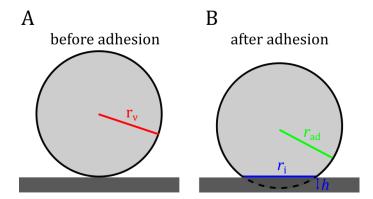


Figure 4.30: Schematic illustration of a GUV before (A) and after (B) adhesion to a surface. The GUV geometry changes from a sphere with a radius r_v to a spherical cap with a radius r_{ad} and a contact area (radius r_i). h describes the height of the spherical cap in (B), which is missing to get a complete sphere.

Marsh theoretically described the observable tension of a fluctuating membrane using equation (4.5):117

$$\sigma = \Delta \sigma - \frac{K_{\rm A} k_B T}{2 A_{\rm v} \kappa} \sum_{q} \left(\frac{\sigma}{\sigma q^2 + \kappa q^4} \right) \tag{4.5}$$

 K_A is the area compressibility modulus, k_B the Boltzmann constant, T the absolute temperature, κ the bending rigidity and A_V the area of a free vesicle. $\Delta \sigma$ depicts the increase in membrane tension of a tension free planar membrane and is proportional to the fractional area expansion ($\Delta A/A_V$) and K_A (equation (4.5)).

$$\Delta \sigma = K_{\rm A} \, \frac{\Delta A}{A_{\rm v}} \tag{4.6}$$

Integration with $q_{\min} = \pi/\sqrt{A_{\rm M}}$ and $q_{\max} = \pi/\sqrt{A_{\rm V}}$, where $A_{\rm M}$ is the area per lipid, the observable tension σ of a fluctuating membrane is then calculated with equation (4.7).

$$\sigma = \Delta \sigma - \frac{K_{A} k_{B} T}{8\pi \kappa} \ln \left(\frac{\frac{\sigma A_{V}}{\pi^{2} \kappa} + 1}{\frac{\sigma A_{M}}{\pi^{2} \kappa} + 1} \right)$$
(4.7)

The denominator of the logarithm becomes close to unity when the membrane tension is low ($\sigma \ll \pi^2 \kappa / A_M \approx 2 \cdot 10^3 \, \text{mN/m}$), simplifying equation (4.7):

$$\sigma = \Delta \sigma - \frac{K_{\rm A} k_B T}{8\pi \kappa} \ln \left(\frac{\sigma A_{\rm v}}{\pi^2 \kappa} + 1 \right) \tag{4.8}$$

Under the assumption of $K_A \gg \sigma \gg \pi^2 \kappa / A_v \approx 10^{-6}$ mN/m, which is valid for relevant membrane tensions, the membrane tension is described with equation (4.9) as $(\sigma A_v) / (\pi^2 \kappa)$ is $\gg 1$.

$$\sigma = \Delta \sigma - \frac{K_{\rm A} k_B T}{8\pi \kappa} \ln \left(\frac{\sigma A_{\rm v}}{\pi^2 \kappa} \right) \tag{4.9}$$

When equation (4.6) and (4.9) are combined, the membrane tension of an adhered vesicle can be calculated using equation (4.10)

$$\sigma = K_{\rm A} \frac{\Delta A}{A_{\rm v}} - \frac{K_{\rm A} k_B T}{8\pi\kappa} \ln\left(\frac{\sigma A_{\rm v}}{\pi^2 \kappa}\right) \tag{4.10}$$

The membrane tension of the adhered GUV minus the tension of the free vesicle (σ_0) enables to compare the obtained membrane tension with different reference states from previous studies.^{118,119} Thereby, the flattening of fluctuations and the dilation of the area per molecule has to be considered.

$$\frac{\Delta A}{A_{\rm v}} = \frac{\sigma - \sigma_0}{K_{\rm A}} - \frac{k_B T}{8\pi\kappa} \ln\left(\frac{\sigma}{\sigma_0}\right) \tag{4.11}$$

The fractional area expansion term can be calculated from the geometries of the GUVs before and after adhesion (Figure 4.30). The fractional area expansion is calculated with equation (4.12).

$$\frac{\Delta A}{A_{\rm v}} = \frac{A_{\rm ad} - A_{\rm v}}{A_{\rm v}} = \frac{4\pi r_{\rm ad}^2 - 2\pi r_{\rm ad}h + \pi r_{\rm i}^2 - 4\pi r_{\rm v}^2}{4\pi r_{\rm v}^2}$$
(4.12)

h is the height of the spherical cap (Figure 4.30). Considering a constant volume of the vesicle during the adhesion process ($V_v = V_{ad}$) results in equation (4.15).

$$V_V = \frac{4}{3} \pi r_{\rm v}^3 \tag{4.13}$$

$$V_{\rm ad} = \frac{4}{3} \pi r_{\rm ad}^3 - \frac{\pi h}{6} (3r_{\rm i}^2 + h^2)$$
 (4.14)

$$\frac{4}{3}\pi r_{\rm v}^3 = \frac{4}{3}\pi r_{\rm ad}^3 - \frac{\pi h}{6}(3r_{\rm i}^2 + h^2) \tag{4.15}$$

Determination of the radius of the immobile, free vesicles r_v and the height of the spherical cap h is experimentally complicated to realize and can be described as a function of r_i and r_{ad} :

$$h = r_{\rm ad} - \sqrt{r_{\rm ad}^2 - r_{\rm i}^2} \tag{4.16}$$

$$r_{\rm v} = \sqrt[3]{r_{\rm ad}^3 - \frac{r_{\rm ad} - \sqrt{r_{\rm ad}^2 - r_{\rm i}^2}}{8} \left(3r_{\rm i}^2 + \left(r_{\rm ad} - \sqrt{r_{\rm ad}^2 - r_{\rm i}^2}\right)^2\right)}$$
(4.17)

Normalization of the radii to r_{ad} gives dimensionless variables and substitution of them in equation (4.12) results in:⁷⁹

$$\frac{\Delta A}{A_{v}} = \frac{4\pi - 2\pi \left(1 - \sqrt{1 - \left(\frac{r_{i}}{r_{ad}}\right)^{2}} + \pi \left(\frac{r_{i}}{r_{ad}}\right)^{2}}{4\pi \left(\sqrt[3]{1 - \frac{2 - \left(\frac{r_{i}}{r_{ad}}\right)^{2}}{\sqrt{1 - \left(\frac{r_{i}}{r_{ad}}\right)^{2}} - 2\sqrt{1 - \left(\frac{r_{i}}{r_{ad}}\right)^{2}}}}\right)^{2} - 1$$
(4.18)

Finally, the lateral membrane tension of the adhered GUV can be calculated numerically by the combination of equation (4.11) and (4.18).

$$\frac{\sigma - \sigma_0}{K_A} + \frac{k_B T}{8\pi\kappa} \ln\left(\frac{\sigma}{\sigma_0}\right) = \frac{4\pi - 2\pi(1 - b) + \pi\left(\frac{r_i}{r_{ad}}\right)^2}{4\pi\left(\sqrt[3]{1 - \frac{2 - \left(\frac{r_i}{r_{ad}}\right)^2 b - 2b}{4}}\right)^2} - 1$$
(4.19)

with
$$b = \sqrt{1 - \left(\frac{r_i}{r_{ad}}\right)^2}$$
.

Rawicz *et al.* measured the area compressibility modulus and the bending rigidity for a DOPC membrane via micropipette aspiration and yielded a K_A of 265 ± 18 mN/m and a κ of 21 k_BT ((0.85 ± 0.1)·10⁻¹⁹ J).¹²⁰ For GUVs weakly adhered to a solid support, pre-tensions (σ_0) were measured to be in a range of 10⁻¹ - 10⁻³ mN/m.¹²¹ Using the different pre-tension regimes, the lateral tension as a function of r_i/r_{ad} first increases slow and then fast when the dilation exceeds the excess area, which is stored in membrane undulations.^{36,119}

Calculation of the membrane tension of the adhered GUVs before and after ENTH addition hence can yield the protein activity as a function of membrane tension dependent of the lipid composition.

4.4.2 Adhesion of biotinylated GUVs on PEGylated surfaces

Adhered GUVs were already used to study membrane-protein interaction or phase separation of lipid bilayers.¹²²⁻¹²⁴ Gleisner *et al.* used Avidin coated glass surfaces to immobilize biotinylated GUVs and therewith analyzed the ENTH-PIP₂ interactions as a function of membrane tension.³⁶ They performed these experiments with adhered GUVs composed of DOPC/DOPE/cap-biotin-DOPE/Atto488-DPPE (67:30:2:1). The membrane tension of these GUVs were regulated by the magnesium chloride (MgCl₂) concentration in the buffer.

In this work, a similar system was used to investigate how the lipid head group PS influences the ENTH binding considering the membrane tension of the GUVs before and after ENTH addition. As PS and PIP₂ exhibit negative charges and are known to

cluster in the presents of divalent ions, ^{125,126} the buffer and the functionalization of the surface were changed. MgCl₂ was removed from the buffer and NeutrAvidin was used instead of avidin to immobilize the vesicles. ³⁶ NeutrAvidin is known to reduce the unspecific interactions compared to avidin. ¹²⁷

Direct functionalization of the surface with NeutrAvidin by simple incubation on a hydrophilized glass substrate led to an inhomogenous occupancy of the protein on the glass slides (Figure 4.31 A), which can be explained by less unspecific interactions of NeutrAvidin with the hydrophilic glass surface. As a homogenous coverage with the protein was necessary to obtain adhesion of biotinylated GUVs, the glass slide surfaces were silanized and PEGylated (PEG= Poly(ethylene glycol)) (cf. 3.2.2). Using biotin-linked PEG (b-PEG) ensured an affinity of NeutrAvidin to the surface and due to four biotin binding sides, additionally binding to the biotinylated vesicles was guaranteed. To ensure the accessibility of the biotin groups on the surface, methoxy functionalized PEGs (m-PEG) were added (ratio m-PEG/b-PEG 1:2). Similar PEG-functionalizations were already used in previous studies. 128,129 With the new functionalization method incubation with DyLight® 594 labeled Neutravidin resulted in a homogenous NeutrAvidin coverage (Figure 4.31 B).

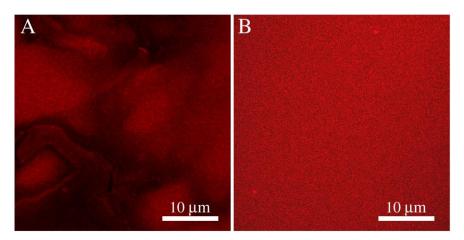


Figure 4.31: Fluorescence images of (A) hydrophilized and (B) biotin-PEG-functionalized glass slides after incubation with DyLight® 594 labeled Neutravidin and subsequent rinsing with buffer. Without PEG-functionalization the protein was removed from the surface by rinsing with buffer resulting in an inhomogeneous coverage of NeutrAvidin. Due to the strong biotin-NeutrAvidin interaction the biotin an even distribution of the labeled protein on the PEG-functionalized surface was observed.

Figure 4.31 clearly shows the specific NeutrAvidin-biotin interaction. In Figure 4.31 A the hydrophilized glass slides were just directly incubated with the protein. After rinsing with buffer NeutrAvidin was partly removed. When the glass slides were first silanized with (3-Glycidyloxypropyl)trimethoxysilane (GOPTS) and then PEGylated with a mixture of m-PEG and b-PEG (cf. chapter 3.2.2), the NeutrAvidin strongly bound to the surface, which was not removed by rinsing with buffer. After the successful coverage of the surface with NeutrAvidin biotinylated GUVs could be added to the substrates. **GUVs** composed of DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1) turned out to be large and stable enough to perform these experiments. Higher concentrations of DOPS or DOPE in combination with the b-PEG-NeutrAvidin surface resulted in instable GUVs, preventing the analysis of adhesion areas and consequently the membrane tension.

Incubation of these GUVs on the PEG-surface for 15 min obtained immobilized GUVs with small adhesion areas (Figure 4.32). Fluorescence images of the adhered GUVs were taken by means of Spinning Disc confocal laser microscopy (SDCLM). To obtain the radii of the adhered GUV and the contact area (cf. 3.2.2), z-stacks of the GUVs with 150-250 nm between the slides were recorded.

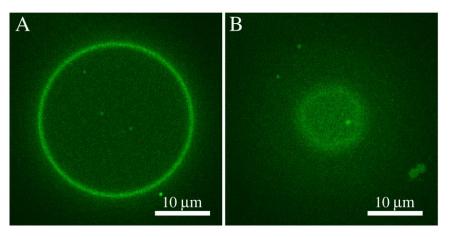


Figure 4.32: z-stacks of an adhered GUVs obtained by SDCLM. (A) Cross section of the adhered GUV with the radius r_{ad} . (B) Fluorescence image of the contact area plane to determine r_i . Lipid composition: DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1).

Calculation of the lateral tension from the geometry of the vesicles is only possible if the ratio r_i/r_{ad} is constant over the experimental time. The radii were determined as

described in chapter 3.2.2. Also after 2 h the adhered vesicles showed sufficient stability as the ratio r_i/r_{ad} did not change within the errors (Figure 4.33).

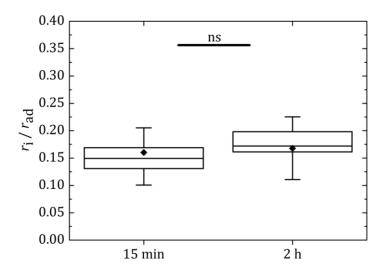


Figure 4.33: Ratios of the contact area radius to the adhered vesicle radius (r_i/r_{ad}) as a function of time for GUVs (DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1)) adhered to a NeutrAvidin covered surface, indicating the sufficient stability of the GUVs after 2 h. Statistical Mann-Whitney U test: ns - not significant (p = 0.3).

25 GUVs were analyzed after subsidence (15 min) and 16 GUVs after 2 h, yielding r_i/r_{ad} values of 0.16 ± 0.04 and 0.17 ± 0.04 , respectively. The lateral membrane tension of the adhered GUVs was then calculated from the ratio of r_i/r_{ad} using equation (4.19) with an area compressibility modulus of $K_A = 265$ mN/m, a bending rigidity of $\kappa = 21$ k_B T and a pre-tension of $\sigma_0 = 10^{-2}$ mN/m (Figure 4.34).

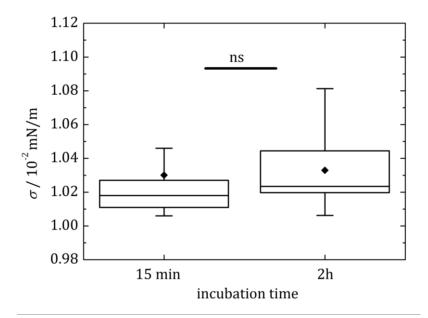


Figure 4.34: Lateral membrane tension of adhered GUVs on PEG-biotin-NeutrAvidin functionalized glass substrates after 15 min and 2 h, showing that the membrane tension did not change during that time. Lipid composition: DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP $_2$ /Atto488-DPPE (76:10:10:2:1:1). Statistical Mann-Whitney U test: ns - not significant (p = 0.2).

The lateral membrane tension of the adhered GUVs after subsidence was $\sigma = (1.03 \pm 0.04) \cdot 10^{-2}$ mN/m and after 2 h the membrane tension of the GUVs was determined with $\sigma = (1.03 \pm 0.03) \cdot 10^{-2}$ mN/m. The membrane tension of the GUVs did not changed significantly, indicating that the GUVs remained stable over time with quite low lateral membrane tensions of the adhered GUVs.

Since the regulation of the lateral membrane tension of the GUVs was not able with MgCl₂ as it was previously done in the literature,³⁶ the ratio of m-PEG to b-PEG on the substrate surface was varied to verify, if the membrane tension can be controlled by this. Changing the ratio of m-PEG/b-PEG from 1:1 to 1:3 did not change r_i/r_{ad} of the vesicles (Figure 4.35). Hence, it was not possible to regulate the lateral membrane tension with the surface functionalization (Figure 4.36).

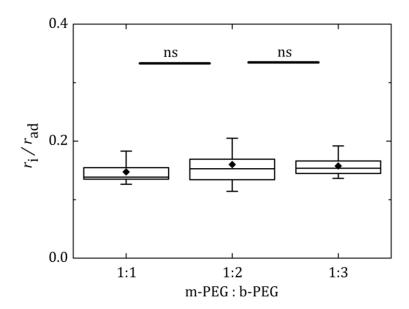


Figure 4.35: Ratio of the contact area radius and the adhered vesicle radius (r_i/r_{ad}) of the adhered GUVs (DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1)) immobilized on PEG-Neutravidin functionalized glass with different ratios of b-PEG and m-PEG. Statistical Mann-Whitney U test: ns - not significant (p > 0.3).

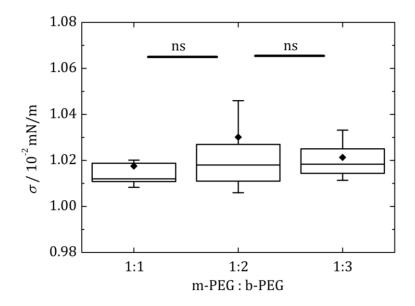


Figure 4.36: Lateral membrane tension of adhered GUVs on PEG-NeutrAvidin functionalized glass substrates with different ratios of m-PEG and b-PEG used for immobilization of the GUVs (DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP $_2$ /Atto488-DPPE (76:10:10:2:1:1)). This shows that the membrane tension did not change by varying the PEG ratio. Statistical Mann-Whitney U test: ns - not significant (p > 0.2).

For each surface functionalization 16-25 GUVs were evaluated. The membrane tensions were determined to be $(1.02 \pm 0.02)\cdot 10^{-2}$ mN/m (m-PEG/b-PEG 1:1), $(1.03 \pm 0.04)\cdot 10^{-2}$ mN/m (m-PEG/b-PEG 1:2) and $(1.02 \pm 0.01)\cdot 10^{-2}$ mN/m (m-PEG/b-PEG 1:3). Although the membrane tension of the adhered GUVs could not be regulated by the surface functionalization, they showed stability over 2 h under iso-osmolar conditions. Therefore, ENTH was added to investigate the effect of the protein binding on the membrane tension.

4.4.3 Effect of ENTH addition to adhered GUVs

Since the ratio of b-PEG and m-PEG showed no influence on the membrane tension of the adhered GUs, the following experiments were done with a ratio of 2:1 (b-PEG/m-PEG). Moreover, it already was shown that with this ratio the PEG surface was covered homogenously with NeutrAvidin.

Adhered GUVs composed of DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP $_2$ /Atto488-DPPE (76:10:10:2:1:1) were incubated with 1 μ M ENTH for 2 h. Every 30 min the solution was stirred to ensure a homogenous distribution of the protein. Afterwards, z-stacks of the ENTH incubated adhered GUVs were taken. Even though the amount of GUVs in the solution was not counted, it was obvious that less GUVs were found in the solution. Furthermore, a lot of lipid material was observed. Additionally imaging of the GUVs showed the rupture of the vesicles (Figure 4.37), indicating that the protein affects the vesicle stability.

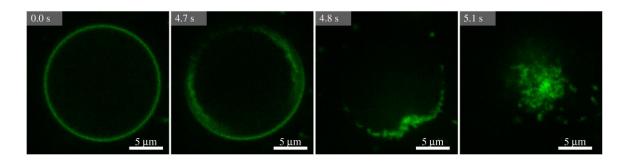


Figure 4.37: Time series of one z-plane of an adhered GUV incubated with 1 μ M ENTH (2 h). During imaging the vesicle ruptures. Lipid composition: DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1).

Since the stability of the adhered GUVs without ENTH addition over 2 h was proven (Figure 4.33), rupturing of the GUVs seemed to be an effect of the ENTH binding to the PIP $_2$ doped GUVs. In order to investigate if lower concentration still affects the GUV stability, the ENTH concentration was decreased to 0.3 μ M. Even at this ENTH concentration the adhered GUVs ruptured during imaging (Figure 4.38). Decreasing the laser intensity did not change this as well.

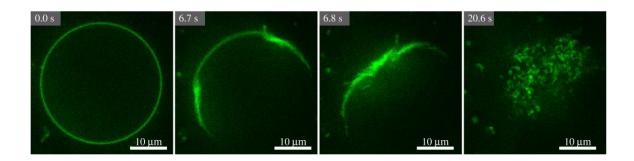


Figure 4.38: Time series of one z-plane of an adhered GUV incubated with 0.3 μ M ENTH (2 h). During imaging the vesicle ruptures. Lipid composition: DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1).

These results showed that the adhered GUVs ruptured during incubation with 1 μ M and 0.3 μ M ENTH, indicating that the ENTH helix insertion induce the destabilization of adhered GUVs.

5 Discussion

To understand how the lipid composition influences the binding behavior of the epsin *N*-terminal homology (ENTH) domain different artificial membrane models and techniques were utilized. Reflectometric interference spectroscopy (RIfS) measurements allow to analyze the adsorption of ENTH and its dissociation constant dependening on the lipid composition. Using atomic force microscopy (AFM) the topology of the membrane surfaces with and without ENTH was analyzed. Moreover, the lipid dependent ENTH-PIP₂ interaction was analyzed by Langmuir-Blodgett trough experiments and fluorescence microscopy. Monolayer penetration experiments allowed to obtain information about the insertion capability of the protein into a monolayer with different lipid compositions by detecting surface pressure changes. Protein activity as a function of membrane tension was investigated on adhered GUVs. In this chapter the results described in chapter 4 were compared to the literature and discussed in more detail.

5.1 Characterization of the ENTH binding to PIP₂ containing solid supported membranes

Phosphatidylinositol-(4,5)-bisphosphate (PIP₂) is the specific receptor lipid for the ENTH domain and consequently involved in the initial step of clathrin-mediated endocytosis (cf. chapter 1.1). In literature, a 1:1 binding of ENTH to PIP₂ is described, which would predict a direct correlation of the amount of bound protein with the receptor concentration.⁷ Stahelin *et al.* already showed that at low PIP₂ concentration (0.5%) the dissociation constant drastically increases compared to higher PIP₂ content (3%).²⁵ However, an influence of the PIP₂ concentration on the ENTH-membrane interaction was not investigated systematically yet and would clarify the relevance of its amount in the cytosolic bilayer leaflet.

5.1.1 Asymmetric distribution of PIP₂ in solid supported lipid bilayers and the influence on ENTH binding

The binding behavior of ENTH to solid supported lipid bi- and monolayers was analyzed using RIfS. To measure the influence of PIP₂ on the protein adsorption, planar and defect-free mono- and bilayers were required. For a pure POPC bilayer an optical thickness of $\Delta OT = 6.4 \pm 0.2$ nm was determined, which can be translated into a physical membrane thickness of $d_{\rm M} = 4.4 \pm 0.2$ nm (equation (3.7)). Kŭcerke *et al.* measured a physical thickness of $d = 3.98 \pm 0.08$ nm for a POPC lipid bilayer.⁸⁸ Other studies revealed membrane thicknesses of 3.68-4.60 nm,¹³⁰⁻¹³² which are in agreement with the determined physical thickness of the POPC bilayer in this work.

Addition of PIP₂ did not change the thickness, ranging from 4.3 to 4.5 nm without a correlation to the PIP₂ content. The size of the PIP₂ head group implies that it protrudes into the aqueous area when the structure is erected. These differences probably get lost in the detection range or in the errors of the measurements. Furthermore a possible hydration of the PIP₂ head group would change the refractive index. An increased refractive index abolishes its exposed character and impedes its detection (cf. equation (3.7)). However hydration of lipid head groups already was shown to decrease the refractive index from ~1.47 to the refractive index of water (n = 1.33). Thus a limited detection range and the errors of the measurements seems more likely to detect significant changes of POPC and POPC/PIP₂ membranes. Physical thicknesses of POPC (PIP₂ hilayers were previously measured via PIFS under

Physical thicknesses of POPC/PIP₂ bilayers were previously measured via RIfS under the same conditions showing values in the same range.^{79,134} This confirmed the successful formation of POPC and POPC/PIP₂ bilayers.

On hydrophobic functionalized SiO₂ wafers spreading of SUVs resulted in monolayers (cf. chapter 4.1.2). For a POPC monolayer the physical thickness of $d_M = 1.6 \pm 0.2$ nm was determined. With increased PIP₂ concentration the d_M -values stayed approximately constant within the error margin. The physical thicknesses varied from 1.4 to 1.8 nm, but no tendency considering the PIP₂ content was observed. Since 4 nm of layer thickness is typical for a bilayer, monolayer thicknesses of about 2 nm are expected.^{88,130} Rossi *et al.* calculated an average thickness of 2.2 ± 0.3 nm for an egg-PC monolayer by means of surface plasmon resonance spectroscopy (SPR).¹³⁵ The values

in this work are slightly smaller. Minor variations can be explained by the different techniques and surface functionalization procedures. In general the RIfS experiments revealed the successful formation of bi- and monolayers on hydrophilized and hydrophobic functionalized silicon surfaces, respectively.

Besides the RIfS experiments, fluorescence images were taken to prove the homogeneity and the mobility of the membranes. Therefore, the lipid bound fluorophore Texas Red® 1,2-dihexadecanoyl-sn-glycero-3-phosphoetanolamine (TxR) was used to stain the membrane. Spreading of SUVs on hydrophilic and hydrophobic functionalized SiO₂ substrates resulted in bi- and monolayer formation, respectively. The fluorescence images of model membranes showed a predominant homogenous distribution of the fluorophore without significant defects. Fluorescence recovery after photobleaching (FRAP) experiments revealed that both systems were mobile, but the diffusion coefficients differed. For the POPC/TxR (99.5/0.5) and POPC/PIP2/TxR (99.5-x/x/0.5, with x = 1, 2, 5, 7, 8 and 10 mol%) bilayers diffusion coefficients of $D = 1.8 - 2.5 \,\mu\text{m}^2/\text{s}$ were measured with the FRAP method. Compared to that, the monolayers exhibited diffusion coefficients from $D = 0.1 \,\mu\text{m}^2/\text{s}$ to $0.4 \,\mu\text{m}^2/\text{s}$. In either case no correlation between the PIP2 content and the diffusion constant was observed. Aditionally, the diffusion constant decreased from bi- to monolayers by a factor of about 7 due to the interactions with the substrate. 136 Between the bilayer and the substrate surface a thin water layer of about 1 nm exists, reducing the interaction with the substrate compared to monolayers. 137-139

Braunger *et al.* determined diffusion coefficients of about 3 μ m²/s for both, bi- and monolayers using tetramethylrhodamine-labeled PIP₂.97 However, they used a dodecyl-trichlorosilane based functionalization for the monolayers. The silane might interact less with the PIP₂ thereby explaining the same mobility of the fluorophore in bi- and monolayers. As they used labeled PIP₂ instead of TxR, the diffusion might also slightly differ. Furthermore, diffusion coefficients can be influenced by lipid composition, lipid packing, temperature and substrate preparation^{100,136,140-142} which could explain the differences. Baumann *et al.* measured diffusion coefficients of POPC/PIP₂ lipid bilayers showing the influence of PIP₂ on the mobility in the bottom leaflet.¹⁴³ Measurements with labeled POPC (1-Palmitoyl-2-{12[7-nitro-2-1,3-benzoxadiazol-4-yl)amino]dodecanoyl}-*sn*-glycero-3-phosphocholine, NBD-POPC) revealed diffusion

constants of $D=1.84-2.46~\mu m^2/s$, comparable to the values observed in this thesis. They changed the PIP2 content from 1% to 10%, which did not influence the mobility of NBD-POPC. Using labeled PIP2 did not change D, but the mobile fractions were significantly reduced, indicating the interactions of PIP2 with the substrate. Compared to that, the mobile fractions (F_m) of NBD-POPC did not change ($F_m=84-91\%$) with varied PIP2 content. In this work, the mobile fractions of TxR in the bilayers showed a slight decrease from $98\pm1\%$ (2~mol% PIP2) to 93 ± 1 (10~mol% PIP2) in contrast to the monolayers with $F_m=98\pm2\%$ for all lipid compositions (cf. Table 4.2). The decreased mobile fractions with increased PIP2 concentration in the bilayers indicated a higher interaction of PIP2 with the hydrophilic substrate. The head groups of the bilayer are mainly interacting with the substrate. PIP2 consists of a larger and highly negatively charged head group compared to POPC (Figure 5.1). Increasing the concentration hence can lead to an increased lipid-surface interaction compared to neutral lipids. In the monolayers the increase of PIP2 does not affect the mobility as only similar hydrocarbon chains were added.

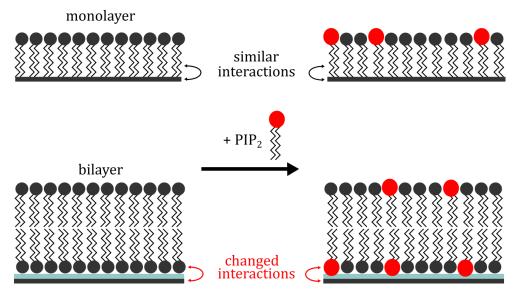


Figure 5.1: Schematic illustration of the different interactions of lipids with substrate surfaces. In the monolayers the increase of PIP₂ does not affect the interactions between each other due to similar structures of the fatty acid chains. Compared to that increasing of the PIP₂ content in bilayers results in increased interactions of the PIP₂ lipid head groups.

Baumann *et al.* also suggest that a larger hydration shell and charge repulsion of PIP₂ result in the reduced bilayer mobility. They also assume that interactions of pits in the negatively charged substrate and the head group lead to a reduction of the mobile fraction.¹⁴³

In this thesis the mobile fractions in the bilayers indicate a higher interaction between the substrate and the bottom leaflet with increased PIP₂ concentration. This may influence the distribution of the receptor lipid within both leaflets of the membrane. Thus, the accessibility of PIP₂ for ENTH binding could be affected. Binding studies of ENTH to supported bilayers via RIfS were reported previously,³⁶ but the influence of the PIP₂ concentration is still not completely understood. For an exact analysis, the PIP₂ molecules have to be entirely accessible. Therefore, besides the experiments on bilayers, also RIfS measurements were performed on monolayers to see if the substrate surface influences the PIP₂ distribution.

ENTH binding experiments

ENTH binding experiments were performed in order to investigate the influence of the receptor lipid concentration on the binding affinity. Incubation of 1 μ M ENTH to pure POPC bi- and monolayers proved the specificity of the ENTH-PIP2 interaction as no increase of the ΔOT was observed after protein addition (Figure 4.7). Itoh *et al.* revealed by co-sedimentation assays that only in the presence of PIP2 a high amount of ENTH binds to liposomes. Liposomes with other inositol species or negatively charged lipids showed a considerably decreased amount of bound protein or even no interaction.²⁴ Kweon *et al.* moreover elucidated with electron paramagnetic resonance (EPR) studies the importance of PIP2 binding for the amphipathic helix formation.¹⁴⁵

After the blind experiment with pure POPC, the binding of ENTH to PIP₂-containing bilayers was analyzed. At low PIP₂ concentrations almost no ENTH binding was visible. An increase of the PIP₂ content led to higher ΔOT values indicating an increased ENTH binding to PIP₂-containing bilayers. Binding of ENTH to PIP₂ and consequently conformational changes in the protein domain leads to the formation of a newly formed helix as well as a binding pocket for the receptor lipid.²⁵ Based on this a 1:1

binding of ENTH to PIP2 is assumed. This is in good agreement with the observed linear increase in ΔOT with rising PIP₂ concentration (cf. Figure 4.9). Even at a PIP₂ concentration of 10 mol% and a $\Delta OT_{\rm ENTH}$ of 1.2 + 0.4 nm no obvious saturation of the curve was observed indicating an incomplete protein occupancy on the surface. Nevertheless, it should be considered that with the RIfS technique the data is averaged over an area of 1 mm². Inhomogeneities hence can lead to underestimation of the physical thickness. Although a homogenous PIP2 distribution was observed by fluorescence microscopy previously, 79,97,134 it is possible that nanoscopic clusters of PIP₂ remained undetected.²² These clusters would prevent a 1:1 stoichiometry due to the size of the protein. Besides this, interactions of the lipids with the surface can influence the protein binding (cf. Figure 5.1). 136 Gleisner achieved full coverage of ENTH on POPC/PIP₂ (90:10) bilayers at a protein concentration of 3 μM using RIfS. At 1 μM a ΔOT of about 1.2 nm was determined,⁷⁹ thus the results in this work are consistent with the literature. As the values were obtained via RIfS, they should be comparable to the data in this thesis (for $c(PIP_2) = 10 \text{ mol}\%$) shown in Figure 4.9. Considering the crystal structure²² and a maximum coverage of proteins (56%) according to the scaled particle theory, 146 a full coverage should be achieved at about $\Delta OT_{\text{ENTH}} = 2.0 - 2.8 \text{ nm}$ ($d_{\text{ENTH}} = 1.4 - 2.0 \text{ nm}$), dependent on the orientation of the protein on the surface. The crystal structure of the bound ENTH state is only known in the presence of the head group of PIP₂, the inositol-1,4,5-triphosphate (IP₃). However, only binding to a membrane results in the formation of a stable amphipathic helix of ENTH¹⁴⁵ and consequently can also change the insertion and protein height. Moreover, interactions of ENTH and PS also can lead to an altered protein structure.

A full coverage of the protein was not yet reached at 10 mol% PIP₂ in a bilayer. PIP₂ has an area per molecule of about 0.7 nm². With respect to the proteins footprint (about 16 nm²)^{22,30} even in the presence of some small clusters, the surface would be completely covered at 10 mol% PIP₂. This in turn shows that a certain amount of PIP₂ seems to be inaccessible. The FRAP experiments on bi- and monolayers, discussed in the previous section, also indicated the uneven distribution of PIP₂ in both leaflets of the bilayer.

PIP₂ exhibits a larger head group than POPC and was shown to have an exposed position in the membrane.^{44,147} Hence a high curvature in the SUV might provoke its location in the outer leaflet of the SUV. After spreading, the outer leaflet of the SUV is facing to the surface,¹⁴⁸ resulting in the accumulation of PIP₂ in the bottom leaflet. This would support the assumption that some PIP₂ molecules are inaccessible for ENTH binding. In monolayers this asymmetric distribution of PIP₂ is prevented due to the missing second leaflet. For this reason ENTH binding experiments were repeated on PIP₂-doped monolayers.

Binding studies of ENTH to POPC/PIP₂ monolayers showed an increase of the optical thickness with increased lipid receptor concentration from $\Delta OT_{\rm ENTH} = 0.5 \pm 0.1$ nm (1 mol% PIP₂) to $\Delta OT_{\rm ENTH} = 2.2 \pm 0.4$ nm (10 mol% PIP₂). Compared to bilayers the $\Delta OT_{\rm ENTH}$ values were about twice as high, indicating that PIP₂ is asymmetrically distributed between both leaflets of a silicon supported bilayer.

In Figure 5.2 a schematic illustration of the predicted protein occupancy on bi- and monolayers due to PIP₂ asymmetry is shown.

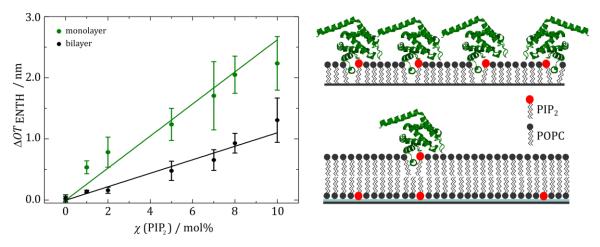


Figure 5.2: (A) Maximum change of the optical thickness upon ENTH binding to PIP₂ doped mono-(green) and bilayers (black) dependent on the PIP₂ content. (B) Schematic illustration of the PIP₂ distribution within the mono- and bilayers and the influence on protein occupancy.

As already mentioned a full protein coverage on the membrane surface is expected at $\Delta OT_{\rm ENTH}$ values of 2.0 – 2.8 nm.^{22,146} For the monolayers $\Delta OT_{\rm ENTH}$ was measured to be 2.2 ± 0.4 nm. Although the results (Figure 5.2) show no obvious saturation of the pro-

tein occupancy at 10 mol% PIP₂, a linear fit through the data indicates a slight overestimation of the expected $\Delta OT_{\rm ENTH}$ value at this PIP₂ concentration. Within the error margins a saturation at 8 or 10 mol% PIP₂ cannot be excluded, thus together with the achieved $\Delta OT_{\rm ENTH}$ value an almost full coverage of ENTH on the monolayer can be implied.

In both membrane systems, bi- and monolayer, a high reversibility of the protein binding was observed. For the bilayer $87 \pm 24\%$ and for the monolayer $75 \pm 14\%$ of the protein bound reversibly. A high reversibility of the ENTH binding ensures the dissociation from the clathrin-coated vesicles (CCV, cf. chapter 1.1) during CME, with the result that recycled proteins are available for the next circle of the process. Hence, the obtained high reversibility in the RIfS experiments are reasonable with regard to the process occurring in biological membranes.

The results show that the amount of bound protein is highly regulated by the PIP₂ concentration. A linear correlation between the protein occupancy and the PIP₂ content furthermore support the assumed 1:1 binding of ENTH to PIP₂ due to the formation of a PIP₂ binding pocket within the protein.

5.1.2 PS dependent binding affinity of ENTH to lipid bilayers

Besides PIP₂ also PS seems to have an influence on the ENTH-membrane interaction. 8,26 This effect was also investigated by means of RIfS. Since the addition of other lipids can change the membrane properties, the formation of POPC/POPS (80:20) and POPC/POPS/PIP₂ (75:20:5) bilayers was analyzed. Membrane thicknesses of $d = 4.2 \pm 0.1$ nm and $d = 4.3 \pm 0.2$ nm were measured for membranes with and without PIP₂, respectively. Experiments with POPC and POPC/PIP₂ (95:5) lipid compositions exhibited membrane thicknesses of $d = 4.4 \pm 0.2$ nm and $d = 4.5 \pm 0.3$ nm showing that there was no influence of the PS lipid head group on the membrane thickness. POPC/POPS (75:25) bilayer thicknesses of 3.81 ± 0.60 nm were previously measured using dual polarization interferometry. The values in this work are similar to those within the error margins indicating the successful formation of a bilayer.

Furthermore, fluorescence images showed a homogenous distribution of TxR in the POPC/POPS (80:20) and POPC/POPS/PIP₂ (75:20:5) bilayers. Using FRAP, diffusion coefficients of $D = 1.9 \pm 0.3 \, \mu m^2/s$ (without PS) and $D = 1.3 \pm 0.4 \, \mu m^2/s$ (with PS) were determined. Increasing the negative charge decreased the mobility, but considering the error margins the differences were negligible. The same trend is observed for the mobile fractions. Without PS the mobile fraction of TxR in the bilayers was $F_{\rm m} = 86 \pm 8\%$. Increasing the negative charge with PS decreased the mobile fraction to $F_{\rm m} = 77 \pm 4\%$.

Zhang *et al.* measured the diffusion coefficient of fluorescent probes with different net charges in biofilms. The negatively charged probes diffused more slowly, which was ascribed to the electrostatic repulsion of the probes with the biofilms. 144 Electrostatic repulsion of the negative charges PIP₂ and PS also can affect the mobility of TxR in the bilayers.

In conclusion, mobile and homogenous bilayers composed of POPC/POPS (80:20) and POPC/POPS/PIP₂ (75:20:5) were generated, enabling the investigation of the ENTH adsorption to these bilayers.

Influence of PS on ENTH binding

In order to investigate if ENTH binds specifically to PIP₂, ENTH was added to POPC/POPS (80:20) bilayers lacking PIP₂, showing no binding of ENTH to PS. Hom *et al.* performed liposome binding assays with POPC/POPE/PIP₂ (75:20:5) vesicles and showed specific binding of the ENTH domain to PIP₂.⁸ In the absence of PIP₂ the domain did not associate with the vesicles, even though PS was present. These results were also supported by co-sedimentation assays performed by Itoh *et al.* No binding of ENTH to negatively charged lipids, like phosphatidic acid (PA), was observed. Furthermore less or even no binding of ENTH occurred to other phosphatidylinositol species.²⁴

The results in this thesis showed that the dissociation constant of ENTH to PIP₂ in the presence of PS decreased from K_D =1.0 ± 0.2 μ M (POPC/PIP₂ (95:5)) to K_D = 0.42 ± 0.05 μ M (POPC/POPS/PIP₂ (75:20:5)). Moreover a higher protein occupancy was observed when PS was present ($\Delta OT_{\rm ENTH,max}$ = 2.5 ± 0.1 nm compared to

 1.6 ± 0.1 nm). This indicates that PIP₂ is necessary for the binding of ENTH to membranes, however PS decreased the dissociation constant of ENTH to PIP₂ doped bilayers and increased the protein occupancy on them.

Gleisner *et al.* measured K_D values by means of SPR, RIfS and spinning disc confocal laser microscopy (SDCLM). 36,79 RIfS experiments revealed a K_D of $0.8\pm0.8\,\mu\text{M}$ for POPC/PIP₂ (90:10). The K_D value is similar to the value obtained in this thesis via RIfS. With densitometry analysis of SDS-PAGE Hom et al. revealed that vesicles (diameter ≈ 1 μm) composed of POPC/POPE/POPS/PIP₂ (63:20:15:2) bound more ENTH than POPC/POPE/PIP₂ (78:20:2) vesicles indicating the influence of PS on the ENTH-PIP₂ interaction.⁸ With SPR experiments they furthermore measured K_D values of 20 ± 2.4 nM in presence of PS and 80 ± 11 nM when PS was absent, both at a pH of 7.4, mimicking the natural pH of mammalian cells. 150 In this thesis the K_D values were larger than the K_D values obtained by Hom et al.⁸ Recruitment of hydrophobic domains to curved membranes is preferred due to the higher binding site density and thus might explain the differences in the binding affinity. 151-153 However, Hom et al. observed that PS clearly enhanced the amount of bound ENTH to PIP2-containing membranes supporting the result obtained in this work.8 Addition of PS clearly increased the binding affinity of ENTH to PIP₂-containing membranes (Figure 4.14). Various other studies revealed values from 0.02 to 6 µM using fluorescence microsisothermal titration calorimetry (cf. Appendix, 10.3). 8,21,25,28,36,79,154,155 The dissociation constants determined in this thesis are comparable to the values in literature.

PS is localized at the cytosolic side of mammalian plasma membranes,³⁷ where ENTH binds to,¹⁵⁶ hence a contribution to the protein-membrane interaction is theoretically conceivable. Stahelin *et al.* suggested that non-specific interactions of ENTH with negative charges recruit the ENTH to the membrane which enables binding of ENTH to PIP₂.²⁵ Binding leads to conformational changes in the protein which also could explain that PS only has an effect on ENTH binding when PIP₂ is present (Figure 4.13). Otherwise, in the absence of PIP₂, no ENTH binding was observed. Thus, PS could enhance the recruitment and subsequently the binding affinity of ENTH. One explanation could be that due to ENTH-PIP₂ interaction a binding site for PS gets accessible or non-specific interactions with positive protein residues facilitate the binding. Since

in the presence of PS a higher protein occupancy on the bilayer was observed, the question arises if besides the helix insertion a crowding mechanism can contribute to a higher amount of bound protein in the presence of PS (Figure 5.3).

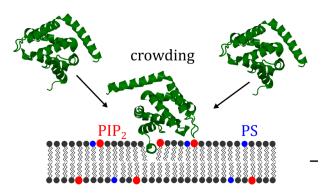


Figure 5.3: Schematic illustration of the ENTH crowding mechanism. The assembly of several ENTH monomers increases the protein occupancy on PS containing membranes in a PIP₂ dependent manner.

Kroppen showed with co-sedimentation assays that ENTH tends to form homo-oligomers in the presence of PS^{26} explaining the higher $\Delta OT_{ENTH, max}$ values (Figure 4.15). Protein-protein interactions can lead to higher surface occupancies or to stacked protein clusters. Both would increase the optical thickness. However, with RIfS experiments the formation of clusters cannot be observed.

PS induced ENTH cluster formation

Since PS showed an influence on the ENTH occupancy, AFM experiments were performed to obtain information about the surface topology after ENTH adsorption. Due to the comparability with the literature^{26,36,151} the lipid composition was changed to DOPC/DOPE/PIP₂/TxR (64.9:30:5:0.1) and DOPC/DOPE/DOPS/PIP₂/TxR (44.9:30:20:5:0.1). AFM experiments to investigate protein adsorption were previously performed with DOPC membranes,^{157,158} showing the applicability for this analysis.

Beforehand, FRAP experiments proved the mobility of these defect free membranes (Figure 4.19) with a $D = 2.5 \pm 0.6 \, \mu m^2/s$ (without PS) and $D = 1.7 \pm 0.4 \, \mu m^2/s$ (with PS), typical for bilayers composed of DOPC and DOPC/DOPS. The decreased D value is explainable with the increased negative charges of PS in combination with

PIP₂ interacting with the hydrophilic surface.^{136,143} Breakthrough experiments additionally verified the formation of bilayers with a membrane thickness of about 4 nm, indicating the successful formation of mobile bilayers.¹⁶²

Incubation with 1 μ M ENTH led to a significant decrease of the membrane mobility resulting in diffusion constants of $D=0.6\pm0.2~\mu\text{m}^2/\text{s}$ and $D=0.3\pm0.2~\mu\text{m}^2/\text{s}$ in absence and in presence of PS, respectively. Reduction of the lateral lipid diffusion is often observed when proteins bind to membranes. Due to binding of proteins, lipids diffuse slower in a protein-lipid complex. Hence decreased diffusion of ENTH-PIP2 complexes can affect the lateral diffusion of TxR in the bilayers.

In the atomic force micrographs (Figure 4.17), structures with an average height of 1.2 ± 0.2 nm were observed in the presence of PS indicating the adsorption of ENTH. The protein occupancy of $6 \pm 1\%$ indicates that only a minor amount of the protein is detected by AFM. Furthermore, no protein clusters were observed in in the absence of PS. At first sight this contradicts the results obtained with RIfS. Although the lipids differ, the head groups are the same and consequently the occupancy should not be affected that much. Previous studies also revealed similar binding behaviors of ENTH to membrane systems composed of POPC or DOPC. 36,154

The low occupancy moreover could be explained by the high reversibility of the ENTH binding detected via RIfS (Table 10.2). During imaging in contact mode lateral shear forces from the tip occurring to the sample can harm or detach the protein. The movement of the cantilever over the sample, hence, could shift the ENTH over the sample (Figure 5.4). This would also explain that no proteins were observed on the membrane in the absence of PS. Thus only PS induced ENTH clusters, which were rigid enough to be scanned in contact mode, were clearly imaged, causing an underestimation of the protein occupancy.

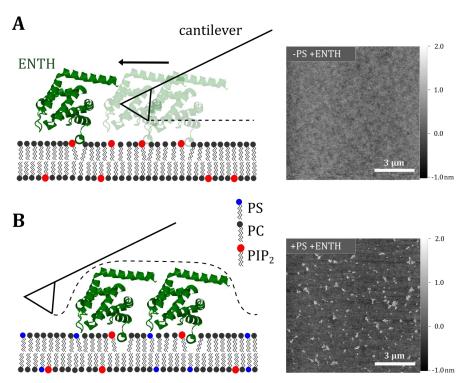


Figure 5.4: Cantilever movement in contact mode during imaging of (A) ENTH monomers and (B) ENTH clusters. Due to the mobile character of ENTH monomers the cantilever shifts the protein over the sample preventing the imaging of single monomers in contrast to immobile clusters.

EPR studies performed by Lai *et al.* indicated an immobilization of the protein due to close proximity of adjacent ENTH monomers, supporting the assumption that the aggregates in the micrographs are immobile and consequently only these were detected. They furthermore hypothesized that ENTH dimers or larger aggregates can occur dependent on the local membrane morphologies. Curvature facilitates the aggregation explaining the minor amount of visible clusters in the micrographs of planar solid supported bilayers.

Previous studies furthermore showed that ENTH tends to form homo-oligomers in the presence of PS.²⁶ Exchanging PS against phosphatidylinositol (PI) to mimic the negative charge¹⁶⁷ did not lead to oligomers. Hence, this supports the assumption that the cluster formation is induced by PS and not only by negative charges.

Although imaging in contact mode obtained low protein occupancies, other imaging modes were unsuitable for this system due to the appearance of inversion effects. ^{168,169} Imaging in tapping mode yielded micrographs, where dark structures in the bilayers were detected. Micrographs obtained in contact mode however revealed

that these dark structures were attributed to ENTH clusters and were a result of the inverted image contrast. In Figure 5.5 atomic force micrographs of a DOPC/DOPE/DOPS/PIP $_2$ /TxR (44.9:30:20:5:0.1) bilayer after incubation with 1 μ M ENTH obtained by QI-, tapping- and contact mode are illustrated. Rahe *et al.* observed inversion of the image contrast of protruded adsorbates on mica and TiO $_2$ in tapping mode depending on the amplitude of the cantilever oscillation and the damping level of the cantilever. Changed attraction of cantilever and adsorbate caused a phase shift of driving force and driven oscillation and consequently incorrect readout of the height. Hence in this work the different interactions between the cantilever with the bilayer and the protein could have led to the inverted contrast in the micrograph during imaging in tapping mode and consequently prevented the analysis of protein heights with this imaging mode.

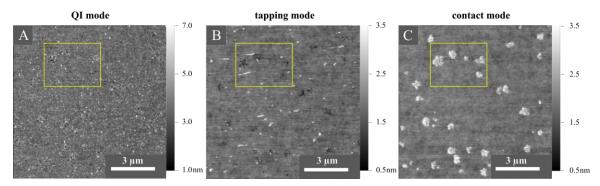


Figure 5.5: Atomic force micrographs of a DOPC/DOPE/DOPS/PIP $_2$ /TxR (44.9:30:20:5:0.1) bilayer after incubation with 1 μ M ENTH obtained by three different imaging modi: QI-, tapping and contact mode. The yellow squares highlight one part of the micrographs to underscore differences. Only when the measurements were performed in contact mode bright cluster structures were observed. Due to the inversion of the image contrast the ENTH clusters falsely were detected as membrane defects.

The atomic force micrographs obtained in contact mode (Figure 4.17 and Figure 5.5) clearly showed that PS induced the formation of ENTH cluster. The height of these clusters was determined to be 1.2 ± 0.2 nm which is smaller than reported in literature. Considering the crystal structure of ENTH at least a protein height of 2.0-2.8 nm was expected. Due to this, the bilayer surface was investigated by fluorescence microscopy to exclude ENTH binding to defects (Figure 4.19). Binding of ENTH to defects would lead to an underestimated protein height. To ensure that these defects

moreover did not appear during the membrane aging process, the bilayers were incubated in buffer for two hours and then imaged again showing no defects. In all cases, with or without PS, no visible defects were observed before and after ENTH addition or in the absence of ENTH after 2h. Hence the low protein height is not resulting from binding to defects and consequently the clusters were attributed to protein adsorption to PIP₂ doped membranes in the presence of PS.

The crystal structure of the bound ENTH is only known in the presence of inositol-1,4,5-trisphosphate (IP₃) (head group of PIP₂) and not when it is bound to the receptor lipid PIP₂. Although binding of ENTH to IP₃ alone results in the formation of the ENTH helix, it is not stable outside the membrane. When ENTH binds to a PIP₂ doped membrane, the helix is stabilized which would facilitate its insertion into the inner leaflet (Figure 5.6). 145

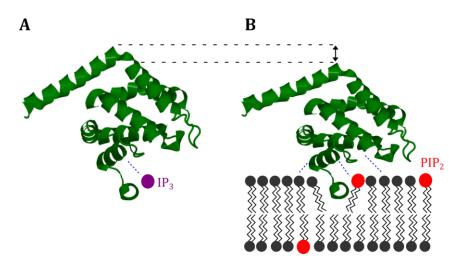


Figure 5.6: Schematic illustration of the different interactions of (A) IP_3 bound ENTH and ENTH bound to a PIP_2 doped membrane. Only when ENTH is bound to the membrane the helix is stabilized, which could lead to a slightly thinner protein layer on the bilayer surface compared to the assumed protein height from the crystal structure.^{21,145}

This would lead to an underestimation of the detected protein height compared to the values evaluated from the crystal structure. However, previous studies determined the physical thickness of ENTH via RIfS, yielding $d_{\rm ENTH}=1.7\pm0.2$ nm. ⁷⁹ Similar values were also obtained in this thesis by means of RIfS. Assuming a maximum of the protein height of 56% according to the scaled particle theory the actual height increases to 3.0 nm. ^{146,171} The differences in the protein heights are mainly explainable by the

used techniques. The protein heights of 1.2 ± 0.2 nm in this thesis were obtained by means of AFM in contact mode. During imaging in contact mode indentation of the protein with the cantilever can underestimate the protein height (Figure 5.7 A). ¹⁶⁶ Furthermore altered interactions of protein with the bilayer, e.g due to increase of the negative charge by PS, could affect the orientation and structure of the bound protein. This would result in a slightly different protein height compared to the protein height obtained from the crystal structure or with other techniques and lipid compositions (Figure 5.7 B and C). ^{22,79} Also other proteins, like synaptotagmin-1, showed different orientations with respect to PIP₂-containing membranes when PS is present. ⁵³ Since the presence of PS was shown to decrease the dissociation constant and increase the protein occupancy (cf. chapter 4.2.2) on the bilayer, altered interactions with the bilayer surface seems reasonable. Thus, the observed structures were attributed to ENTH clusters.

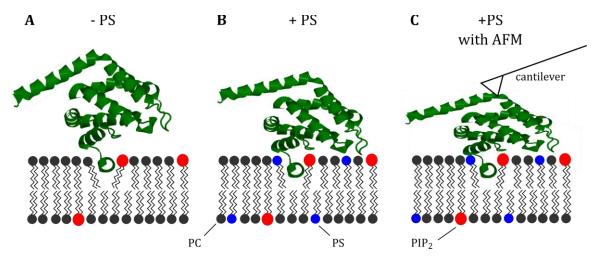


Figure 5.7: Predicted protein heights of ENTH for different conditions. (A) ENTH bound to PIP2-doped membranes. (B) ENTH bound membranes containing PIP2 and PS. (C) Indentation of the ENTH bound to PIP2/PS membranes by the cantilever. Due to lipid dependent protein-membrane interactions the protein height can change. Moreover imaging in AFM can indent the protein resulting in the detection of a decreased protein height.

The results clearly show that PS has a significant impact on the binding behavior of ENTH to PIP₂ doped membranes. First, the dissociation constant of ENTH to those bilayers is decreased and the protein occupancy on the surface increased. Moreover the AFM experiments revealed that PS induces the formation of ENTH clusters in the

presence of PIP₂. Hence, this shows that not only PIP₂ regulates the recruitment of ENTH to the membrane, but also PS highly affects the binding structure of ENTH to membranes.

5.1.2.1 The ENTH R114A mutant

There are different regions within the ENTH domain involved in protein-protein and protein-lipid interactions. Especially the amino acid R114 was found to interact with the membrane and to be involved in the formation of ENTH homo-oligomers. ^{26,110} Lai *et al.* performed coarse-grained molecular dynamics (CG-MD) simulation and EPR experiments with the ENTH domain. Their results indicated an interaction of the amino acid R114, located in the unstructured loop between helix 6 and helix 7, with the membrane. ¹¹⁰ This interaction further recruits ENTH and stabilizes the ENTH membrane docking geometry. Hence, the amino acid R114 can be involved in the cluster formation of ENTH.

In order to investigate the role of the amino acid R114 in the formation of clusters, as observed in the presence of PS (Figure 4.17), RIfS and AFM experiments were performed with the mutant R114A. The results showed a decreased dissociation coefficient of the mutant on POPC/POPS/PIP₂ bilayers and a higher protein occupancy on the membrane compared to bilayers lacking PS. Furthermore similar surface occupancies of ENTH R114A and ENTH wt were determined indicating that the mutant is still binding to PIP₂ doped membranes. However, the dissociation constant of the mutant significantly increased compared to the wt showing the contribution of the amino acid R114 in the binding process of ENTH. The obtained values for the ENTH wt and mutant are listed in Table 5.1.

Table 5.1: Summary of the fit results of the Langmuir adsorption isotherms for ENTH wt and ENTH R114A. The values are given as parameter ± SE.

lipid	K_{D}		ΔOT ENTH, max	
composition	ENTH wt	ENTH R114A	ENTH wt	ENTH R114A
POPC/ PIP ₂	$1.0 \pm 0.2 \mu\text{M}$	$2.9 \pm 0.8 \mu\text{M}$	1.6 ± 0.1 nm	1.5 ± 0.2 nm
(95:5)				
POPC/POPS/PIP ₂	$0.42 \pm 0.05 \mu\text{M}$	$1.0 \pm 0.3 \mu\text{M}$	2.5 ± 0.1 nm	2.7 ± 0.2 nm
(75:20:5)				

In general, the dissociation constants lie within the same range as the values found in literature (cf. Appendix, Table 10.3). Stahelin *et al.* mutated the ENTH domain at different positions. Mutations, like L6Q and M10Q strongly affected the binding affinity and the penetration ability into monolayers as these amino acids were localized in the α_0 helix.²⁵ The R114A mutant in turn exhibited a decreased dissociation constant compared to the wt, but showed a similar penetration ability into POPC/POPE/PIP2 (77:20:3) as the ENTH wt. This amino acid is not located at the PIP₂ binding site, but the close contact to the membrane facilitates the contribution in non-specific electrostatic interactions.^{23,25,110}

Lack of the positively charged protein residue (R114) or the negatively charged PS in the membrane led to a decreased binding affinity (Figure 4.21) and hence support the assumption that the amino acid R114 is involved in further recruitment of ENTH when PS was present. As the occupancy of ENTH R114A in presence of PS was similar to that of ENTH wt (cf. Table 5.1), it seems likely that still interactions of PS and the mutant occur in a PIP2 dependent manner, but ENTH clustering is prevented (Figure 4.22). Although the K_D of the mutant is higher as the value for the ENTH wt, the mobility of single ENTH R114A molecules seemed to be higher than the mobility of the ENTH wt clusters, hence the resolution of individual ENTH R114A monomers (and also individual ENTH wt monomers in the absence of PS, Figure 4.17) in contact mode was not possible.

Previous studies revealed the contribution of the amino acid R114 of ENTH in recruitment and membrane deformation. Kroppen showed that the mutation of this amino

acid did not change the molecular structure of ENTH and still was binding to the membrane in presence of PIP₂ and PS.²⁶ However, the ability to form tubular structures got lost proving the importance of this amino acid.²⁶ Moreover, co-sedimentation assays revealed the contribution of R114 in homo-oligomer formation. To investigate if the mutant also prevents cluster formation on solid-supported membranes, AFM experiments were performed. In all cases (DOPC/DOPE/PIP₂/TxR (64.9:30:5:0.1) and DOPC/DOPE/DOPS/PIP₂/TxR (44.9:30:20:5:0.1)) homogenous and planar bilayers with a roughness of 152±7 pm (without PS) and 197±47 pm (with PS) were observed after the incubation with ENTH R114A. Even in the presence of PS no clusters appeared (Figure 4.22) meaning that the lack of the arginine in position 114 prevented the formation of clusters. This also shows that the presence of PS and the amino acid R114 are required for the cluster formation which also can indicate an interaction of PS and R114.

Yoon *et al.* already revealed that interactions between two ENTH monomers can occur in the presence of PIP₂ and PS.²⁸ Although they did not investigate the impact of PS, they performed vesicle binding studies with GUVs composed of POPC/POPE/POPS/PIP₂/Rh-PE (46.5:30:20:0.5, Rh-PE = 1,2-dipalmitoyl-sn-glycerol-3-phosphoethanolamine-*N*-(lissamine rhodamine B sulfonyl)) showing the appearance of tubulation induced by ENTH crowding. This was attributed to the interaction of K23 and E42 of two ENTH monomers. When the amino acids were mutated (K23E and E42K) the ENTH lost the ability to induce tubulations. Their vesicle binding studies were supported by EPR studies which indicated an interaction of two α_0 helices (ENTH wt) in an antiparallel manner meaning a well-ordered clustering of the proteins. In combination with the results observed in this thesis, a contribution of PS and R114 in the accessibility of the amino acids K23 and E42 for protein-protein interactions (Figure 5.9) seems likely.

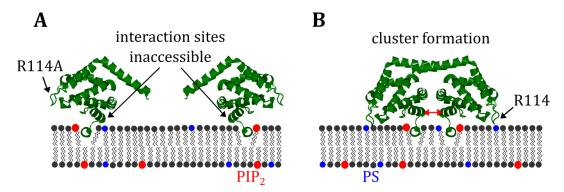


Figure 5.8: Orientation and interactions of two ENTH monomers in the (A) absence and (B) presence of PS. When PS is lacking less ENTH-membrane interactions occur. Due to the presence of PS interactions of ENTH with the membrane increase. Changed penetration depth and angles can lead to a facilitated interaction of K23 and E42 resulting in cluster formation.

Yoon *et al.* furthermore determined the insertion depth of the ENTH helix (usually 13 Å² from the phosphate group) as well as the penetration angle of the helix (11° to the membrane normal) in the presence of PS.²⁸ Lack of PS and/or R114 may change the orientation of the helix in the membrane which could lead to a different accessibility of the amino acids K23 and E42 for protein-protein interactions. Consequently less or no clusters were formed preventing imaging by AFM due to the mobile character of ENTH monomers (cf. chapter 5.1.2).

In summary, it can be stated that interactions of the amino acid R114 with the membrane or/and other ENTH monomers in the presence of PS are involved in the ENTH cluster formation.

5.2 Lipid dependent penetration ability of ENTH to lipid monolayers

Binding of ENTH to membranes illustrates the initial step of clathrin-mediated endocytosis (cf. chapter 1.1). As the ENTH-PIP₂ binding leads to the formation of a newly formed α_0 helix which can insert into the inner leaflet of the membrane, penetration experiments can help to understand the role of different lipids in this process. In the previous chapters, the importance of PIP₂ and PS on the membrane recruitment and cluster formation was discussed. To analyze if these lipids also have an effect on the

penetration ability of ENTH, the influence of PIP₂ and PS on lipid packing densities and the protein surface activity will be discussed in more detail.

5.2.1 Receptor lipid dependent penetration ability of ENTH

In order to investigate the lipid packing of monolayers composed of POPC/POPE (80:20) and POPC/POPE/PIP $_2$ (80-x:20:x, with x = 1, 5, 10 mol%), isotherms were recorded. The results show that the isotherms were shifted to smaller area per molecule values with increased PIP $_2$ content. It has to be considered that the area per molecule values (A_M) was averaged over all lipids in the ratio of their molecular fraction. The monolayers of POPC/POPE lacking PIP $_2$ showed a similar compression behavior as the monolayer with 5 mol% PIP $_2$. Although the fatty acid chains of PIP $_2$ molecules in sum exhibit lower degree of saturation than the fatty acid chains of POPC or POPE, ⁴⁴ different studies revealed similar area per molecule values for PIP $_2$ and POPC at 30 mN/m (about 70 Å $_2$). ^{172–175} Furthermore, it could be assumed that the negatively charged head group of PIP $_2$ should increase the lipid packing due to electrostatic repulsion. ¹⁷³

However, despite the negative charges, PIP₂ molecules can form hydrogen bonds which counteract the repulsion of the negative charge. ^{176,177} This would explain that the area per molecule values of POPC and PIP₂ are in the same range. Moreover, the increase in the PIP₂ content could increase the amount of hydrogen bond, leading to a higher lipid packing and a shift of the isotherms to smaller area per molecule values. Besides the increased amount of hydrogen bonds, the shift in the surface pressurearea per molecule diagram can also be an effect of decreased amounts of lipids in the monolayer as they submerge into the subphase. If more PIP₂ molecules submerge into the subphase, less molecules would be available for ENTH-PIP₂ interactions within the monolayers. Hence the critical surface pressure and consequently the penetration ability could be underestimated.

In general, no collapse of monolayers composed of pure POPC, POPE or PIP₂ until a surface pressure of 31 mN/m was observed in the literature.^{172-174,178,179} In this work, only the isotherms of POPC/POPE/PIP₂ (79:20:1) monolayers showed irregularities

or even collapsed between 27-31 mN/m. This means that the critical surface pressure seems to be in this region and that further compression is not possible due to monolayer collapse. Furthermore, the area per molecule molecule per area ($A_{\rm m}$) of the lipids in the POPC/POPE monolayers was slightly lower compared to the $A_{\rm m}$ in the 1 mol% PIP2 monolayer which would support the assumption that a certain amount of PIP2 is necessary for the formation of stable hydrogen bonds. Thus, collapse appeared earlier even with a looser lipid packing compared to monolayers with 5 and 10 mol% PIP2(Figure 4.23).

It is only possible to measure the penetration ability of ENTH when the monolayer stability is ensured. Especially the experiments with PIP2 required the monolayer stability to guarantee that enough PIP2 molecules were accessible for interactions with the protein. Johnston *et al.* assumed the stability of different phospholipid monolayers when the decay rate did not exceed 1 mN/m per minute. Penetration experiments were only performed with POPC/POPE monolayers as monolayers composed of DOPC and DOPC/DOPE were not stable (data not shown). The decay rate in this work was smaller compared to Johnston *et al.* (decay rate < 1.5 mN/m per hour) for all used lipid mixtures. Although submerging of lipids into the subphase cannot be excluded completely, the monolayers were stable over 1 h once 30 mN/m was reached. This enabled the performance of monolayer penetration experiments. (cf. sections 3.3.5 and 4.3.1).

After addition of the protein, the surface pressure (π) changed dependent on the initial value (π_0) . This increase is attributed to the insertion of the α_0 helix. The change in surface pressure $(\Delta\pi)$ was inversely proportional to π_0 , whereby the critical surface pressure (π_c) was determined as the x-intercept. As the increments of π often were irregular, no rate constants could be calculated. These irregularities were also found in penetration experiments with different proteins Addition were attributed to the uneven distribution of the ENTH in the subphase. In this setup, the mixing of the subphase after protein addition was only possible with a syringe, otherwise the monolayer was disturbed. Thus, distribution of the protein mainly took place by diffusion, which also increased the time to reach $\Delta\pi_{max}$.

Stahelin *et al.* e.g. used another Langmuir-Blodgett trough setup where a stirrer was used to achieve a homogenous distribution of the protein in the subphase.²⁵ They

monitored $\Delta\pi$ at a stirrer velocity of 60 rpm reaching a maximum after 20 min. In contrast to that, $\Delta\pi_{max}$ here was reached after 10-120 min which could be explained by the different techniques to distribute the protein in the subphase. Previous experiments determined a critical surface pressure of π_c = 21 mN/m for a POPC/POPE (80:20) monolayer.^{8,25} Also changes in the pH of the buffer did not influence the penetration ability essentially. In this thesis, ENTH also showed a poor penetration ability into POPC/POPE monolayers with π_c = 23.3 ± 0.8 mN/m. As cell membranes and large unilamellar vesicles exhibit similar lipid packing densities compared to Langmuir-Blodgett monolayers at surface pressures in a range of 30-35 mN/m,^{113,114,186} the protein is only able to penetrate natural bilayer if π_c exceeds 30 mN/m. This in turn means that the ENTH domain is not able to penetrate POPC/POPE bilayers which is reasonable as the ENTH helix is only formed in the presence PIP₂.²¹

For monolayers with PIP₂, π_c increased showing an influence of the receptor lipid on the surface activity of the protein. For monolayers with 1 mol% PIP₂, π_c slightly increased (π_c = 24.1 ± 0.8 mN/m), but within the error of the measurements it was not significantly different. At 10 mol% PIP₂, π_c changed to 28.7 ± 0.6 mN/m and at 5 mol% to 30.3 ± 0.4 mN/m. In both cases the π_c value was approximately around 30 mN/m, indicating a high penetration capability into these monolayers. Hence, this means that ENTH is a membrane-active protein domain capable of penetrating the membrane when a certain amount of PIP₂ is present.

 π_c changed from 24.1 ± 0.8 mN/m for 1 mol% PIP₂ to values aroun 30 mN/m for 5 and 10 mol% PIP₂ showing that the penetration ability of ENTH is a function of the PIP₂ content. Furthermore, there is the possibility that for 1 mol% PIP₂ molecules submerge into the subphase during compression, resulting in less ENTH binding and penetration into the monolayer. In some cases, the isotherms of the monolayers with 1 mol% PIP₂ showed a collapse at high surface pressures and this would support the assumption. A maximum of π_c was observed when ENTH was added to monolayers containing 5 mol% PIP₂. Since at higher PIP₂ concentrations more ENTH-PIP₂ interactions are expected due to previous results via RIfS (cf. chapter 4.1.3), the lipid packing could diminish a further insertion of the helix at 10 mol% PIP₂ (Figure 5.9 A). Moreover it is possible that PIP₂ molecules submerged into the subphase or that PIP₂

clusters could prevent the accessibility for ENTH binding (Figure 5.9 B and C). In all cases a reduced penetration of the ENTH helix would result.

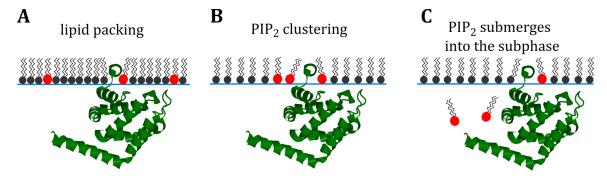


Figure 5.9: Schematic illustration of the ENTH penetration into monolayers composed of POPC/POPE/PIP₂ (70:20:10). Three penetration mechanisms, which would explain the limited helix penetration into those monolayers, are conceivable: (A) Insertion is prevented due to the dense lipid packing. (B) PIP₂ clusters diminish the accessibility for ENTH binding. (C) PIP₂ molecules submerge into the subphase preventing the ENTH-PIP₂ binding within the monolayer.

Stahelin *et al.* and Hom *et al.* determined π_c values of 32 mN/m and 31 mN/m for POPC/POPE/PIP₂ (77:20:3) and POPC/POPE/PIP₂ (78:20:2) monolayers at pH 7.4.8,25 Hom *et al.* moreover showed that at pH of 6 π_c decreased to 34.5 mN/m which was attributed to the facilitated ENTH binding to PIP₂ due to protonation of histidine residues (H73 and H43) in the protein. The experiments in this work were performed at pH 7.4 and therefore the measured π_c values are in good agreement with the literature, although the $\Delta\pi$ values at lower initial surface pressures were slightly smaller. Stahelin *et al.* used 3 mol% PIP₂ and Hom *et al.* 2 mol% PIP₂ instead of 5 mol% PIP₂ which could explain the small differences.8,25 Moreover, even small changes in temperature and pH can lead to changes in the surface pressure due to the high sensitivity and also can affect the measurements.187

Although the experiments with POPC/POPE/PIP₂ (79:20:1) monolayers showed a poor penetration ability of ENTH, RIfS measurements indeed showed specific binding of ENTH to supported POPC/PIP₂ (99:1) monolayers with $\Delta OT = 0.5 \pm 0.1$ nm (cf. Figure 4.9). One possibility for this difference could be the stabilization of the monolayer by the used substrate and consequently a higher amount of accessible PIP₂. On the other hand, the ENTH domain is recruited to POPC/POPE/PIP₂ monolayers (and POPC/PIP₂ monolayer for RIfS), but only a minor amount of the proteins penetrated

them. This would explain the low π_c of ENTH to 1 mol% PIP₂ monolayers. Furthermore, the protein penetration is only feasible up to an upper limit of the molecular packing. Many studies reported an optimum of the surface pressure in enzyme velocity and also the influence of the lipid packing density on protein penetration. Hence, this could support the assumption that the lipid packing in monolayers with 10 mol% prevents further helix insertion.

In conclusion, the results show that PIP_2 has a high impact on the penetration ability of ENTH as πc significantly increases for monolayers.

5.2.2 PS dependent penetration ability of ENTH

Lipids with charged head groups like PS can have essential impact on different protein-membrane interactions.^{192,193} One example is the protein domain ENTH, whose binding affinity is increased in presence of PS (cf. chapter 4.2). Moreover, it was shown that PS induces the formation of ENTH clusters (Figure 4.17).

In order to investigate if PS can additionally affect the penetration ability of ENTH, Langmuir-Blodgett through measurements were performed. First, the compression behavior of monolayers in presence and absence of PS was analyzed. No differences in the isotherms were observed when the subphase was changed from water to HEPES buffer, even though the ionic strength can have an influence on the lipid packing.^{173,194} Furthermore no collapse of the monolayer was detected until the surface pressure of 32 mN/m.

The monolayers lacking PS showed denser lipid packing (Figure 4.28) than the other monolayers with PS as they were shifted to lower area per molecule values.

Addition of PIP₂ to the PS monolayer even increased the area per molecule due to the incremented negative charges. These negative charges can repulse each other,¹⁹⁵ which would increase the area per molecule. Although PIP₂ molecules can form intermolecular hydrogen bonds among themselves and counteract the negative repulsion^{173,174} the additional negative charges of PS seemed to diminish this counteraction.

As all monolayers showed a high stability within one hour (percentage loss of π < 1.5 mN/m per hour), ¹¹⁵ penetration experiments of the ENTH domain into those monolayers were performed and analyzed. Again, an irregular increase of π after ENTH incubation was observed (cf. sections 4.3.1, 4.3.2 and 5.2.1) due to inhomogeneous protein distribution preventing the determination of rate constants. The increase of π indicated the penetration of the ENTH helix.¹⁸¹ At high area per molecule values, the lipid packing density is low whereby insertion or surface activity of the protein can be measured, resulting in an inverse proportionality of $\Delta \pi$ to π_0 . This behavior was also detected for the ENTH penetration experiment into PS-containing monolayers. The ENTH domain showed low penetration abilities into monolayers of POPC/POPE (80:20) with π_c = 23 mN/m and POPC/POPE/POPS (60:20:20) with π_c = 25 mN/m due to the lack of PIP₂. Addition of PS slightly increased the penetration capability, especially at low area values. Consequently the negative charge of PS influences the penetration ability of ENTH in absence of PIP2, even though the insertion into cell membranes seems unlikely $(\pi_c < 30 \text{ mN/m})$. 113,114,186 Compared to the POPC/POPE monolayers the POPC/POPE/POPS monolayers had a looser lipid packing (Figure 4.28) which would explain the slightly higher surface activity of the ENTH. RIfS experiments (Figure 4.13) showed that the negative charge of PS is not sufficient for ENTH-binding as the interaction is specific. ^{24,164} Hence, the low penetration ability is reasonable.

The experiments with POPC/POPE/POPS/PIP2 (55:20:20:5) monolayers revealed slightly higher π_c values with 32 mN/m compared to POPC/POPE/PIP2 (75:20:5) monolayers with π_c = 30 mN/m. Within the error margin, the differences are negligible indicating that the penetration ability of ENTH into POPC/POPE/POPS/PIP2 monolayers is not significantly different compared to POPC/POPE/PIP2 bilayers. In Figure 5.10, π_c is plotted against the area per molecule at π = 20 mN/m, showing that the lipid packing did not affected the ENTH helix penetration significantly. For a better overview, only values for 0 and 5 mol% PIP2 with and without PS were considered. The surface pressure of 20 mN/m was chosen to ensure that the monolayers did not collapse.

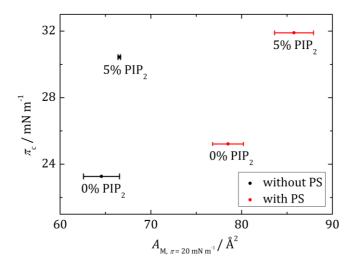


Figure 5.10: Illustration of π_c against the lipid packing at π = 20 mN/m for the lipid compositions POPC/POPE/PIP₂ (80-x:20:x, with x = 0 or 5 mol%) and POPC/POPE/POPS/PIP₂ (60-x:20:20:x, with x = 0 or 5 mol%) indicating that PS only slightly affects the penetration capability of ENTH although the lipid packing is decreased. Subphase: HEPES buffer.

The area per molecule increased about 15-20 $\mbox{Å}^2$ when PS was present. Due to repulsive interactions, the area per molecule incremented resulting in a larger area for the lipids.

In different studies using vesicles, it was shown that the outer leaflet of the vesicles has a looser lipid packing at higher membrane curvatures. ^{196–198} Moreover, the recruitment of hydrophobic helices to curved structures is preferred, ^{153,199} thus penetration into a monolayer with a looser lipid packing was expected. Although the penetration is preferred when the lipid packing is low, the results show that the penetration is mainly a function of the PIP₂ concentration at a pH of 7.4. Hom *et al.* revealed that the penetration capability of ENTH increased with decreased pH due to protonated histidine residues interacting with the monolayer. ⁸ Therefore, changing the pH value might also increase penetration in case of PS. Furthermore the preferred helix insertion into membranes with looser lipid packing may facilitate penetration in curved structures, but not in flat Langmuir-Blodgett monolayers.

In contrast to this, binding studies by means of RIfS, PIP₂ and PS containing bilayers showed a higher protein occupancy than the bilayers without PS (Figure 4.15). This indicates that ENTH adsorption to PIP₂ doped membranes is facilitated, however the penetration ability is not affected. Hence, formation of clusters induced by PS could

be a result of protein-protein interactions, where only a few ENTH monomers penetrate the monolayer (or the outer monolayer in supported bilayers) and the other monomers are interacting with the bound ENTH monomer without penetrating the membrane (Figure 5.11 A). As this cluster formation is not observed in the absence of PS (Figure 4.17), binding to PS could establish new binding sites for protein-protein interactions.

Yoon *et al.* assumed an orientation of two ENTH helices in an antiparallel manner and additionally interactions of the amino acids K23 und E42 in the presence of PS.²⁸ Thus two binding structures of ENTH are conceivable. In Figure 5.11 the different feasible membrane-protein and protein-protein interactions in the presence of PS are illustrated.

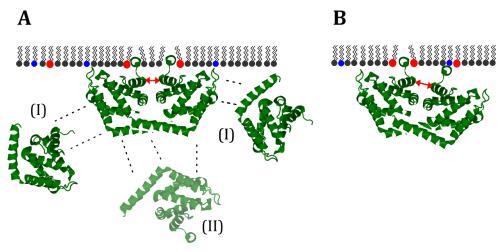


Figure 5.11: Schematic illustration of predicted ENTH-ENTH interactions resulting in protein clusters. Since penetration of ENTH into the monolayers in the presence of PS is not significantly increased although cluster formation was observed by means of AFM, interaction of the ENTH with PS could result in new interaction sites within the ENTH monomers for protein-protein interactions (A). In cases of (I) and (III) the result would be a higher protein occupancy on the surface. Case (II) seems less likely as the protein height observed via AFM showed lower monomer heights. (B) Furthermore, interactions of the helices in an antiparallel manner still can occur, when only one helix inserts into the monolayer. This would indicate a higher occupancy, but no further penetration.

The results so far show that the lipids PIP_2 and PS have an impact on the ENTH-membrane interaction. Increase of PIP_2 and addition of PS increased the amount of bound ENTH on bi- and monolayers. PIP_2 specifically binds the ENTH and induces the formation of a new α_0 helix which can insert into the membrane. PS additionally induced the formation of protein clusters. This in turn shows that the lipid composition

contributes to the regulation of protein-dependent membrane deformation during CME.

5.3 Adhered GUV as a membrane model to investigate proteinmembrane interactions

Adhesion of membranes and regulation of the membrane tension are substantial factors in biological processes, e.g. in exo- and endocytosis.^{200,201} In the previous chapters the impact of PIP₂ and PS on the binding structure of ENTH in protein-membrane and protein-protein interactions was discussed in detail. In order to investigate if the ENTH can moreover affect the mechanical properties of membranes dependent on the lipid composition, adhered GUVs were utilized as membrane models.

The interactions of ENTH with PIP₂ as a function of lateral membrane tension were already analyzed previously,⁷⁹ thus in the next chapter the influence of ENTH on the membrane tension in the presence of PS was investigated.

5.3.1 Adhesion of biotinylated GUVs on PEG-functionalized surfaces

The formation of adhered GUVs requires a suitable surface to generate a contact area between the GUV and the substrate. Immobilization of GUVs via biotin-avidin interactions was already realized previously. Gleisner *et al.* incubated avidin on hydrophilized glass substrates.³⁶ Addition of biotinylated GUVs composed of either DOPC/DOPE/cap-biotin-DOPE/PIP₂/Atto488-DPPE (66:30:2:1:1) or DOPC/DOPE/cap-biotin-DOPE/Atto488-DPPE (67:30:2:1) resulted in immobile vesicles with a contact area to the surface. The lateral membrane tensions of these adhered GUVs ranged from 0.01 to 1.02 mN/m.³⁶ Regulation of the membrane tension was realized by varying the concentration of MgCl₂ in the surrounding buffer and was attributed to stronger interactions of PIP₂ and the Mg²⁺ ions. Thereby, smaller membrane tensions were determined for GUVs lacking PIP₂ compared to GUVs with PIP₂.

In this thesis, the impact of PS on the membrane tension of GUVs was investigated using similar conditions as Gleisner *et al.*³⁶ Since negative charges of PS and PIP₂ are known to cluster in the presence of MgCl₂, regulation of the membrane tension with the MgCl₂ concentration was excluded.^{125,126} Moreover, it was shown that NeutrAvidin exhibits less non-specific interactions than avidin in immunoassays while maintaining a high binding affinity to biotin ($\sim 1.3 \cdot 10^{-9} \, \mu M$).^{127,202} Thus, it is assumed to be promising for biosensing applications. To reduce unspecific interactions of the negative charges of PS/PIP₂ and avidin to a minimum, avidin was replaced by the deglycosylated protein NeutrAvidin.

However, incubation of NeutrAvidin on hydrophilized glass substrates resulted in an inhomogenous distribution of the protein on the surface (Figure 4.31) and immobilization of the vesicles via NeutrAvidin-biotin binding was not possible as the protein was removed from the surface during rinsing. Avidin has its isoelectric point at pH 10.5 and it is positively charged at a pH of 7.4.²⁰³ This promotes the electrostatic interactions of avidin with the surface. Compared to that, NeutrAvidin exhibits an isoeletric point of 6.3, which leads to less non-specific interaction with the substrate surface, while the high binding affinity to biotin is still existing^{60,127} and would explain the inhomogenous coverage of NeutrAvidin on the hydrophilized glass surface.

On this account, the strategy for NeutrAvidin immobilization had to be changed. A method to immobilize vesicles bases on the work of Piehler *et al.* was used. ¹²⁹ Silanization of glass substrates with (3-Glycidyloxypropyl)trimethoxysilane (GOPTS) resulted in a surface with epoxide moieties. Depending on the added reagent, the functionalization on the substrate surface was regulated (Figure 3.3). J. Thiart and Schwamborn *et al.* utilized this technique and coupled amino-terminated Poly(ethylen glycol) (PEG) groups to the surface. ^{128,204} This solvent-free coupling with amino groups of bifunctional PEGs enabled immobilization of vesicles on these surfaces. Using biotin-PEG-NH₂ (b-PEG) ensured the affinity for NeutrAvidin and methoxy-PEG-NH₂ (m-PEG) was used to obtain a lateral distance between the biotin groups to be accessible for binding (Figure 5.12).



Figure 5.12: PEG-based surface functionalization. After silanization and PEGylation of glass substrates, the biotin-PEG (b-PEG)/methoxy-PEG (m-PEG) covered surface was incubated with NeutrAvidin, enabling the immobilization of biotinylated GUVs on these surfaces.

In contrast to the studies of J.Thiart and Schwamborn *et al.*, b-PEG and m-PEG species with less PEG groups were used (M(m-PEG) = 350 g/mol and M(b-PEG) = 600 g/mol instead of M(m-PEG) = 2000 g/mol and M(b-PEG) = 3000 g/mol). PEG is known to reduce the adhesion of proteins or other molecules and is used to protect surfaces. 205,206 Thereby, the PEG density on the surface and the chain length can influence the adsorption of molecules. 207 Since an ordered structure of PEGs on the surface cannot be guaranteed, biotin molecules could be less accessible. Therefore, the amount of PEG groups was reduced to diminish the repelling character of the PEG surface.

Incubation of these m-PEG/b-PEG surfaces with labeled NeutrAvidin yielded an even distribution in the fluorescence intensity of the labeled protein on the surface (Figure 4.31) indicating a homogenous coverage of NeutrAvidin on the substrates. As the b-PEGs were covalently bound to the silane on the glass surface, the NeutrAvidin also was immobilized due to the high affinity to biotin and was not removed by rinsing with buffer.

Since a homogenous NeutrAvidin occupancy on the PEG surfaces was realized, the system was assumed to be suitable to immobilize biotinylated GUVs via further NeutrAvidin-biotin interactions. Therefore, GUVs composed of DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1) were added to the system resulting in adhered GUVs on the b-PEG-NeutrAvidin surface(Figure 5.13).

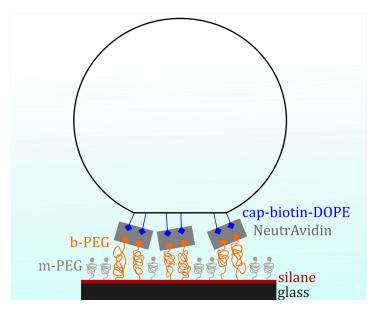


Figure 5.13: Schematic illustration of the PEG-based surface functionalization. After silanization and PEGylation of glass substrates, the biotin-PEG (b-PEG)/methoxy-PEG (m-PEG) covered surface was incubated with NeutrAvidin enabling the immobilization of biotinylated GUVs on these surfaces.

These GUVs were stable over 2 h, which was proved by determination of the $r_{\rm i}/r_{\rm ad}$ values of the adhered GUVs directly after immobilization (15 min) and after 2h (Figure 4.33). The $r_{\rm i}/r_{\rm ad}$ values did not significantly change within the deviation (from 0.16 ± 0.04 to 0.17 ± 0.04) indicating that the GUVs were stable over time. Most likely the stability was maintained due to iso-osmolar conditions using a humidity chamber. The membrane tensions of these GUVs were measured to be $(1.03 \pm 0.04) \cdot 10^{-2}$ mN/m (after 15 min) and $(1.03 \pm 0.03) \cdot 10^{-2}$ mN/m (after 2 h). The values in this thesis are comparable to cellular tensions which are in the range of 0.002-0.3 mN/m depending on the cell type. $^{201,208-210}$ Compared to the membrane tensions of GUVs adhered to avidin coated glass surfaces measured by Gleisner *et al.*, the values of this work lay in the low membrane tension range.

To mimic membrane tensions of different cell types, the lateral membrane tension of the adhered GUVs should be regulated. Due to this, the amount of biotin on the surface was varied by changing the ratio of m-PEG and b-PEG to see if the amount of accessible biotin molecules could be increased. NeutrAvidin has 4 binding sites for biotin, 2 on each side (cf. Figure 5.13),⁵⁸ thus a higher concentration of biotin on the surface could increase the coverage of NeutrAvidin on the surface. However the membrane tension of the GUVs was not affected by variations in the m-PEG/b-PEG ratio (Figure

4.34). Since NeutrAvidin has a high binding affinity to biotin, the surface could be completely covered at a ratio of 1:1.^{127,202} Thus further raise of the b-PEG concentration would not affect the protein occupancy and the adhesion of biotynilated GUVs. Another possibility to regulate the membrane tension could be achieved by changing the lipid composition. Nevertheless, it was not tried to vary the membrane tension as a function of the lipid composition because the search for an adequate lipid composition, resulting in stable GUVs, was challenging. However, the increase of biotin-coupled lipids may facilitate binding to the NeutrAvidin surfaces and consequently increase the membrane tension.

Moreover, the variation of the negative charges (PIP₂ and PS) in the lipids could have an impact on the membrane tensions. Raucher *et al.* already showed with experiments on NIH-3T3 cells that PIP₂ can control the adhesion.²¹¹ Raise of the PS concentrations in combination with this surface functionalization though, was not possible as these vesicles were unstable or too small to determine the membrane tension based on the geometry. Although GUVs with higher proportions of negative charges (~25%) were already generated,^{212,213} incorporation of negative charges is limited due to repulsive interactions and cluster formations.^{214,215} As higher PS concentrations were already realized in free vesicles, it is more likely that the PEG functionalization and the adhesion (higher membrane tensions) also had an impact on the GUV stability when a higher PS content was incorporated in the vesicles.

Thus, before the lipid composition was further adjusted, the impact of ENTH on the membrane tension of GUVs composed of DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP₂/Atto488-DPPE (76:10:10:2:1:1) was investigated in detail.

5.3.2 Effect of the ENTH binding to PS-containing adhered GUVs

Once stable adhered GUVs were generated and the membrane tension was determined, the ENTH binding activity on these membranes was analyzed. After incubation of GUVs composed of DOPC/DOPS/Cholesterol/cap-biotin-DOPE/PIP $_2$ /Atto488-DPPE (76:10:10:2:1:1) with 1 μ M ENTH for 2h, the solution contained less GUVs and

a lot of lipid material outside the vesicles was observed. The amount of vesicles before and after ENTH addition was not counted, but the difference was obvious. During imaging of the remaining GUVs, the vesicles ruptured. This only happened when ENTH was present (Figure 4.37). In the absence of ENTH, no rupturing event occurred within 2 h (Figure 4.33) indicating that the rupture of the vesicles is affected by the protein adsorption to the PIP₂ doped GUVs. Reduction of the ENTH (0.3 μ M) concentration still resulted in GUV rupturing (Figure 4.38) showing that even significantly lower concentrations are sufficient to induce GUV bursting.

In clathrin-mediated endocytosis (CME), binding of the ENTH domain to PIP₂ results in the formation of a new amphipathic helix which inserts into the inner leaflet of the membrane.²¹ The insertion of the helix disturbs the membrane and consequently can influence the lipid packing in the membrane.⁷ In this work an analogous behavior for the adhered GUVs might cause an insertion into the outer leaflet.

Therefore, different reasons for the presence of less GUVs and the vesicle rupture are conceivable. First of all, the helix insertion can cause membrane defects due to the disturbance of the lipid packing. ^{216,217} This could destabilize the vesicles, leading to the breakdown of the osmolarity gradient and consequently rupture of the vesicles. Different studies revealed that insertion of an amphipathic helix is preferred to small hydrophobic lipid bilayer packing defects. ^{218–220} This is also consistent with the fact that ENTH favors binding to curved membranes as they exhibit more packing defects. ²²¹ Preferred insertion of the ENTH helix into those membranes can facilitate membrane deformation or induction of membrane defects. Although the bilayer structure of GUVs seem to be flat in the lipid scale, they exhibit a higher curvature compared to flat Langmuir-Blodgett monolayers. Thus, a slightly higher insertion rate of the ENTH helix can be assumed. In the adhered GUVs the PS could additionally increase the insertion of the ENTH helix, even though it was not shown in Langmuir-Blodgett monolayers.

Furthermore, the helix insertion reduces the bending rigidity of the membrane. A reduced bending stiffness lowers the energy which is required for membrane deformation.⁷⁹ Due to this, the formation of tubules or even vesicles out of the GUV is facilitated (Figure 5.14) which leads to the destabilization of the vesicles.

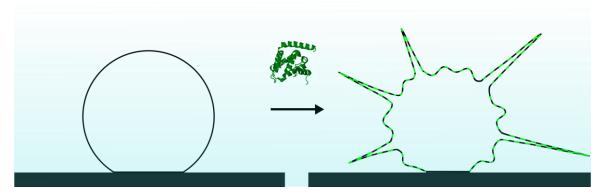


Figure 5.14: Binding of ENTH (green) to adhered GUVs resulting in the formation of tubules, which furthermore destabilizes the adhered GUV.

During the formation of these structures the size of the vesicle decreases until all lipid material is dissipated. This would explain why usually less GUVs were found in the solution after ENTH addition. Tubulation of vesicles due to ENTH adsorption was already reported in the literature. Yoon *et al.* showed that incubation of free vesicles composed of POPC/POPE/POPS/PIP₂/Rh-PE (46.5:30:20:3:0.5) with 0.5 μ M ENTH resulted in the appearance of tubular structures. ²⁸ Gleisner *et al.* showed with studies on adhered GUVs that at low membrane tensions (0.08 mN/m) tubular structures were observed after the ENTH addition (1 μ M). In contrast to that adhered GUVs flattened at high membrane tensions when ENTH was present. The adhered GUVs in this thesis exhibited even lower membrane tensions (0.01 mN/m), indicating that tubulation seems likely.

However, some studies revealed that tubulation only occurs, when the protein coverage of the surface is above 20%. 30 At high ENTH concentrations and 1 mol% PIP₂ concentration, a protein coverage of maximum 23% should be reached considering the footprint of ENTH (16 nm^2) 22 and the surface demand of the receptor lipid (0.7 nm^2). 44 At a concentration of 1 μ M this decreases to 16% when the K_D of 0.42 μ M is noted, which should not induce tubulations. Nevertheless, tubulation was already observed at lower protein coverage (11.4%) even in the absence of PS⁷⁹

Addition of PS might facilitate the tubulation. It was shown that PS has a significant impact on the binding as a higher binding affinity and a higher protein occupancy on the bilayer surface were measured in the presence of PS (cf. chapter 4.2.2). This was

attributed to the cluster formation induced by PS. Protein crowding was shown to destabilize membrane shapes.³⁵ Cluster formation and the insertion of the ENTH helix in one region of the GUVs would lead to instabilities, hence the rupturing seems reasonable.

Furthermore, heterogeneities of the lipid distribution within the GUVs could lead to a higher amount of PIP₂ and PS in some vesicles. Different studies revealed that the heterogeneities of the lipid compositions in one GUV preparation via electroformation can vary from 2 to 13%.^{222,223} If more PIP₂ and PS are present in some GUVs compared to others (excluding PIP₂ clustering), this would lead to a higher amount of bound ENTH and would explain the reduced amount of GUVs as they burst earlier due to defects caused by the helix insertion of ENTH.

Krishnan *et al.* showed that the interactions of PIP₂ with the actin-binding protein profilin led to the destabilization of GUVs.²²⁴ From 8-10 mol% PIP₂ they recognized the start of PIP₂ cluster formation. Further increase of the concentration destabilized the GUVs due to the local membrane deformation by the clusters. If some vesicles exhibit high local accumulations of PIP₂, it can lead to destabilization.

Considering that some GUVs were imaged in the first seconds and only ruptured during the time-series, perhaps also the laser irradiation had an influence on the GUV stability. Brückner *et al.* induced changes in the GUV shape with UV light. Other experiments varied the area/volume ratio by temperature dependent vesicle shape changes. Although the energy of UV (10-400 nm) and visible light (here \sim 600 nm) differ and the temperature changes by the laser should be negligible, an impact of the laser illumination cannot be excluded as GUVs rupture during imaging.

As the amount of GUVs during incubation with ENTH seemed to decrease, a combination of different effects is assumed as its cause. First of all, the helix insertion disturbs the bilayer and leads to lipid packing defects which destabilize the GUV. Secondly, the presence of PS induces protein crowding, facilitating membrane deformation and increasing membrane defects. Finally, the laser illumination leads to the rupture of the remaining destabilized vesicles.

In conclusion, it was shown successfully that the analysis of adhered GUVs on b-PEG-NeutrAvidin functionalized surfaces is possible. However, the system is only suitable to a limited extent to investigate the impact of ENTH on the mechanical properties of these adhered GUVs in the presence of PS. Reducing the PS concentration or/and repelling character of the PEG surface (e.g. further decrease of the amount of PEG groups) might improve the system.

6 CONCLUSION

Membrane remodeling processes which occur during endocytosis, are regulated by the complex interactions between proteins and lipid membranes. The most important pathway for endocytosis is the clathrin-mediated endocytosis (CME), where adaptor proteins like epsin induce a membrane invagination by specific binding to the receptor lipid phosphatidylinositol-(4,5)-bisphosphate (PIP₂). In this thesis the lipid dependent binding of the epsin *N*-terminal homology (ENTH) domain to artificial membrane models was analyzed.

The binding studies of ENTH to supported lipid bilayer and monolayers doped with PIP $_2$ revealed an asymmetric distribution of the specific receptor lipid within both leaflets of a bilayer. Increasing the PIP $_2$ concentration resulted in an increased optical thickness and consequently a higher occupancy on the membrane surface. Only in case of monolayers nearly a complete coverage was achieved with an optical thickness of 2.2 ± 0.4 nm.

Addition of lipids with the negatively charged head group phosphatidylserine (PS) increased the binding affinity of ENTH to PIP $_2$ containing membranes. Since the ENTH did not bind to POPC/POPS bilayers in the absence of PIP $_2$, the specific binding to PIP $_2$ and not to the negative charge by non-specific electrostatic interactions was proved. The dissociation constant of ENTH binding to a POPC/POPS/PIP $_2$ bilayer was about twice as high as to a POPC/PIP $_2$ bilayer. Moreover, the maximum in optical thickness upon protein binding increased from 1.6 ± 0.1 nm (without PS) to 2.5 ± 0.1 nm (with PS) at the same PIP $_2$ concentration, showing a relevant effect of PS on the binding behavior of ENTH to those membranes. Analysis of the surface topography before and after ENTH adsorption indicated that this increased occupancy resulted from cluster formation of ENTH monomers when PS was present. In the absence of PS no clusters or single molecules were detected on the bilayer surface due to the high mobility of the single ENTH monomers.

Monolayer penetration experiments were performed to investigate the penetration ability of ENTH into monolayers as a function of the surface pressure dependending on the lipid composition. The critical surface pressure is the maximum value, where

no insertion of proteins into the monolayer is possible anymore. Increasing the PIP₂ content first increased the penetration capability and reached a maximum at 5 mol% PIP₂ with π_c = 30.3 ±0.4 mN/m.

The lipid packing was then changed by the addition of PS, only slightly increasing the critical surface pressure to π_c = 32.0 ± 0.9 mN/m, although a higher binding affinity and protein coverage on PIP₂/PS doped membranes was observed. This indicates that PS is inducing protein clusters, which only partly penetrate the membrane.

Adhered GUVs as membrane models enabled to measure ENTH activity as a function of membrane tension. The membrane tension of GUVs adhered to NeutrAvidin functionalized substrates was determined to be 0.01 mN/m. After incubation with ENTH, these GUVs ruptured induced by the insertion of the ENTH helix and consequently appearance of lipid packing defects in the membrane.

In conclusion, it was shown that ENTH binding to membranes is highly regulated by the lipid composition. Furthermore, the results revealed that the membrane deformation in CME induced by ENTH is an interplay of the ENTH helix insertion and the crowding mechanism induced by PIP₂ and PS.

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10 APPENDIX

Devices

Atomic force microscopy

Nanowizard 4 JPK Instruments AG, Berlin, GER

Fluorescence microsopy

spinning disc confocal

LSM 880 Examiner Carl Zeiss, Jena, GER
Objektiv WPlan APOChromat 63x/1.0 Carl Zeiss, Jena, GER

custom stand based on ix73 Olympus, Tokyo, JAP

spinning disc Yokogawa CSU-X Rota Yokogawa, Wehr, GER

camera iXON 897Ultra Andor Technology Ltd., Belfast, UK

AOTF: TF525-250-6-3-GH18A Gooch&Housego PLC, Illminster, UK

piezo P-721-CDQ Physik Instrumente, Karlsruhe, GER

filter wheel Rotr Andor Technology Ltd

filter LP496 ET AHF Analysentechnik, Tübingen,

GER

custom build

filter ET525/50 AHF Analysentechnik, Tübingen,

GER

objective CFI Plan 100XW Nikon, Düsseldorf, GER

100×/1.1

Langmuir-Blodgett trough

RK 1 Standard Langmuir Trough Riegler & Kirstein GmbH, Potsdam,

GER

Thermostat VWR VWR International, Darmstadt

Reflectometric Interference Spectroscopy (RIfS)

NanoClalc-2000, SD2000 spectrometer Ocean Optics, Dunedin, FL, USA
Flame Spectrometer Ocean Optics, Dunedin, FL, USA
Halogen light source Ocean Optics, Dunedin, FL, USA

Ismatec 795C peristaltic pump IDEX Health & Science, Wertheim,

GER

UV Vis Spectroscopy

NanoDrop 2000c Thermo Scientific, Wilmington, USA

Electroformation

Frequency generator Aglient 333220A Aglient Technology, St. Clara, USA

Other devices

Micro scale Satorius CP225D Satorius, Göttingen, GER
Magnetic mixer MR 3001 K Heidolph, Schwabach, GER

Osmometer

pH-meter Calimatic 766 Knick, Berlin, GER

Plasma cleaner Zepto Diener Electronic, Ebbhausen, GER

Thermomixer compact Eppendorf, Hamburg, GER

Tip sonifier Sonoplus HD2070 Bandelin, Berlin, GER

Ultra sonic bath Sonorex RK 255 H Bandelin, Berlin

Ultrapure water system Millipore, Billerica, MA, USA

Vacuum drying oven VD 23 Binder, Tuttlingen, GER

Vortexer Heidolph, Schwabach; GER

NanoDrop 2000c ThermoFisher Scientific, Waltham,

MA, USA

Software

ImageJ http://rsbweb.nih.gov/ij

Origin Pro 8.5G OriginLab Cooperation, Northhamp-

ton, USA

ZEN black 2.3 SP1 Carl Zeiss, Jena, GER

MATLAB 2017b MathWorks, Natrick, MA, USA

Gwyddion 2.45 http://gwyddion.net

SpectraSuite Ocean Optics Germany GmBH, Ost-

fildern, GER

JPK Data Processing Instruments AG, Berlin, GER

RuK Trough Control Riegler & Kirstein GmbH, Potsdam,

GER

Protein sequence

The sequence of ENTH was obtained from the Protein Data Bank.²³

- 1 HNYSEAEIKV REATSNDPWG PSSSLMSEIA DLTYNVVAFS EIMSMIWKRL
- 51 NDHGKNWRHV YKAMTLMEYL IKTGSERVSQ QCKENMYAVQ TLKDFQYVDR
- 101 DGKDQGVNVR EKAKQLVALL RDEDRLREER AHALKTKEKL AQTA

Optical thickness changes upon ENTH addition to POPC/PIP $_2$ bi- and monolayers

Table 10.1 Averaged maximum changes in optical thickness after ENTH addition to POPC/PIP₂ bi- and monolayers with different PIP₂ concentrations.

Lipid	bilayer	Number of	monolayer	Number of
Composition	$\Delta OT_{\rm ENTH}$ / nm	measurements	$\Delta OT_{\rm ENTH}$ / nm	measurements
POPC/PIP ₂		n		n
100:0	0.03±0.01	5	0.06±0.04	3
99:1	0.1±0.1	4	0.5±0.2	4
98:2	0.2±0.1	5	0.8±0.3	5
95:5	0.5±0.2	5	1.2±0.3	4
93:7	0.7±0.2	4	1.7±0.6	3
92:8	0.9±0.2	5	2.0±0.4	3
90:10	1.3±0.4	8	2.2±0.5	6

Reversibility of the ENTH binding to POPC/PIP₂ bi- and monolayers

Table 10.2: Reversibly of the ENTH binding to $POPC/PIP_2$ bi- and monolayer in %. The PIP_2 content was varied from 1-10 mol%. The errors are given as standard deviations.

	Reversibility of the ENTH binding/%		
Lipid composition	Lipid bilayer	Lipid monolayer	
POPC/PIP ₂ (99:1)	70 ±40 (n = 4)	83 ±15 (n = 4)	
POPC/PIP ₂ (98:2)	88 ±24 (n = 5)	73 ± 15 (n = 4)	
POPC/PIP ₂ (95:5)	85 ±15 (n = 6)	80 ± 11 (n = 4)	
POPC/PIP ₂ (93:7)	99 ±21 (n = 4)	74 ± 23 (n = 3)	
POPC/PIP ₂ (92:8)	89 ±14 (n = 7)	$70 \pm 6 \ (n = 3)$	
POPC/PIP ₂ (90:10)	85 ±6 (n = 4)	68 ± 4 (n = 5)	

Dissociation constants of the ENTH binding to PIP₂ containing membranes

Table 10.3: Summary of all dissociation constants K_D of ENTH to PIP₂ doped membranes found in literature.

Binding assay	Lipid composition	K_{D}
Liposome assay ¹⁵⁵	PC/PE/PIP ₂ (70:20:10)	618.4 nM ± 0.009%
SPR ²⁵	POPC/POPE/PIP ₂ (77:20:3)	23 ± 7 nM
	POPC/POPE/PIP ₂ (79.5:20:0.5)	250 ± 60 nM
SPR ⁸	POPC/POPE/PIP ₂ (78:20:2)	80 ± 11 nM
	POPC/POPE/POPS/PIP ₂ (63:20:15:2)	20 ± 2.4 nM
SPR ²⁸	POPC/POPS/PIP ₂ (77:20:3)	49 ± 8 nM
SPR ⁷⁹	POPC/PIP ₂ (90:10)	0.5 ± 0.1 μM
	POPC/POPE/PIP ₂ (60:30:10)	0.69 ± 0.08 μM
RIfS ^{36,79}	POPC/PIP ₂ (90:10)	0.8 ± 0.2 μM
Calorimetric titra-	PC/PE/Cholesterol/diC8PIP2	0.85 μΜ
tion ²¹	(40:40:10:10)	
Fluorescence mi-	DOPC/DOPS/PIP ₂ TxR® (79:15.5:1)	6 μΜ
croscopy ¹⁵⁴		
Fluorescence mi-	DOPC/DOPE/PIP ₂ (60:30:10)	0.5 ± 0.2 μM
croscopy ⁷⁹		

Abbrevations and symbols

Angström

A_M Area per molecule

 $A_{\rm M,20}$ Area per molecule at 20 mN/m

*A*_v Area of a free giant unilamellar vesicle

 $A_{\rm ad}$ Are of an adhered giant unilamellar vesicle

AFM Atomic force microscopy

b-PEG Biotin polyethylene glycol

b-PEG-NH₂ Biotin polyethylene glycol amine

BSA Bovine serum albumin

C Concentration

CCV Clathrin-coated vesicles

Cf. compare

CG-MD Coarse-grained molecular dynamics
CLSM Confocal laser scanning microscopy

CME Clathrin-mediated endocytosis

D Diffusion coefficientd Cuvette thickness

 d_{M} Physical membrane thickness d_{ENTH} Physical thickness of ENTH

DHPE 1,2 Dihexadecanoyl-*sn*-glycero-3-phosphoethanolamin

DIPS 1,2-dilinoleoyl-*sn*-glycero-3-phospho-*L*-serine

DPPE 1,2-dipalmitoyl-*sn*-glycero-3-phosphoethanolamine

DOPC 1,2-dioleoyl-*sn*-glycero-3-phosphocholine

DOPE 1,2-dioleoyl-*sn*-glycero-3-phosphoethanolamine

DOPS 1,2-dioleoyl-*sn*-glycero-3-phospho-*L*-serine

DPPS 1,2-dipalmitoyl-*sn*-glycero-3-phospho-*L*-serine

EDTA Ethylenediaminetetraacetic acid

EGFR Epidermal growth factor receptor

ENTH Epsin *N*-terminal Homology

EPR Electron paramagnetic resonance

F Force

*F*_m Mobile fraction

FRAP Flourescence recovery after photobleaching

FWHM Full width at half maximum

GOPTS (3-Glycidyloxypropyl)trimethoxysilane

GUV Giant unilamellar vesicle

I Intensity h Height h Hour H_2O Water

H₂O₂ Hydrogen peroxide

HMDS Hexamethyldisilane

HEPES (4-(2-hydroxyethyl)-1-piperazine)-ethanesulfonic acid

Hz Hertz

ITO Indium tin oxidek Spring constant

*K*_A Area compressibility modulus

KCl Potassium chloride

*K*_D Dissociation constant

kDa Kilodalton

KH₂PO₄ Monopottasium phosphate

LC Liquid-condensed
LE Liquid-expanded

LUV Large unilamellar vesicle

M Molecular mass

MeOH Methanol

MgCl₂ Magnesium chloride

Min Minute

MLV Multilamellar veicles

m-PEG Methoxy polyethylene glycol

m-PEG-NH2 Methoxy polyethylene glycol amine

n Refractive index, number of measurements

NA Numerical aperture

nm Nanometer

Na Sodium

NaCl Sodium chloride

Na₂HPO₄ Disodium phosphate

NaN₃ Sodium azide

NH₃ Ammonia

OT Optical thickness

PBS Phosphate buffered saline

PEG Polyethylene glycol

PEM Photoelastic modulator

PIP₂ Phosphatidylinositol-4,5-bisphosphate

POPC 1-Palmitoyl-2-oleoyl-*sn*-glycero-3-phosphocholine

POPE 1-Palmitoyl-2-oleoyl-*sn*-glycero-3-phosphoethanolamine

POPS 1-Palmitoyl-2-oleoyl-*sn*-glycero-3-phosph

*r*_{ad} Radius of an adhered giant unilamellar vesicle

 r_i Radius of the contact area

 $r_{\rm v}$ Radius of a free giant unilamellar vesicle

 $r_{1/2}$ Fresnel coefficients

RIfS Reflectometric interference ppectroscopy

rms Root mean square ROI Region of interest

rpm Revolutions per minute

RT Room temperature

s Second

s Traveling distance

SAM Self assembled monolayer

SD Standard deviation

SDCLM Spinning disc confocal laser microscopy

SLB Supported lipid bilayer

SiO₂ Silicon dioxide

SPR Surface plasmon resonance

SUV Small unilamellar vesicle

t Time

T Temperature

TiO₂ Titanium dioxide

U Voltage

UV Ultraviolett

V Volt

V Volume Vis Visible

wt Wild type

Zc Vertical deflection
Zp Piezo position

 α Angle of incident light

 γ Surface tension δ Phasendifferenz Bending rigidity

 φ Angle of refracting light

heta Contact angle λ Wavelength

 σ Membrane tension π Surface pressure

 π_{c} Critical surface pressure π_{0} Initial surface pressure

χ Molar fraction

Materials

Braungläschen OCHS GmbH, Göttingen, GER

Cantilever

Eppendorf cups Eppendorf, Hamburg, GER Eppendorf pipettes Eppendorf, Hamburg, GER

Hamilton syringe ThermoScientific, Braunschweig, GER

ITO-slides Präzisions Glas & Optik GmbH, Isersloh, GER

Klemmen ALCO, Arnsberg, GER

Copper stripes Präzisions Glas & Optik GmbH, Isersloh, GER

Menzel-Objektträgergläschen ThermoScientific, Braunschweig, GER Parafilm American National Can, Chicago, GER

Petri dishes Sarstedt, Nürnbrecht, GER
Pipette tips Sarstedt, Nürnbrecht, GER

Pump tubes Tygon Ismatec, Wertheim-Mondfeld, GER

RIfS chamber custom-built, OC Werkstatt, Göttingen, GER

SiO₂ substrates Sitron Inc, KOR

Test tubes

VWR International, Darmstadt, GER

Chemicals

Aceton Carl Roth, Karlsruhe, GER

Dylight® 594 NeutrAvidin Thermo Fisher, Waltham, USA

Aluminiumoxid Sigma-Aldrich, Taufkirchen, GER

Ammonia solution (25%) VWR International, Darmstadt, GER

Argon Linde, München, GER

Biotin-PEG-NH₂ Thermo Fisher, Waltham, USA

BSA. Protease free Carl Roth, Karlsruhe, GER

Chloroform VWR International, Darmstadt, GER

Cholesterol Sigma-Aldrich, Taufkirchen, GER

DOPC Avanti Polar Lipids, Alabaster, USA

DOPE Avanti Polar Lipids, Alabaster, USA

DOPS Avanti Polar Lipids, Alabaster, USA

Ethanol, p.a. VWR International, Darmstadt, GER

Glucose Carl Roth GmbH, Karlsruhe, GER

GOPTS Sigma Aldrich, Taufkirchen, GER

HEPES Carl Roth GmbH, Karlsruhe, GER

H₂O₂ (30%) Grüssing GmbH, Filsum, GER

KCl Carl Roth GmbH, Karlsruhe, GER

KH₂PO₄ Merck, Darmstadt, GER

KOH Merck, Darmstadt, GER

Methanol VWR International, Darmstadt, GER

MgCl₂ · 6 H₂O Merck, Darmstadt, GER

Mucasol® Sigma-Aldrich, Taufkirchen, GER

NaCl Merck, Darmstadt, GER

Na citrate Merck, Darmstadt, GER

NaN₃ Merck, Darmstadt, GER

Na₂HPO₄ VWR International, Darmstadt, GER

NeutrAvidin Thermo Fisher, Waltham, USA

Nitrogen Linde, München, GER

Parafilm Pechiney Plastic Packaging, Chicago, USA

Pipette tips Sarstedt, Nürnbrecht, GER

PIP2 Avanti Polar Lipids, Alabaster, USA
POPC Avanti Polar Lipids, Alabaster, USA
POPE Avanti Polar Lipids, Alabaster, USA
POPS Avanti Polar Lipids, Alabaster, USA

Hydrochlorid acid Merck, Darmstadt, GER

Sucrose Carl Roth GmbH, Karlsruhe, GER
TxR Sigma-Aldrich, Taufkirchen, GER

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2006-2009 Allgemeine Hochschulreife at the Gymnasium Athenaeum

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Scientific contributions

	publications
2019	Kroppen, B, <u>Teske, N.</u> , Kondratiuk, I., Denkert, N., Mukherjee, I.,
	Yambire K. F., Jaipuria, G, Zweckstetter, M., Milosevic, I., Steinem,
	C., Meinecke M. (2019) Lipid-dependent oligomerization of the
	epsin1 ENTH domain drives membrane deformation and is im-
	portant for clathrin-mediated endocytosis, in preparation.
2017	Teske, N. & Sibold, J., Schumacher, J., Teiwes, N. K., Gleisner, M.,
	Mey, I., Steinem, C. (2017) Continuous pore-spanning lipid bi-
	layers on silicon oxide-coated porous substrates. Langmuir
	2017 , 33, 14175–14183.
2016	Gleisner, M., Kroppen, B., Fricke, C., <u>Teske, N.</u> , Kliesch, TT., Jans-
	hoff, A., Meinecke, M., Steinem, C. (2016) Epsin N-terminal ho-
	mology domain (ENTH) activity as a function of membrane ten-
	sion. J. Biol. Chem. 2016 , 291, 19953-19961.
2010	oral presentations
2019	SFB 803 Winter school. <i>Lipid dependent ENTH-membrane inter-</i>
221	actions.
2017	SFB 803 Autumn school. Continuous pore-spanning bilayers on
	silicon oxide coated porous substrates.
	Selected poster presentations
2018	Biennial Meeting of the German Biophysical Society, Düsseldorf.
	Investigation of biomimetic membrane interactions with the
	ENTH domain of epsin 1 and substrate surfaces.

2017	19th IUPAB and 11th EBSA Congress, Edinburgh. <i>Interactions of</i>
	biomimetic membrane models with the ENTH domain of epsin 1
	and substrate surfaces.
2016	International Symposium "Biological Dynamics from Micro-
	scopic to Mesoscopic Scales", Grimma. Measurements of ENTH ac-
	tivities in lipid membrane systems.

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